

## Exercises 7

(to hand in: January 7, 2015 in the lecture)

In the following, we say that a state  $\rho$  is *diagonal* on a quantum system  $S$  with Hamiltonian  $H_S$  if and only if  $[\rho, H_S] = 0$ . This is equivalent to the fact that there is some (energy) eigenbasis of  $H_S$  such that the matrix representation of  $\rho$  in this basis is a diagonal matrix. The *Gibbs state at inverse temperature*  $\beta$  is  $\gamma_\beta := \exp(-\beta H_S)/Z$ , where  $Z \in \mathbb{R}$  is such that  $\text{tr}(\gamma_\beta) = 1$ .

Problem 26 (Gibbs states are completely passive): (8 points)

In the lecture, we have learned that a state  $\rho_S$  is passive (with respect to some Hamiltonian  $H_S$ ) if and only if  $\rho_S = \text{diag}(p_1, \dots, p_n)$  in some eigenbasis of  $H_S$ , and  $E_i > E_j \Rightarrow p_i \leq p_j$ . A state  $\rho_S$  is called completely passive if  $\rho_S^{\otimes n}$  is passive for all  $n$ . (In this statement, the Hamiltonian on the system  $S^{\otimes n}$  is taken to be  $H_S \otimes \mathbb{1} \otimes \dots \otimes \mathbb{1} + \mathbb{1} \otimes H_S \otimes \mathbb{1} \otimes \dots \otimes \mathbb{1} + \dots + \mathbb{1} \otimes \dots \otimes \mathbb{1} \otimes H_S$ .) Show that all Gibbs states  $\gamma_\beta$  (for all inverse temperatures  $\beta \geq 0$ ) are passive and completely passive.

Problem 27 (Passive and completely passive states in small dimensions): (8 points)

- (i) Argue that every density matrix of full rank can be seen as Gibbs state if you choose the Hamiltonian correctly. Or stated differently, every density matrix can be written as  $\rho = e^{-\beta H}/Z$  if  $H$  and  $\beta \geq 0$  are chosen appropriately. (You don't have to be totally mathematically rigorous here.)
- (ii) Show that for a *fixed* diagonal Hamiltonian  $H$  of size two, every diagonal density matrix of full rank is a Gibbs state for some appropriate inverse temperature  $\beta \geq 0$ . Use Exercise 26 to conclude that all passive states in two dimensions are completely passive.
- (iii) Find an example in three dimensions of a state (and a corresponding Hamiltonian) that is passive but not completely passive.

Problem 28 (General properties of thermal operations): (6 points)

Consider thermal operations  $\Phi : S \rightarrow S$ , where  $S$  is a finite-dimensional quantum system.

- (i) Show that all thermal operations leave the Gibbs state of the system invariant, i.e.  $\Phi(\gamma_S) = \gamma_S$ .
- (ii) Show that the image of any diagonal state under thermal operations is diagonal. Thus, no superpositions of energy eigenstates can be created by thermal operations.

Problem 29 (Transitions with a small environment):

(8 points)

Consider a two-dimensional quantum system  $S$  and a two-dimensional heat bath  $B$  with respective Hamiltonians

$$H_S = \begin{pmatrix} 0 & 0 \\ 0 & \Delta E \end{pmatrix} \quad H_B = \begin{pmatrix} E & 0 \\ 0 & E + \Delta E \end{pmatrix}.$$

Let  $\beta \geq 0$  be some arbitrary but fixed inverse temperature.

- (i) Show that the Gibbs states of the system and the bath are identical, i.e.  $\gamma_S = \gamma_B$ .
- (ii) Let the system be in the diagonal state  $\rho_S = \text{diag}(p, 1 - p)$ , where  $0 \leq p \leq 1$ . Show that under thermal operations the transition  $\rho_S \mapsto \gamma_S$  is possible, and give a unitary that (together with a partial trace) achieves this transition.
- (iii) What other state transitions are possible with this environment? Give the set of achievable states as a subset of the set of two-component probability vectors. (Of course the states are all matrices, but since they are all diagonal (cf. Exercise 28), we can represent them as probability vectors.)

Hint: Argue that we obtain transitions from  $\text{diag}(p, 1 - p)$  to  $\text{diag}(p', 1 - p')$ , and that

$$\begin{pmatrix} p' \\ 1 - p' \end{pmatrix} = B \begin{pmatrix} p \\ 1 - p \end{pmatrix}$$

for some matrix  $B$ . Use known properties of  $B$  to determine all possible  $p'$ .