



Testing quantum theory by generalizing noncontextuality

Markus P. Müller^{1,2,3} and Andrew J. P. Garner^{1,2}

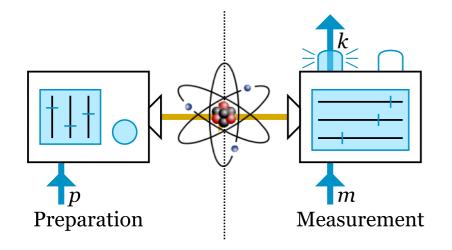
¹Institute for Quantum Optics and Quantum Information (IQOQI), Vienna ²Vienna Center for Quantum Science and Technology (VCQ), Vienna ³Perimeter Institute for Theoretical Physics (PI), Waterloo, Canada





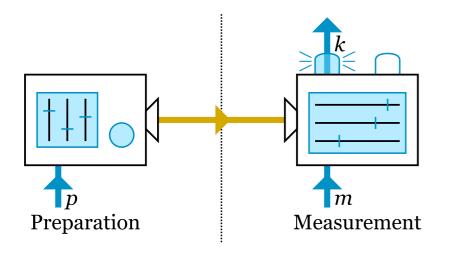


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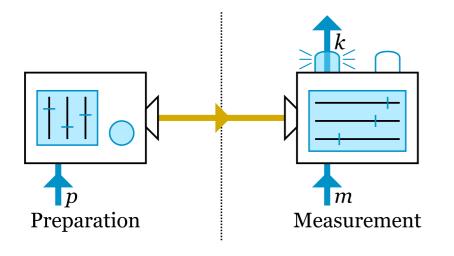
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Could the resulting data **falsify QT** w/o assumptions on devices or physics?

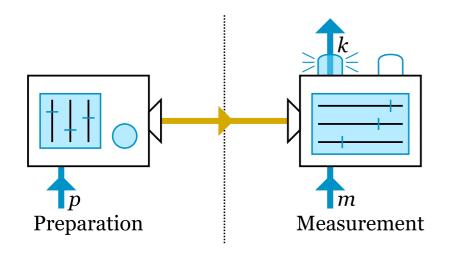
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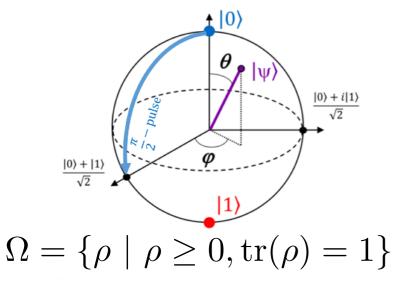
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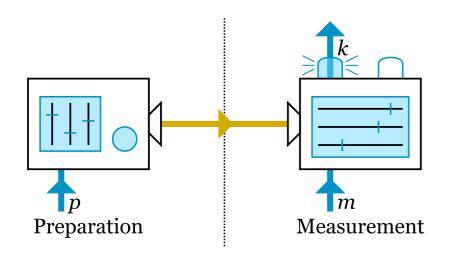


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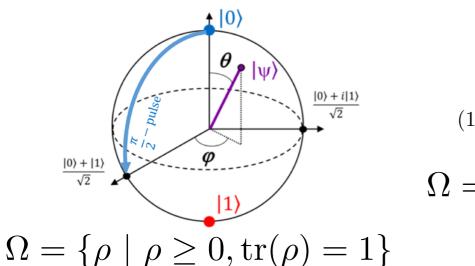


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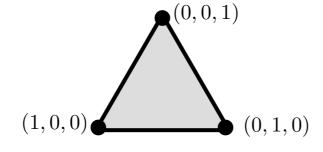


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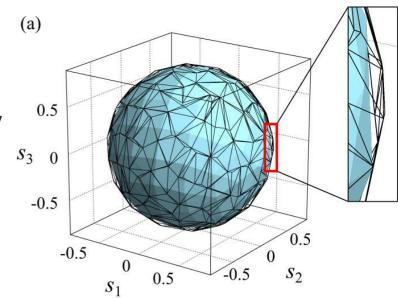


$$\Omega = \{ p = (p_1, \dots, p_n) \mid p_i \ge 0, \sum p_i = 1 \}$$

- noisy qubits etc.
- QT w/ superselection rules
- ... ?

Overview

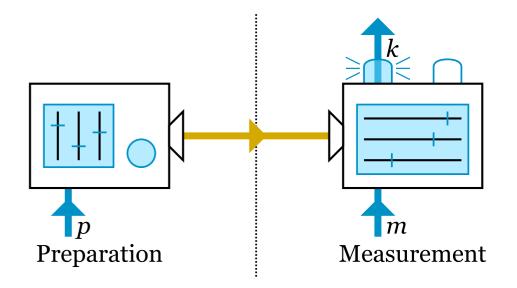
1. GPTs and theory-agnostic tomography



2. Contextuality, simulations, and embeddings

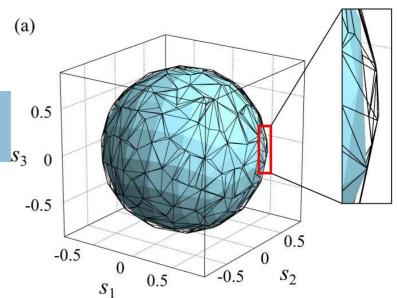
3. Exact embeddings into quantum theory

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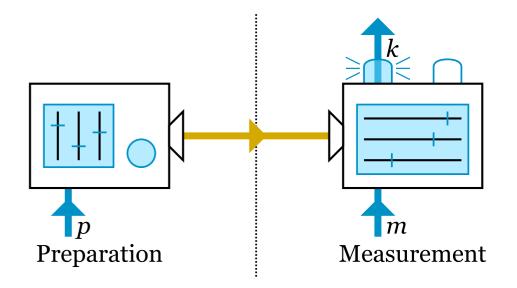
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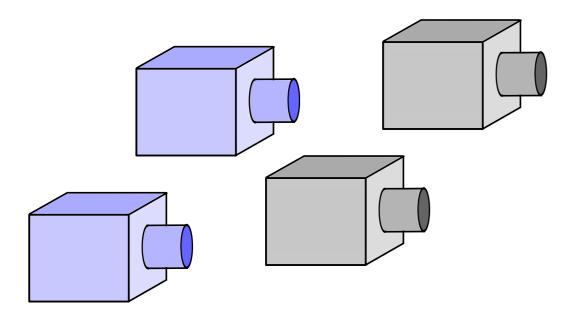


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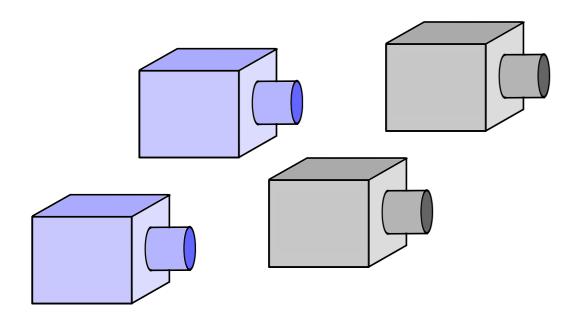
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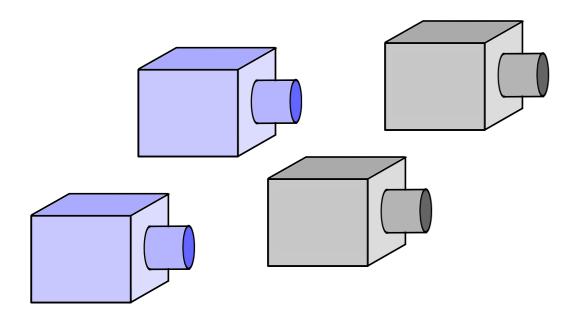


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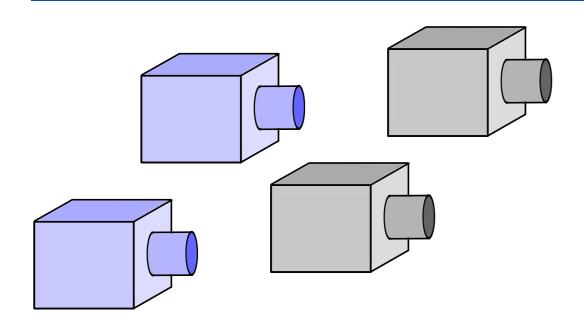
 $P_1 \sim P_2$ if they give identical probabilities for all outcomes of all accessible measurements.



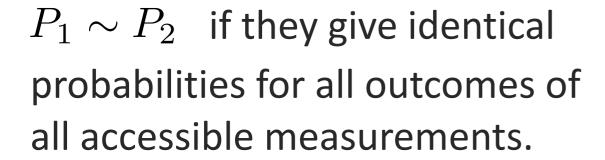
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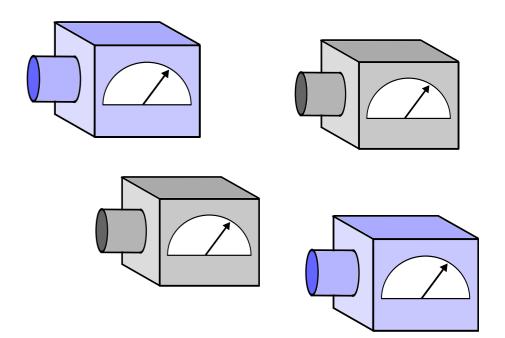
State ω_P = equivalence class of preparation procedures



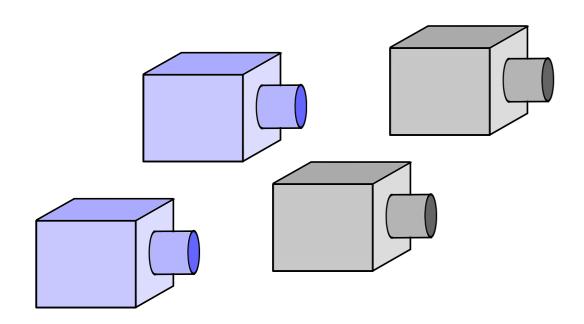




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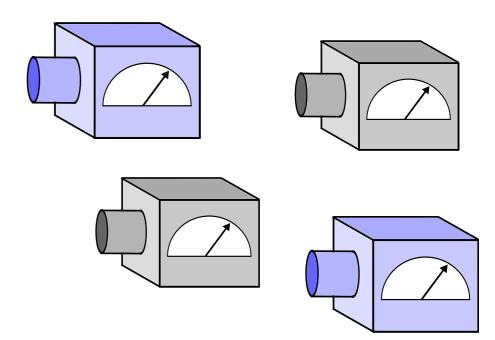
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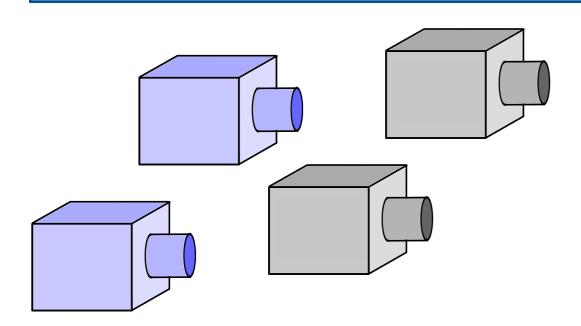
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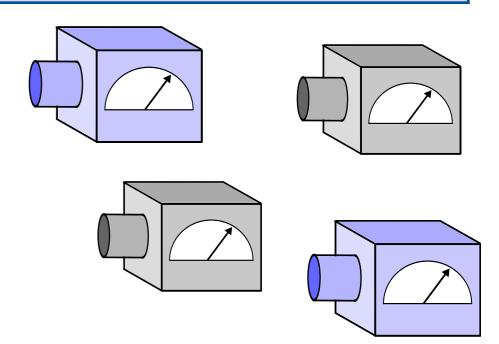
$$(k_1,M_1) \sim (k_2,M_2)$$
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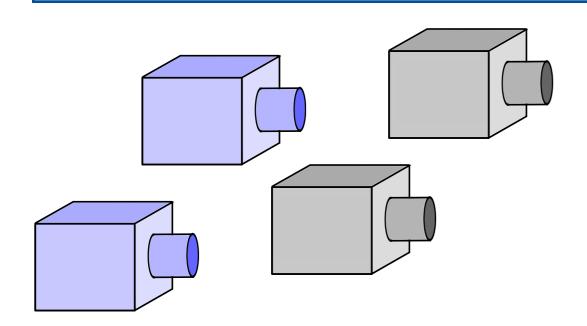
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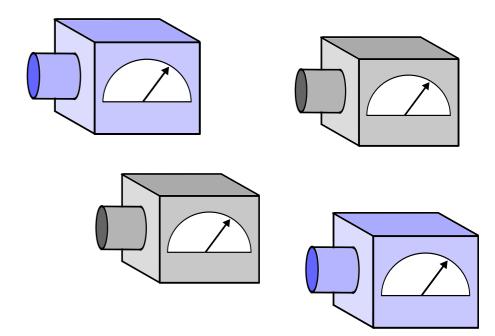
Effect $e_{k,M}$ = equivalence class of outcome-measurement pairs





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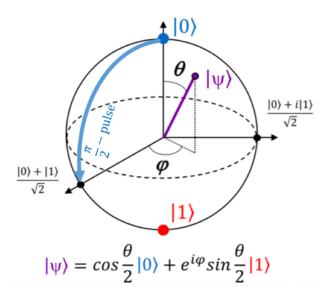
GPT $\mathcal{A} = (A, \Omega_A, E_A)$ = (vector space over \mathbb{R} , normalized states, effects).

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Quantum theory (QT): Q_n

$$A = \mathbb{H}_n(\mathbb{C})$$
 (complex Hermitian $n \times n$ matrices)

$$E_A = \{E \mid 0 \le E \le 1\}$$
 (POVM elements)

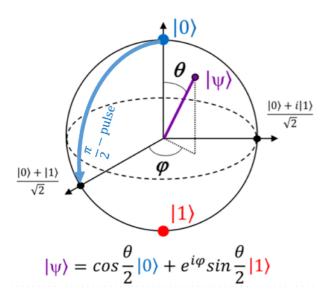
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$$A^* \simeq A$$
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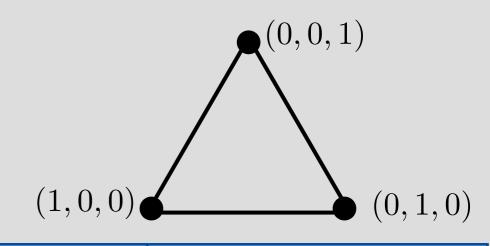
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Classical probability theory (QT): \mathcal{C}_n

$$A = \mathbb{R}^n \simeq A^*$$

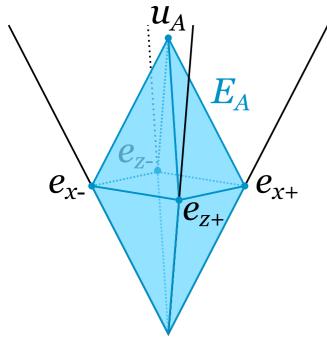
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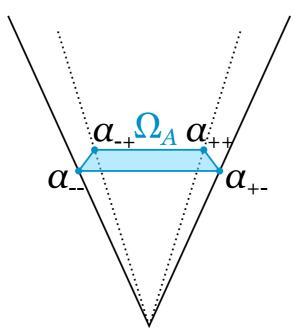


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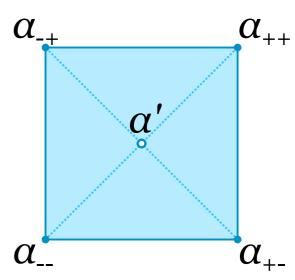
The gbit
$$\mathcal{A}=(\mathbb{R}^3,\Omega_A,E_A)$$



a) Cone of effects A_+

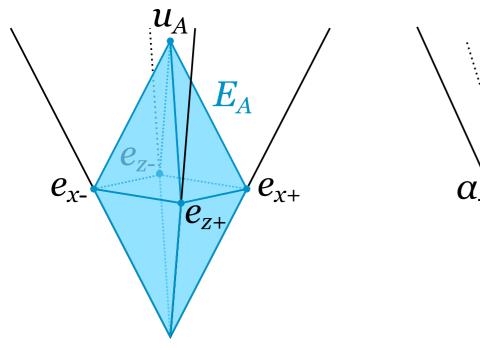


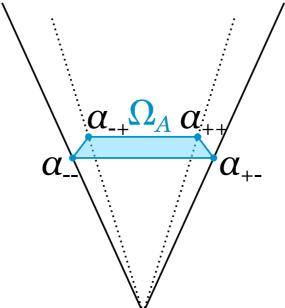
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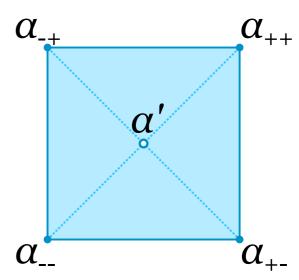


c) Normalized states Ω_A

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The four pure states $\alpha_{\pm,\pm}$ are **pairwise** perfectly distinguishable, but **not jointly** \Longrightarrow this cannot be a classical or quantum system.

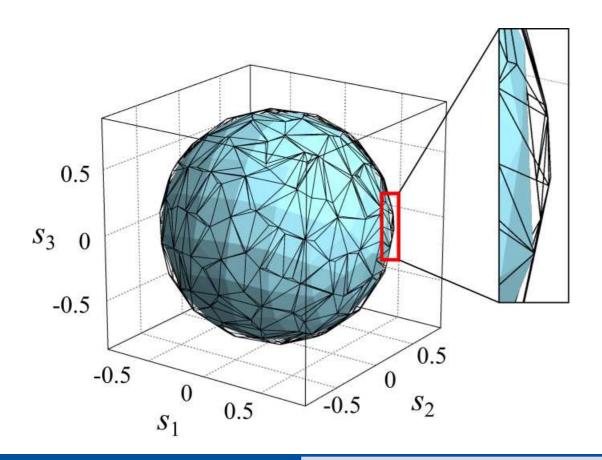
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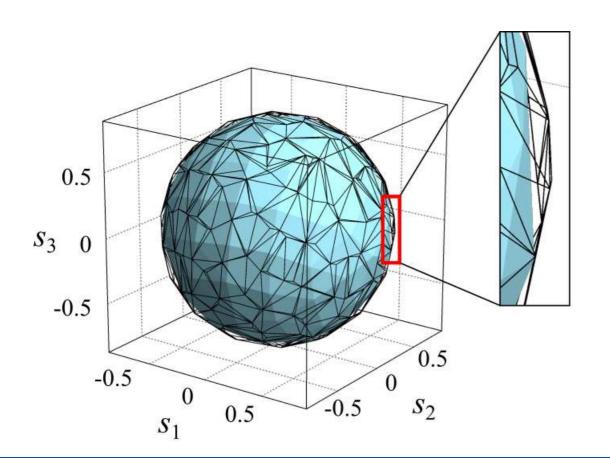


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Tomographic completeness loophole: can never be sure that we probed the system *completely*.

1. Theory-agnostic tomography

What if we just see a (low-dimensional) "shadow"?

Let's **drop** the **tomographic completeness** assumption.

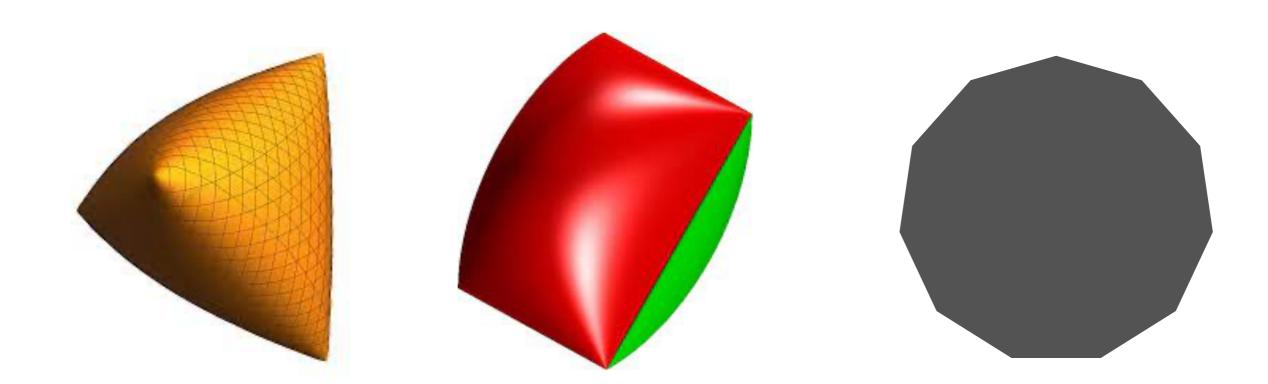
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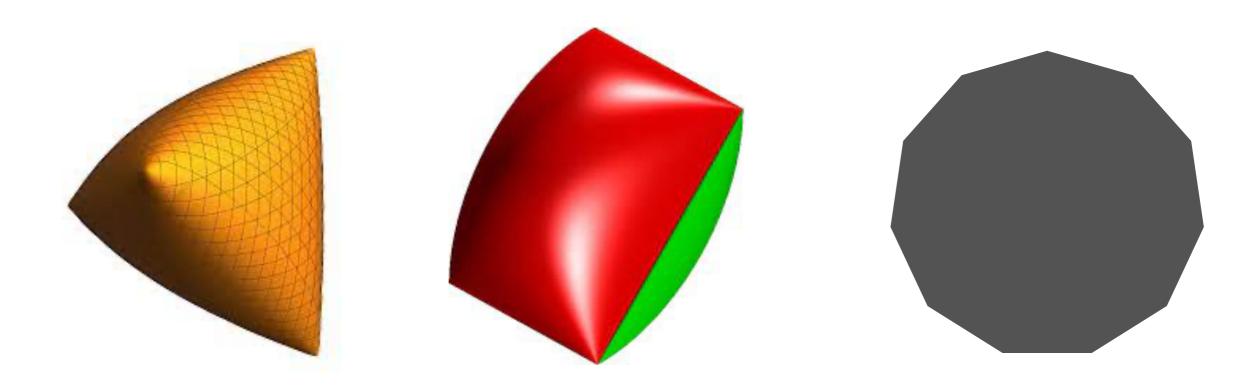


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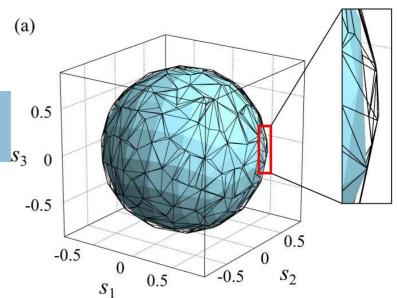
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Is **fundamental QT** a plausible explanation of a given **effective GPT**?

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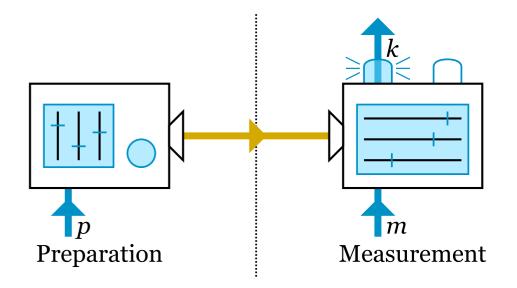
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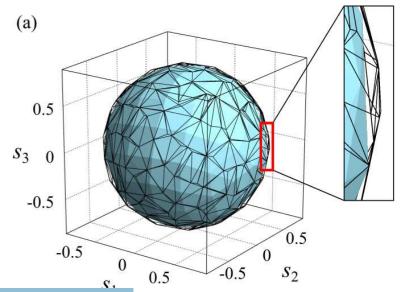
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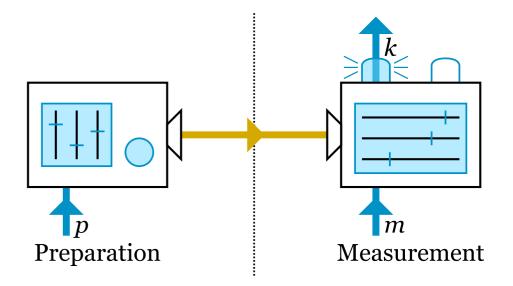
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Contextuality for preparations, transformations and unsharp measurements

R. W. Spekkens*

Perimeter Institute for Theoretical Physics, 35 King St. North, Waterloo, Ontario N2J 2W9, Canada (Dated: Feb. 25, 2005)



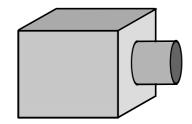
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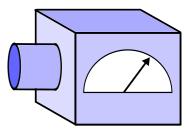
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Recall the notion of an operational theory.





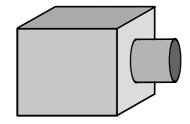
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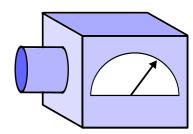
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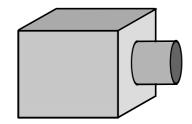
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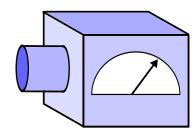
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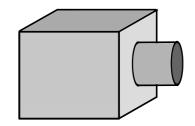
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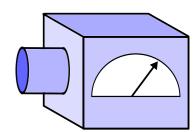
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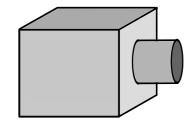
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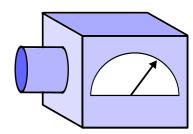
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2. Simulations, embeddings, ...

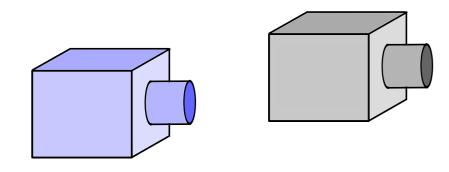
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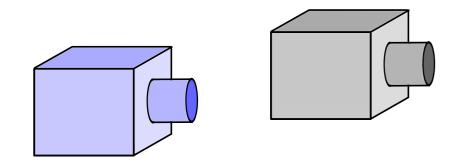
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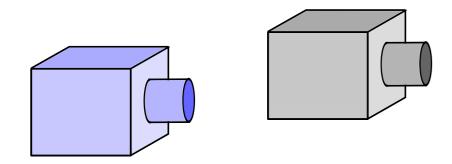


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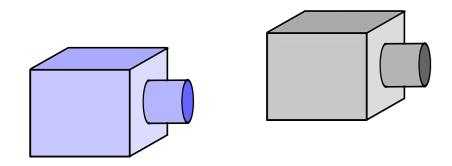
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Theorem: Ontological models of QM-systems must be preparation-contextual (and, assuming outcome-determinism for sharp meas., measurement-contextual).

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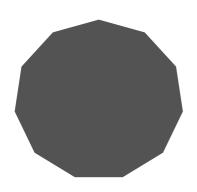
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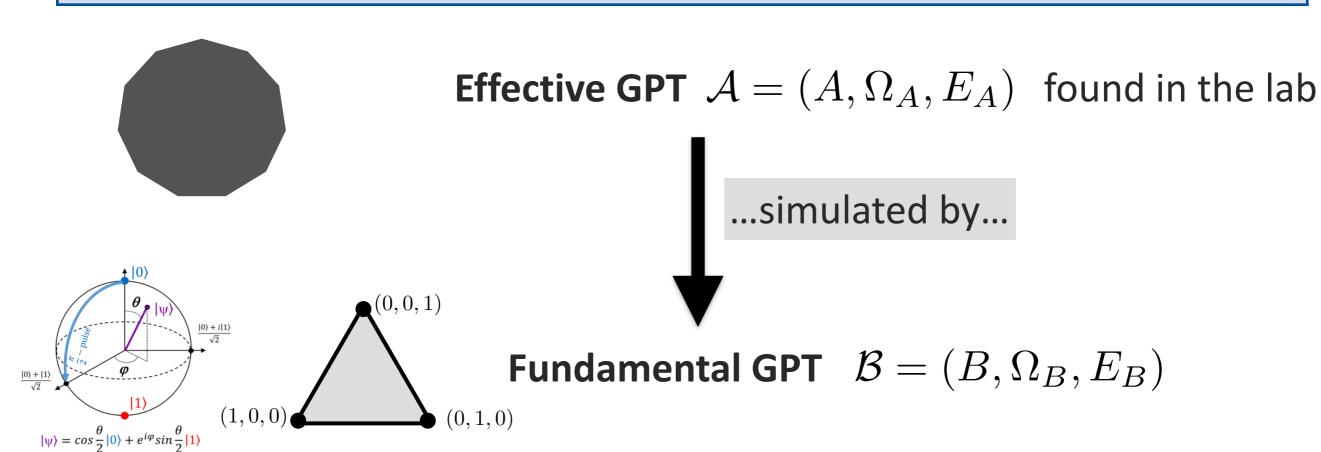
Intuition: Contextual models are implausible because they are **fine-tuned**: operationally, $P \sim P'$, but ontologically, $\mu_P \neq \mu_{P'}$.

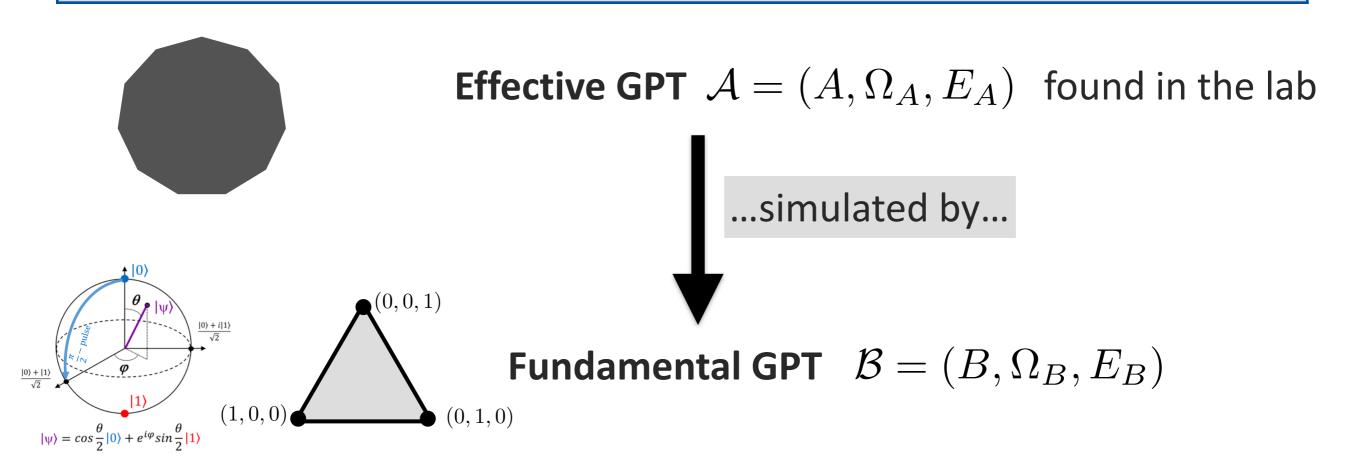
An instance of Leibniz' principle of the "identity of the indiscernibles".

2. Simulations, embeddings, ...



Effective GPT $\mathcal{A}=(A,\Omega_A,E_A)$ found in the lab





Effectively preparing state ω_A means **fundamentally** preparing some ω_B , but ω_B may depend on the preparation *procedure*, i.e. the *context*. Collect all those states into a set $\Omega_B(\omega_A) := \{\omega_B\}$.

2. Simulations, embeddings, ...

Definition. An arepsilon -simulation of effective GPT $\mathcal A$ by fundamental GPT $\mathcal B$:

Effective state ω_A \longrightarrow set of simulating states $\Omega_B(\omega_A)$,

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such that all outcome probabilities are reproduced up to ε :

$$|(\omega_A, e_A) - (\omega_B, e_B)| \le \varepsilon$$
 for all $\omega_B \in \Omega_B(\omega_A), e_B \in E_B(e_A)$

and, essentially, mixtures are valid simulations of mixtures (see paper).

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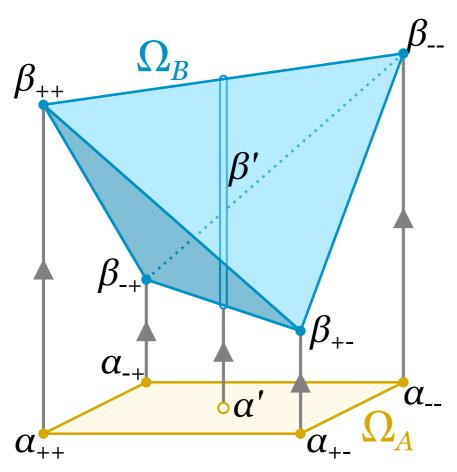
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Special case A = QT, B = classical probability theory:

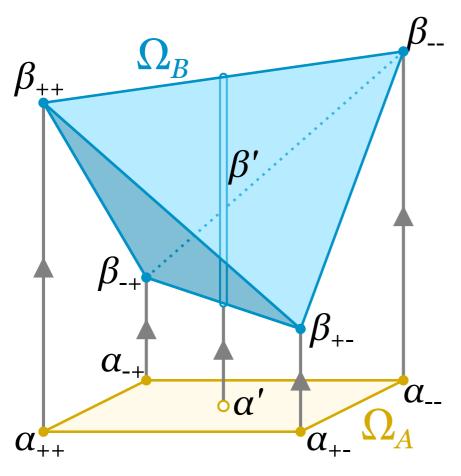
Simulations are **ontological models**, and univalence = **noncontextuality**.

2. Simulations, embeddings, ...

Example ("Holevo projection"): simulating the gbit $\mathcal{A} = (\mathbb{R}^3, \Omega_A, E_A)$ with a classical 4-level system \mathcal{C}_4 .



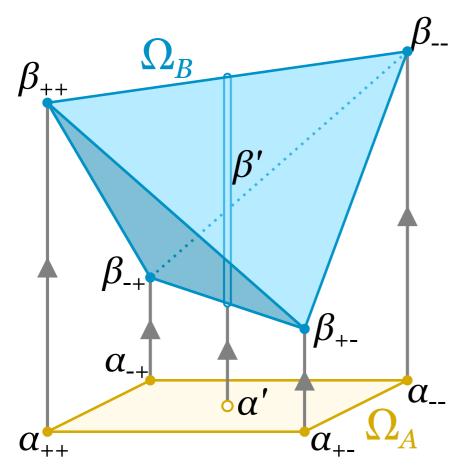
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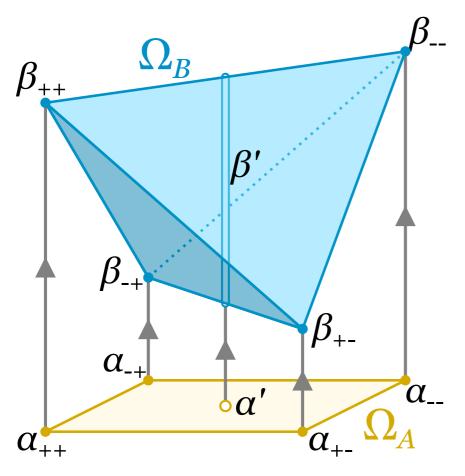
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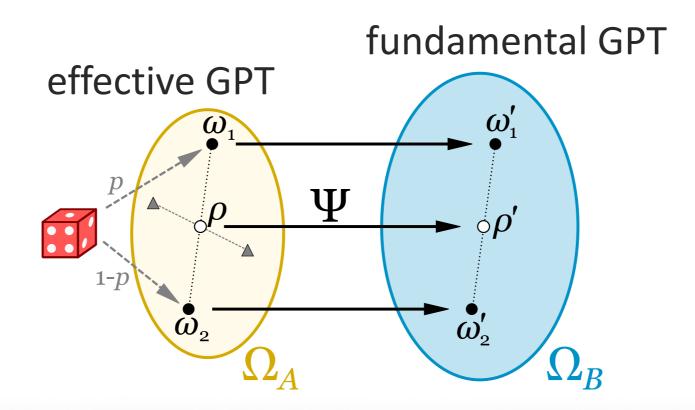
(Preparation) contextuality = multivalence:

the fundamental state β' does not only depend on α' , but *must* also depend on the way it has been prepared.

This is an instance of implausible fine-tuning: the statistical differences among the fundamental states are miraculously *exactly "washed out"* on the effective level.

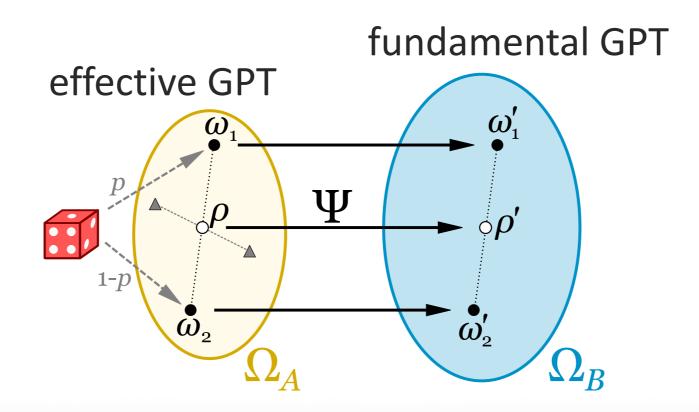
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An ε -embedding consists of two linear maps Ψ and Φ such that

- Ψ maps the normalized states of A into those of B,
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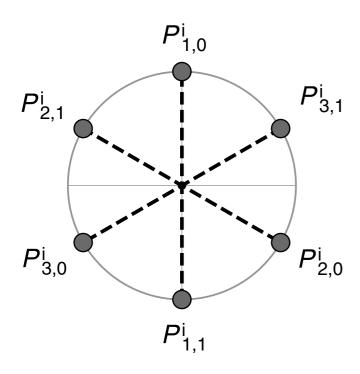
2. Simulations, embeddings, ...

Noncontextual inequalities and approximate embeddings

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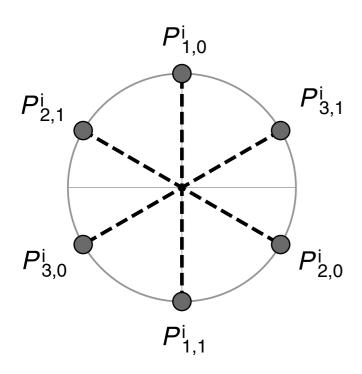
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Quantitative statement:

$$A := \frac{1}{6} \sum_{t \in \{1,2,3\}} \sum_{b \in \{0,1\}} P(b \mid p_{t,b}, m_t) \le \frac{5}{6}.$$

Noncontextual inequalities and approximate embeddings

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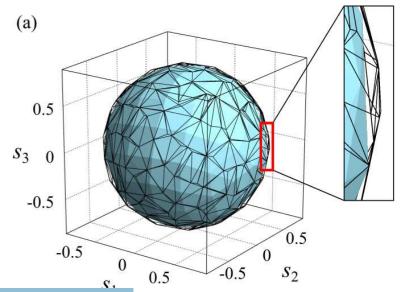
$$P_{2,0}^{i}$$
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These imply bounds on the approximate embeddability into classical:

Example 1. Let $\varepsilon < \frac{1}{6}$. Then the rebit (and thus, also the qubit) cannot be ε -embedded into any C_n .

Overview

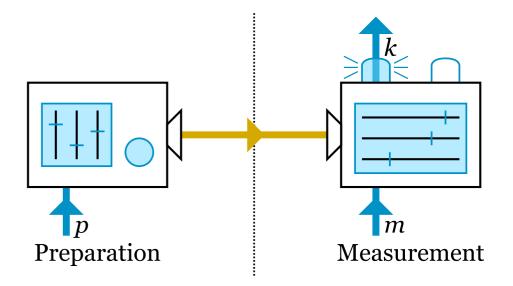
1. GPTs and theory-agnostic tomography



2. Contextuality, simulations, and embeddings

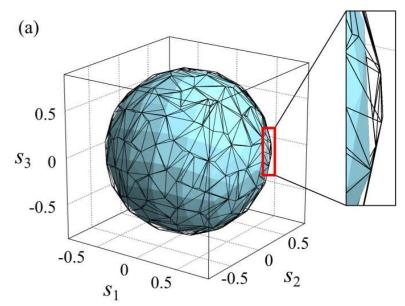
3. Exact embeddings into quantum theory

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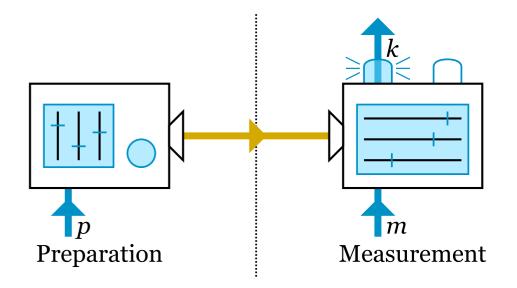
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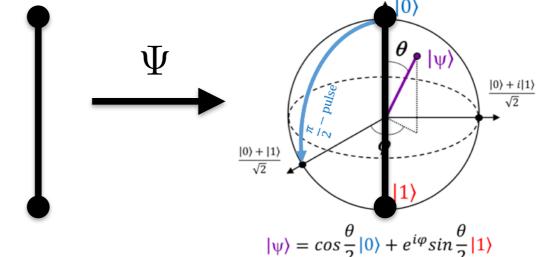
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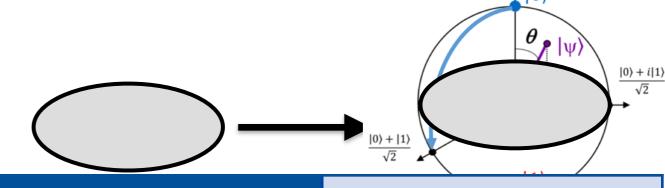
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$$(e_1, \dots, e_n) \xrightarrow{\Phi} \left(\begin{array}{ccc} e_1 & \dots & 0 \\ 0 & \ddots & 0 \\ 0 & \dots & e_n \end{array}\right).$$



$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1$$

Similarly, QT over the real numbers can be embedded into QT.



3. Exact embeddings into QT

Focus on the "unrestricted GPTs" where all vectors yielding valid probabilities on all states are effects: $\mathcal{A} = (A, \Omega_A, E_A)$

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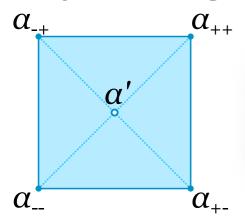
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Non-exact embeddings into quantum theory

Example: the gbit

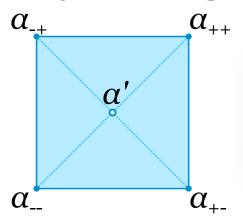


Example 2. Let $\varepsilon \leq 0.1014$. Then the gbit cannot be ε -embedded into any \mathcal{Q}_n or \mathcal{Q}_{∞} .

There is no better-than-10% univalent ("noncontextual") simulation by QT.

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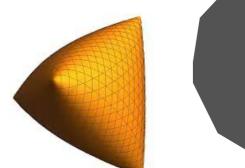
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Also shown in our paper:

can use known results on Bell inequalities to certify nonembeddability.

Impractical and inefficient, but "proof of principle".



Multivalence / nonembeddability / contextuality is hard to obtain

We prove a couple of results:

- If A arises from B via coarsegraining, noise, or generalized decoherence, then A can be exactly embedded into B.
- Embeddability is **transitive**: If $A \mapsto B \mapsto C$ then $A \mapsto C$.

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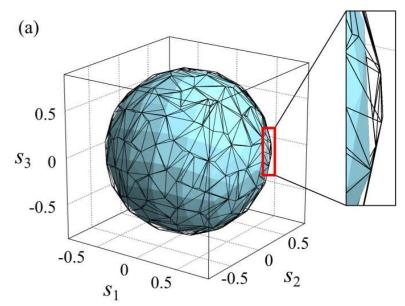
Example consequences:

- Any future physical theory that decoheres to QT must also be contextual.
- "Statistical-mechanics-like" fundamental theories correspond to univalent ("noncontextual") simulations.

3. Exact embeddings into QT

Overview

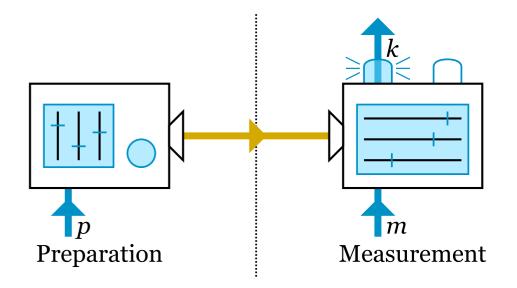
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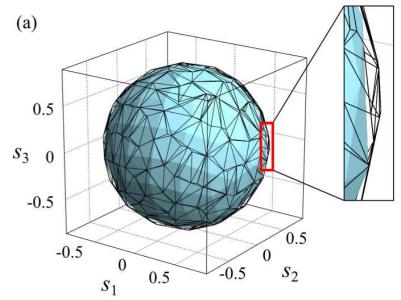
3. Exact embeddings into quantum theory

4. An experimental test of QT



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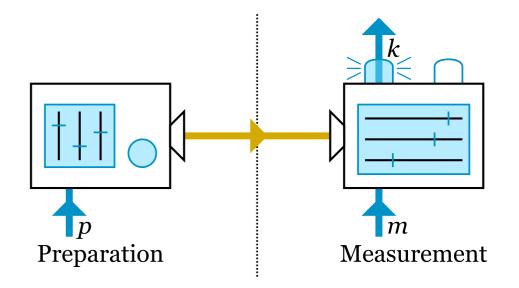
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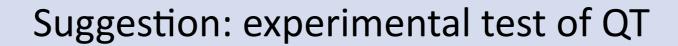


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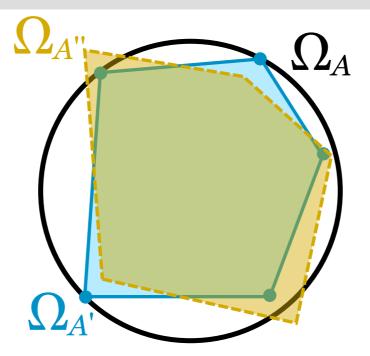


Suggestion: experimental test of QT

Perform theory-agnostic tomography on an effective physical system in your laboratory.

Test whether the resulting effective GPT is $\varepsilon-$ embeddable into QT, where ε is a function of the experimental uncertainty.

If, surprisingly, "no", then this challenges QT.

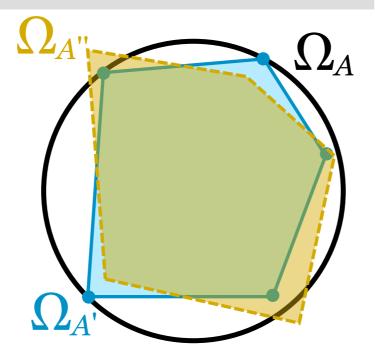


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A quantum explanation of the result is then similarly implausible as a classical (contextual) explanation of the quantum state space.

What is a physical system?

Similarly as

V. Gitton and M. P. Woods, arXiv:2209.04469

we think that, <u>for our purpose</u>, we should not think of physical systems as "spatially localized objects in the world", but rather as being **defined by an experimental scenario**.

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The implications of this for **contextuality** are interesting and to be debated (thanks to Rob Spekkens, David Schmid, and others so far).

Summary

Have generalized Spekkens' notion of generalized noncontextuality:
 "Processes that are statistically indistinguishable in an effective theory
 should not require explanation by multiple distinguishable processes
 in a more fundamental theory."

 Result: Several mathematical insights, a new experimental test of QT, a conceptual discussion of the notion of a "physical system".
 Looking forward to discussions and comments.

arXiv:2112.09719 (update soon)

Thank you!