

Testing quantum theory by generalizing noncontextuality

Markus P. Müller^{1,2,3} and Andrew J. P. Garner^{1,2}

¹Institute for Quantum Optics and Quantum Information (IQOQI), Vienna

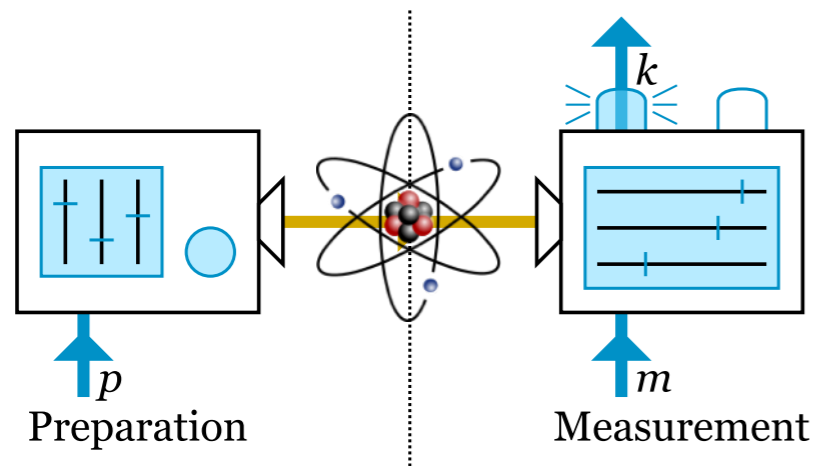
²Vienna Center for Quantum Science and Technology (VCQ), Vienna

³Perimeter Institute for Theoretical Physics (PI), Waterloo, Canada



Two motivations

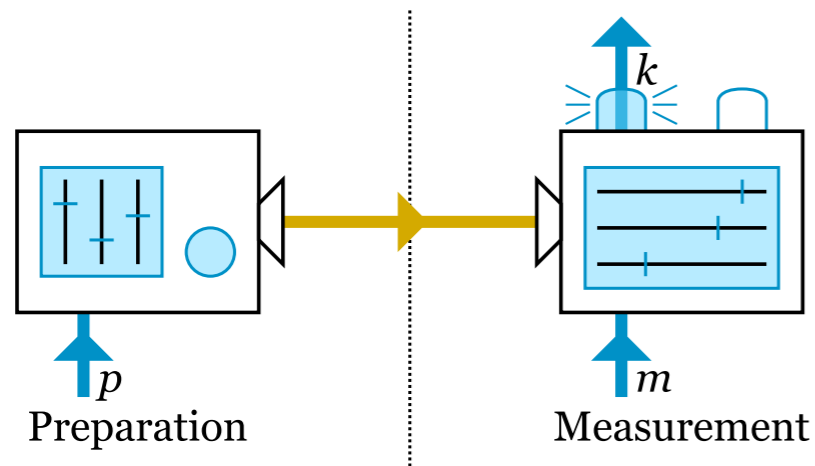
Suppose we **prepare** and **measure** a physical system in all ways accessible to us.



Could the resulting data **falsify QT**?

Two motivations

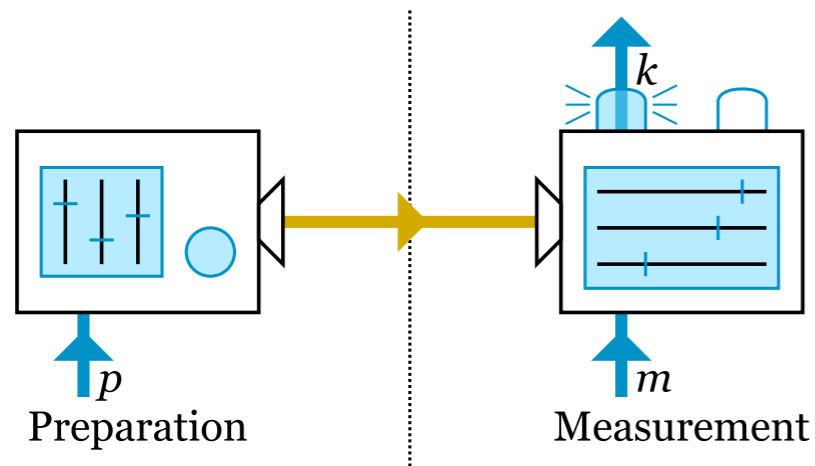
Suppose we **prepare** and **measure** a physical system in all ways accessible to us.



Could the resulting data **falsify QT** w/o assumptions on devices or physics?

Two motivations

Suppose we **prepare** and **measure** a physical system in all ways accessible to us.

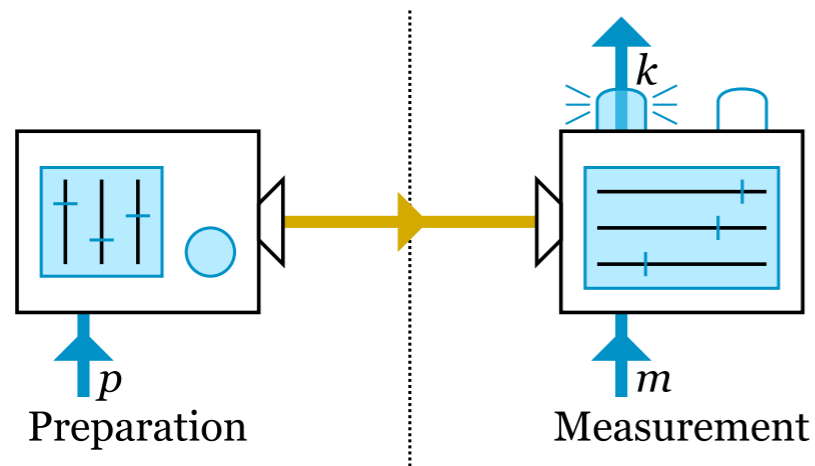


Could the resulting data **falsify QT** w/o assumptions on devices or physics?

If Nature is **fundamentally quantum**, which **effective probabilistic theories** can we reasonably expect to encounter?

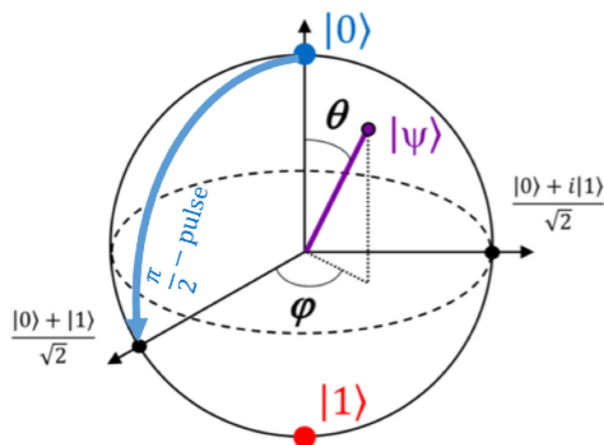
Two motivations

Suppose we **prepare** and **measure** a physical system in all ways accessible to us.



Could the resulting data **falsify QT** w/o assumptions on devices or physics?

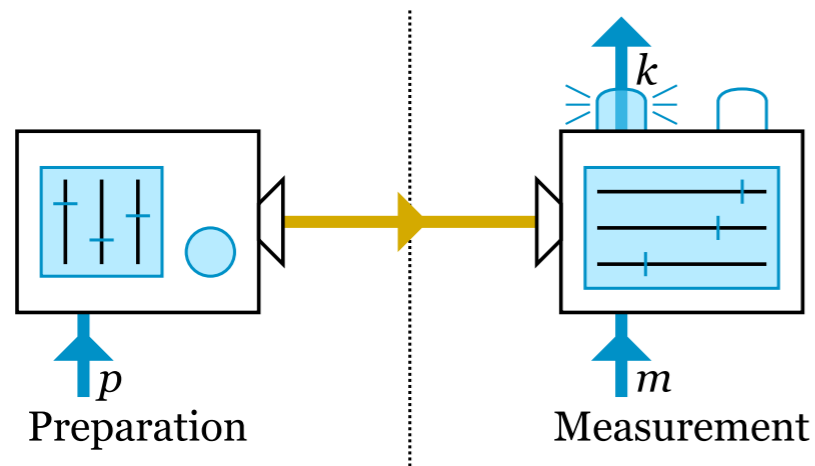
If Nature is **fundamentally quantum**, which **effective probabilistic theories** can we reasonably expect to encounter?



$$\Omega = \{\rho \mid \rho \geq 0, \text{tr}(\rho) = 1\}$$

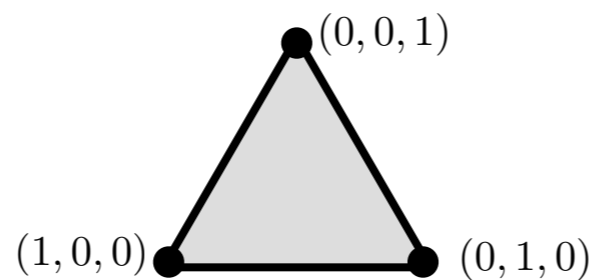
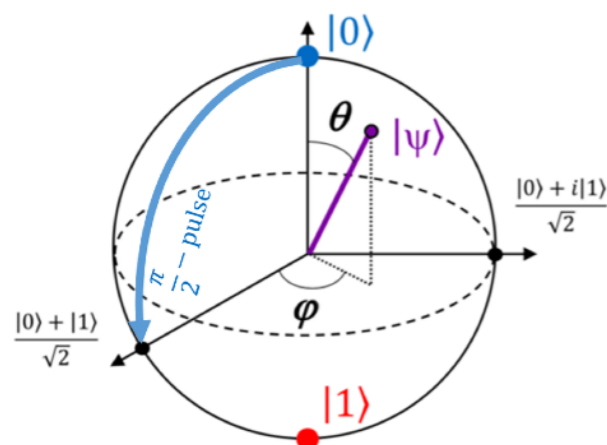
Two motivations

Suppose we **prepare** and **measure** a physical system in all ways accessible to us.



Could the resulting data **falsify QT** w/o assumptions on devices or physics?

If Nature is **fundamentally quantum**, which **effective probabilistic theories** can we reasonably expect to encounter?



$$\Omega = \{p = (p_1, \dots, p_n) \mid p_i \geq 0, \sum p_i = 1\}$$

- classical probability theory
- noisy qubits etc.
- QT w/ superselection rules
- ... ?

$$\Omega = \{\rho \mid \rho \geq 0, \text{tr}(\rho) = 1\}$$

Unambiguously testing / falsifying QT is really hard!

Science

 AAAS

Ruling Out Multi-Order Interference in Quantum Mechanics

Urbasi Sinha *et al.*

Science **329**, 418 (2010);

DOI: 10.1126/science.1190545

Unambiguously testing / falsifying QT is really hard!

Science

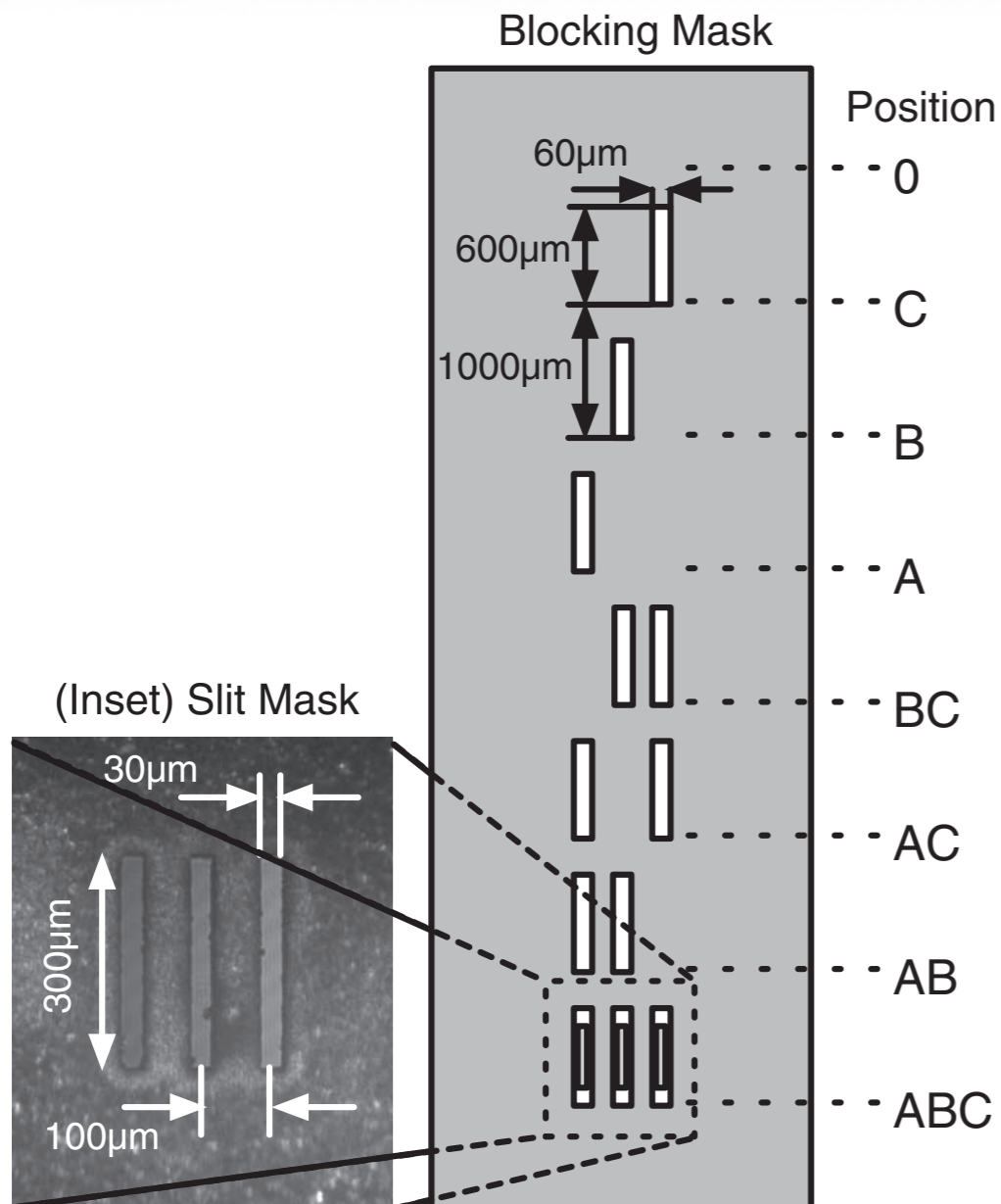
AAAS

Ruling Out Multi-Order Interference in Quantum Mechanics

Urbasi Sinha *et al.*

Science **329**, 418 (2010);

DOI: 10.1126/science.1190545



$$\begin{aligned}
 I_{ABC} &:= P_{ABC} - (P_A + P_B + P_C + I_{AB} + \\
 &\quad I_{BC} + I_{AC}) \\
 &= P_{ABC} - P_{AB} - P_{BC} - P_{AC} + P_A + \\
 &\quad P_B + P_C
 \end{aligned} \tag{5}$$

In QT, only **pairs of paths** interfere (Sorkin 1994)

$$\Rightarrow I_{ABC} = 0.$$

Unambiguously testing / falsifying QT is really hard!

Science

AAAS

Ruling Out Multi-Order Interference in Quantum Mechanics

Urbasi Sinha *et al.*

Science **329**, 418 (2010);

DOI: 10.1126/science.1190545

Non-classical paths in interference experiments

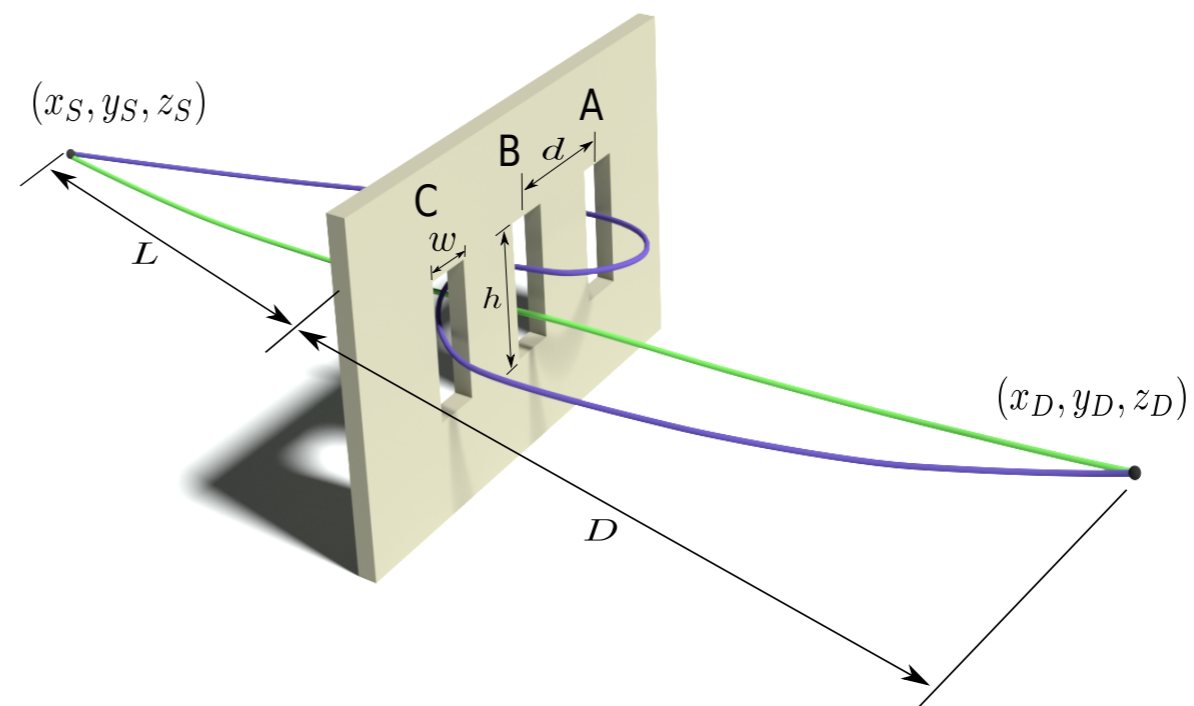
Rahul Sawant¹, Joseph Samuel¹, Aninda Sinha², Supurna Sinha¹ and Urbasi Sinha^{1,3}

¹*Raman Research Institute, Sadashivanagar, Bangalore, India.*

²*Centre for High Energy Physics, Indian Institute of Science, Bangalore, India.*

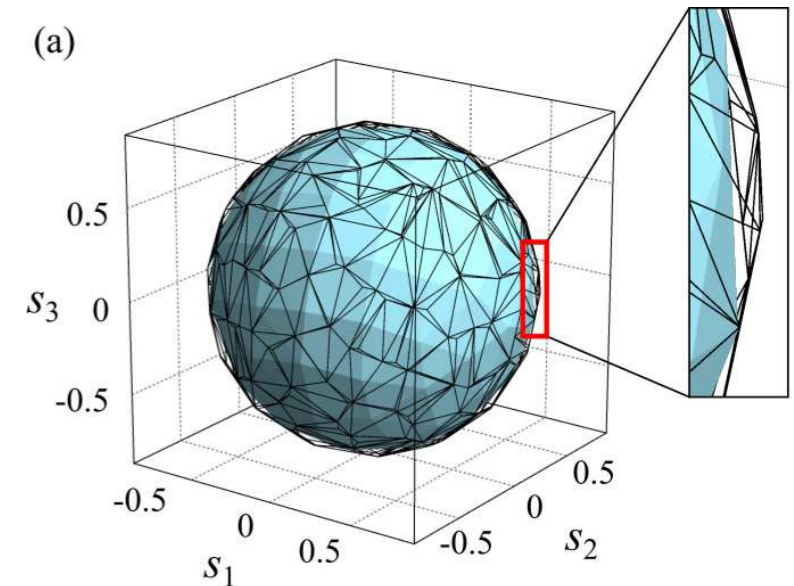
³*Institute for Quantum Computing, 200 University Avenue West, Waterloo, Ontario, Canada.*

* *To whom correspondence should be addressed; E-mail: usinha@rri.res.in.*



Overview

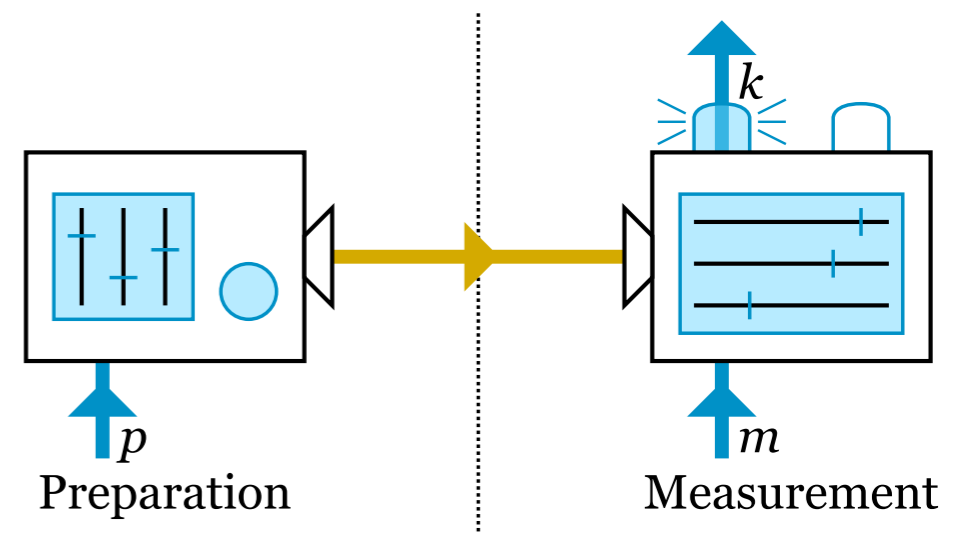
1. GPTs and theory-agnostic tomography



2. Contextuality, simulations, and embeddings

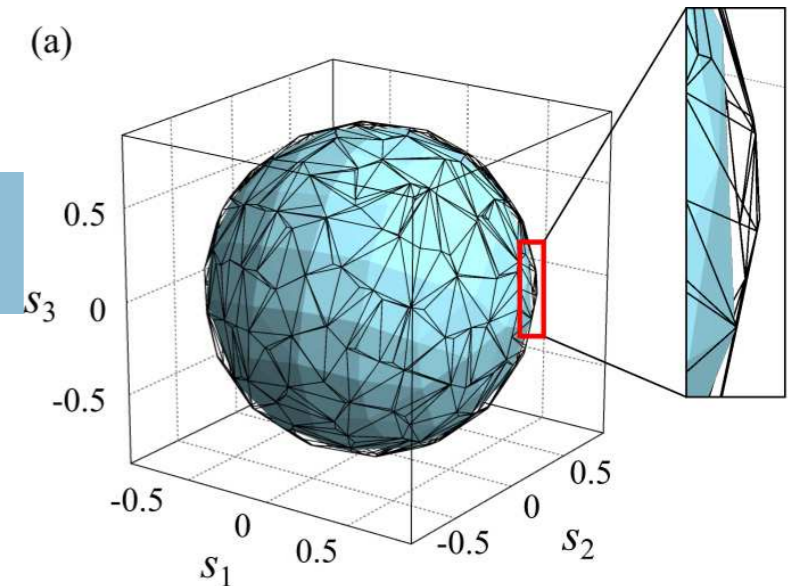
3. Exact embeddings into quantum theory

4. An experimental test of QT



Overview

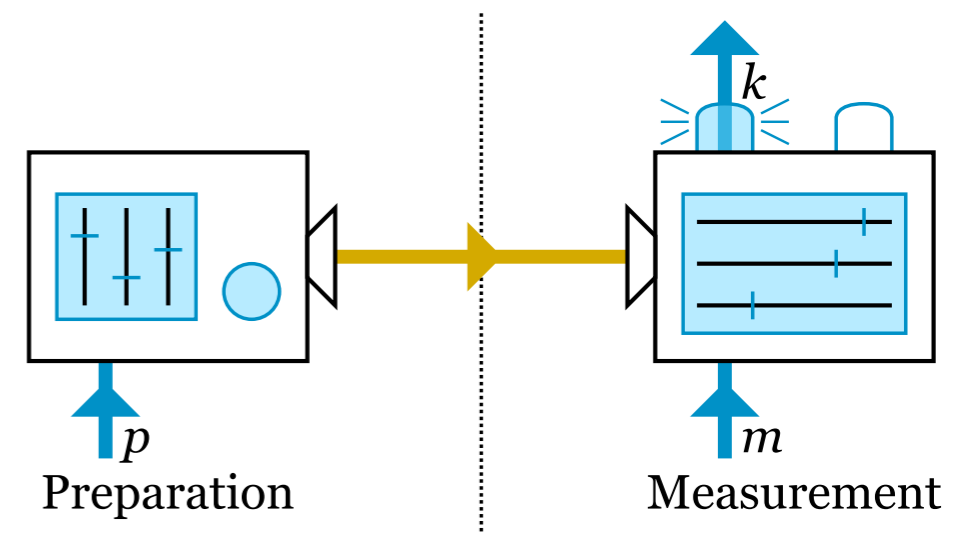
1. GPTs and theory-agnostic tomography



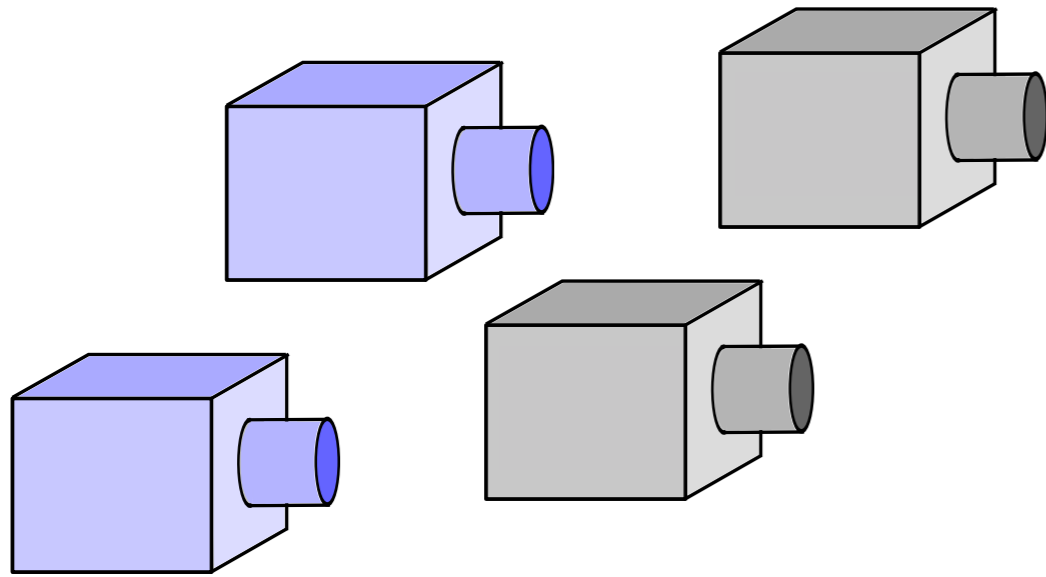
2. Contextuality, simulations, and embeddings

3. Exact embeddings into quantum theory

4. An experimental test of QT

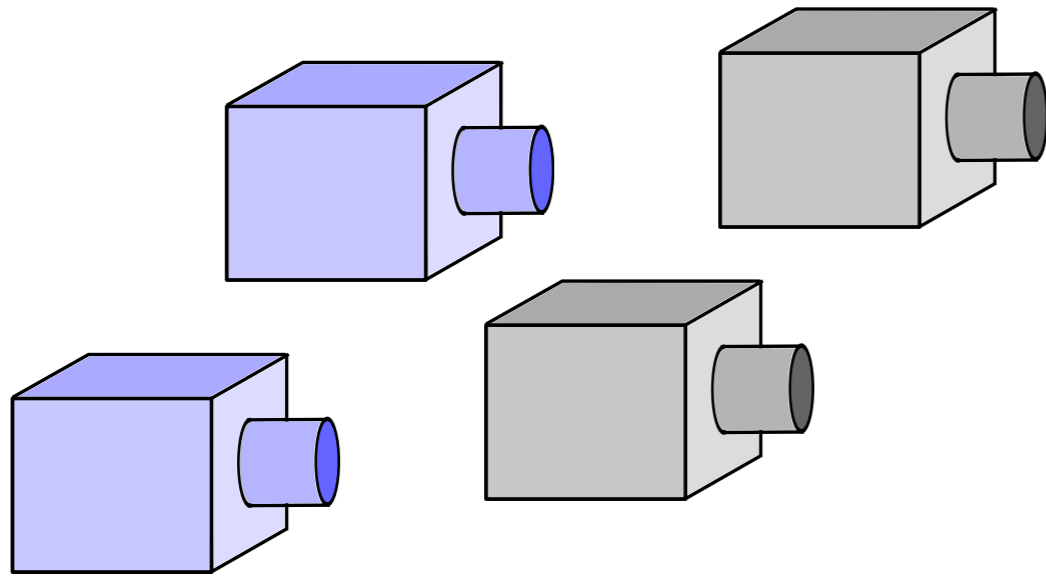


Operational theories



(all accessible preparation procedures)

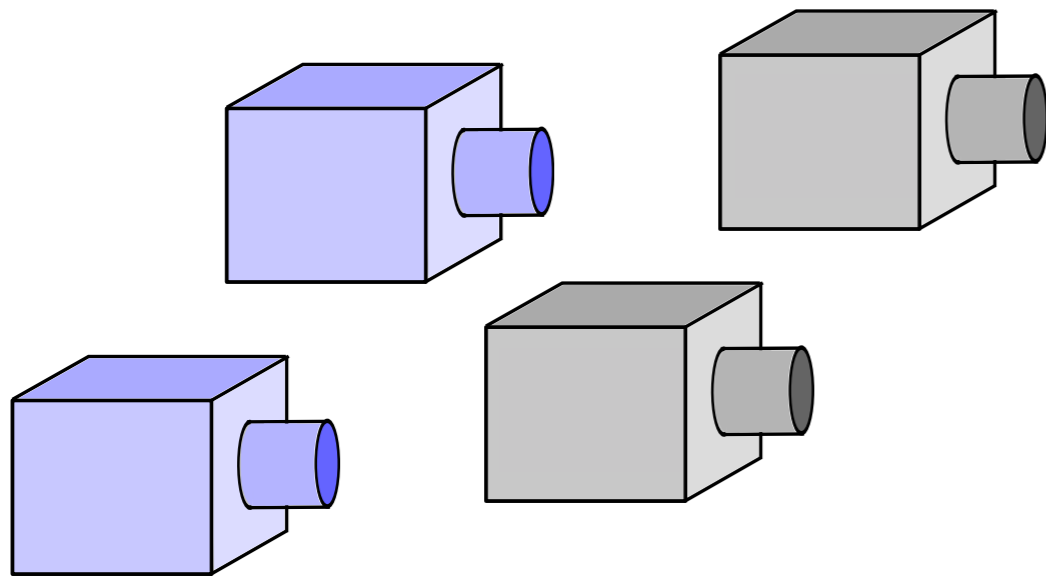
Operational theories



(all accessible preparation procedures)

$P_1 \sim P_2$ if they give identical probabilities for all outcomes of all accessible measurements.

Operational theories

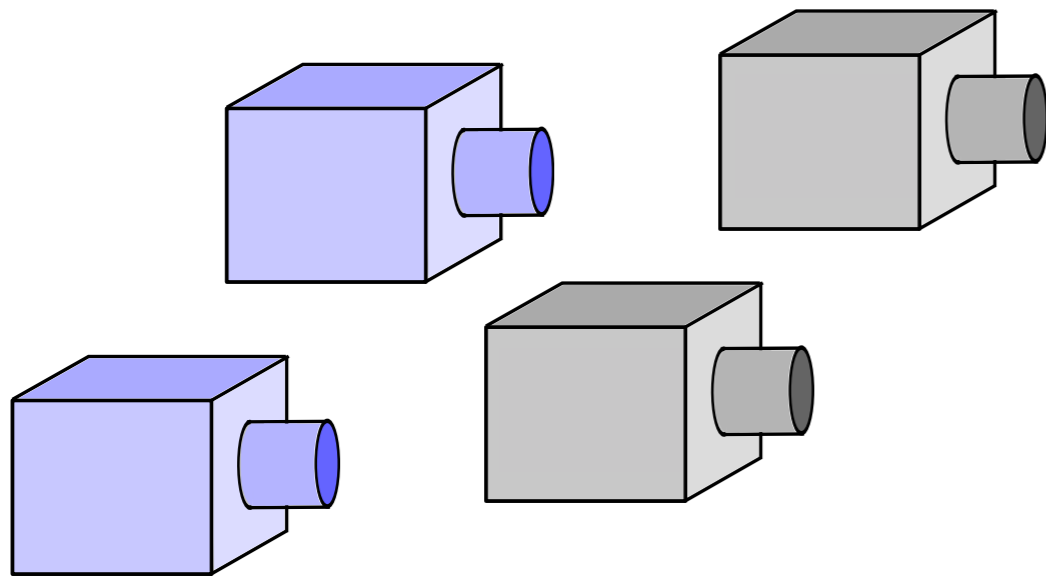


(all accessible preparation procedures)

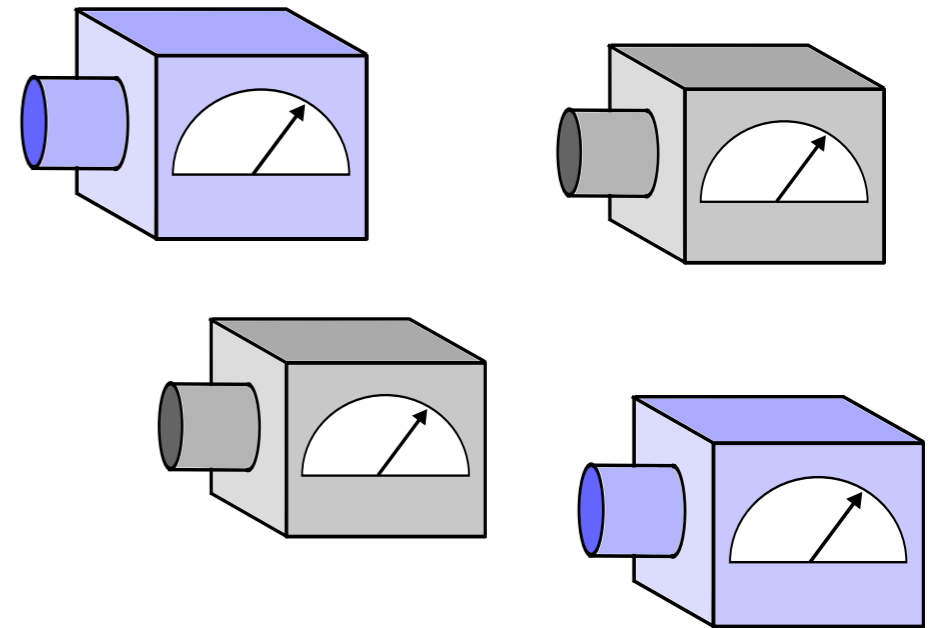
$P_1 \sim P_2$ if they give identical probabilities for all outcomes of all accessible measurements.

State ω_P = equivalence class of preparation procedures

Operational theories



(all accessible preparation procedures)

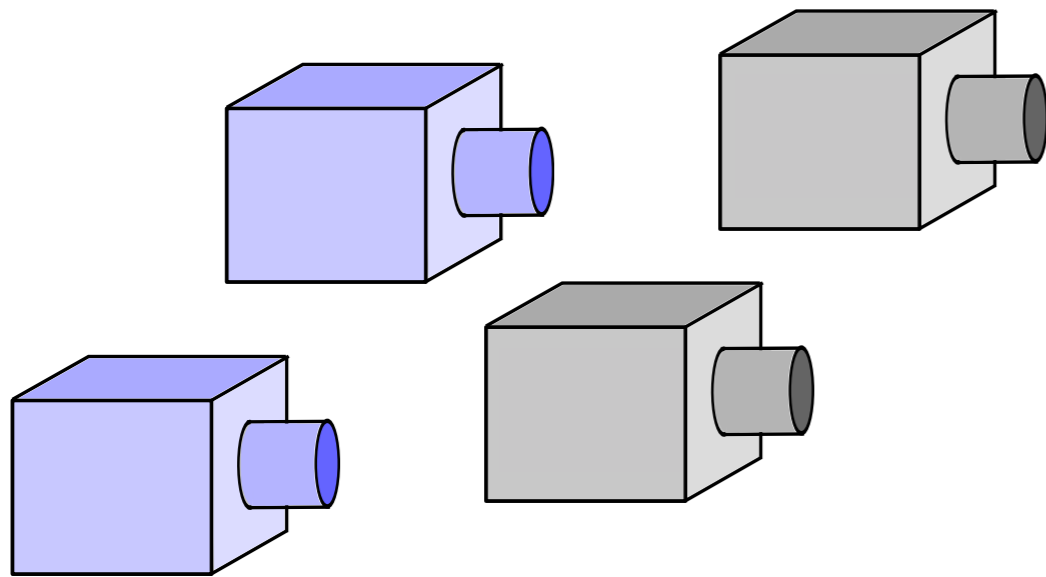


(all accessible measurement procedures)

$P_1 \sim P_2$ if they give identical probabilities for all outcomes of all accessible measurements.

State ω_P = equivalence class of preparation procedures

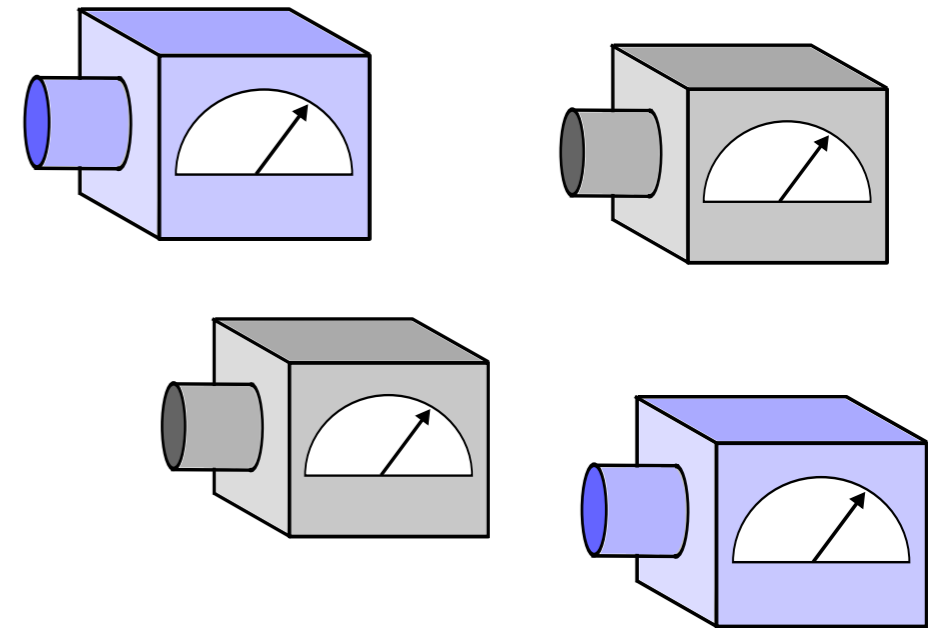
Operational theories



(all accessible preparation procedures)

$P_1 \sim P_2$ if they give identical probabilities for all outcomes of all accessible measurements.

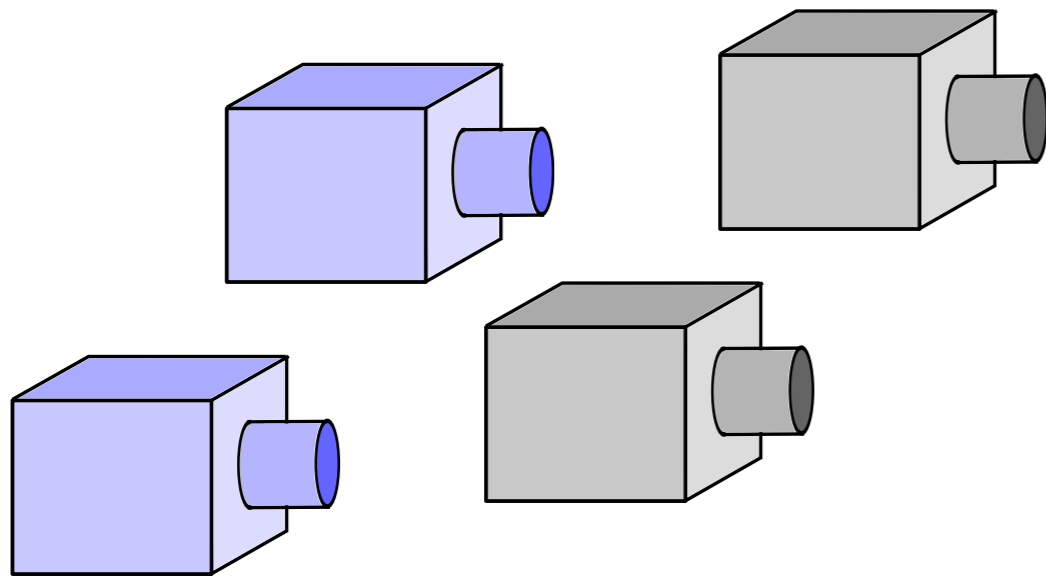
State ω_P = equivalence class of preparation procedures



(all accessible measurement procedures)

$(k_1, M_1) \sim (k_2, M_2)$ if $\text{Prob}(k_1|M_1, P) = \text{Prob}(k_2|M_2, P)$ for all accessible preparations P .

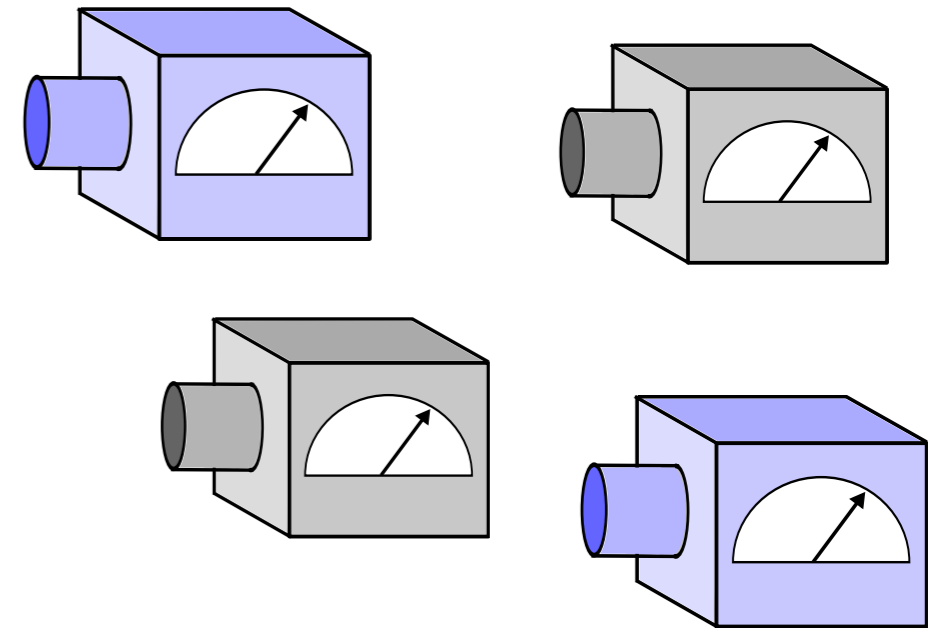
Operational theories



(all accessible preparation procedures)

$P_1 \sim P_2$ if they give identical probabilities for all outcomes of all accessible measurements.

State ω_P = equivalence class of preparation procedures

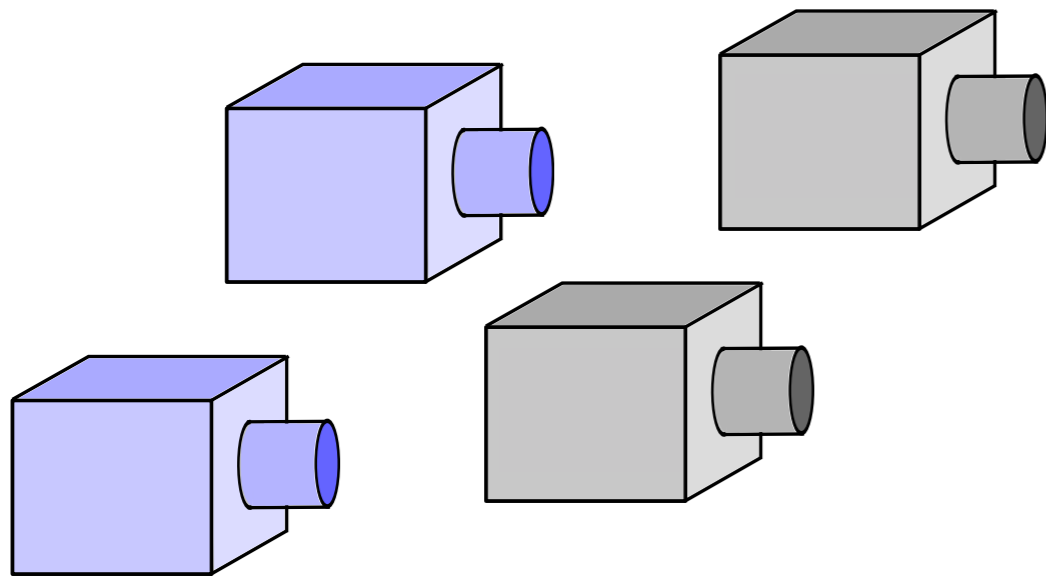


(all accessible measurement procedures)

$(k_1, M_1) \sim (k_2, M_2)$ if $\text{Prob}(k_1|M_1, P) = \text{Prob}(k_2|M_2, P)$ for all accessible preparations P .

Effect $e_{k,M}$ = equivalence class of outcome-measurement pairs

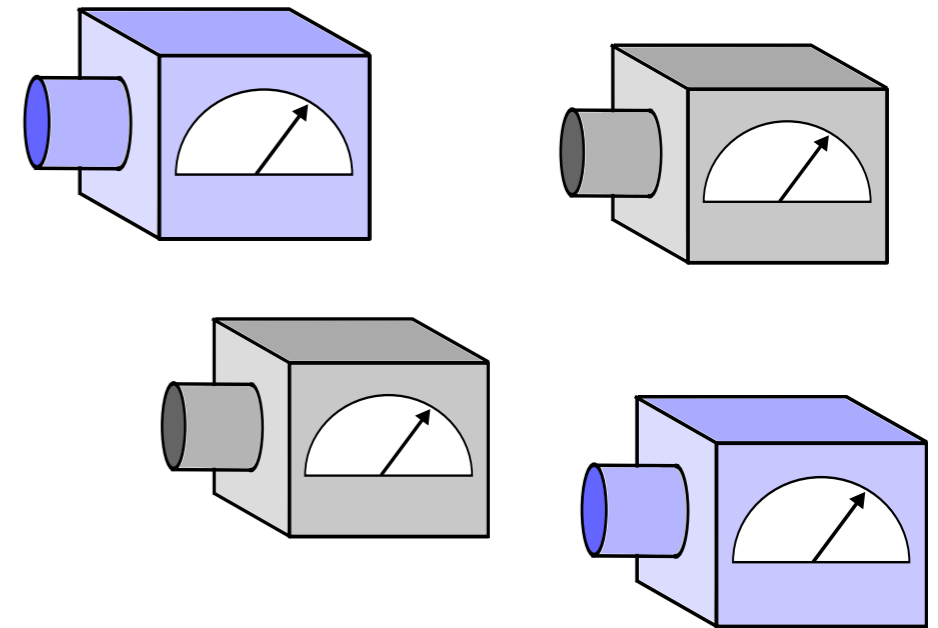
Operational theories



(all accessible preparation procedures)

$P_1 \sim P_2$ if they give identical probabilities for all outcomes of all accessible measurements.

State ω_P = equivalence class of preparation procedures



(all accessible measurement procedures)

$(k_1, M_1) \sim (k_2, M_2)$ if $\text{Prob}(k_1|M_1, P) = \text{Prob}(k_2|M_2, P)$ for all accessible preparations P .

Effect $e_{k,M}$ = equivalence class of outcome-measurement pairs

$$\text{Prob}(k|P, M) = \langle \omega_P, e_{k,M} \rangle \quad (e_{k,M} \in A, \omega_P \in A^*).$$

General probabilistic theories

$$\text{Prob}(k|P, M) = \langle \omega_P, e_{k,M} \rangle \quad (e_{k,M} \in A, \omega_P \in A^*).$$

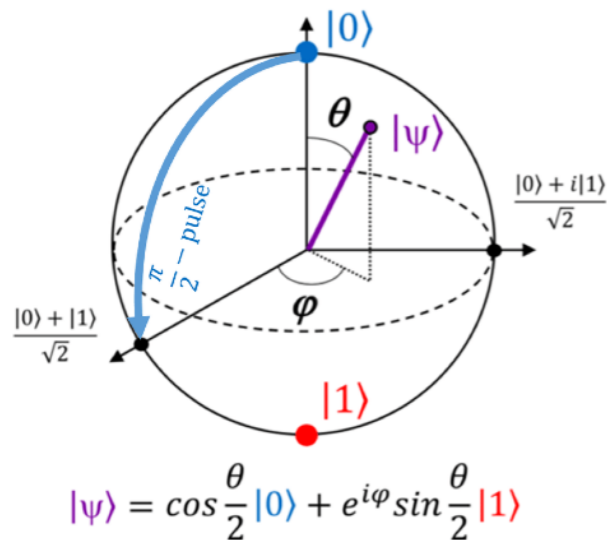
General probabilistic theories

GPT $\mathcal{A} = (A, \Omega_A, E_A)$ = (vector space over \mathbb{R} , normalized states, effects).

$$\text{Prob}(k|P, M) = \langle \omega_P, e_{k,M} \rangle \quad (e_{k,M} \in A, \omega_P \in A^*).$$

General probabilistic theories

GPT $\mathcal{A} = (A, \Omega_A, E_A) = (\text{vector space over } \mathbb{R}, \text{normalized states, effects}).$



Quantum theory (QT): \mathcal{Q}_n

$A = \mathbb{H}_n(\mathbb{C})$ (complex Hermitian $n \times n$ matrices)

$E_A = \{E \mid 0 \leq E \leq \mathbf{1}\}$ (POVM elements)

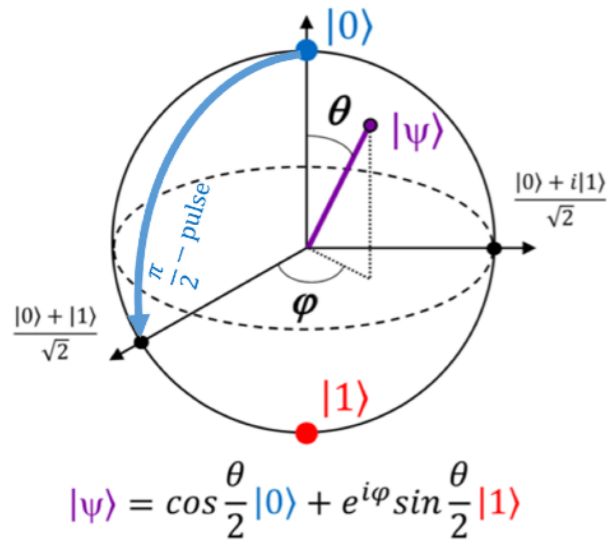
$\Omega_A = \{\rho \mid \rho \geq 0, \text{tr}(\rho) = 1\}$ (density matrices)

$A^* \simeq A$ via $\langle X, Y \rangle = \text{tr}(XY)$.

$$\text{Prob}(k|P, M) = \langle \omega_P, e_{k,M} \rangle \quad (e_{k,M} \in A, \omega_P \in A^*).$$

General probabilistic theories

GPT $\mathcal{A} = (A, \Omega_A, E_A) = (\text{vector space over } \mathbb{R}, \text{normalized states, effects}).$



Quantum theory (QT): \mathcal{Q}_n

$A = \mathbb{H}_n(\mathbb{C})$ (complex Hermitian $n \times n$ matrices)

$E_A = \{E \mid 0 \leq E \leq \mathbf{1}\}$ (POVM elements)

$\Omega_A = \{\rho \mid \rho \geq 0, \text{tr}(\rho) = 1\}$ (density matrices)

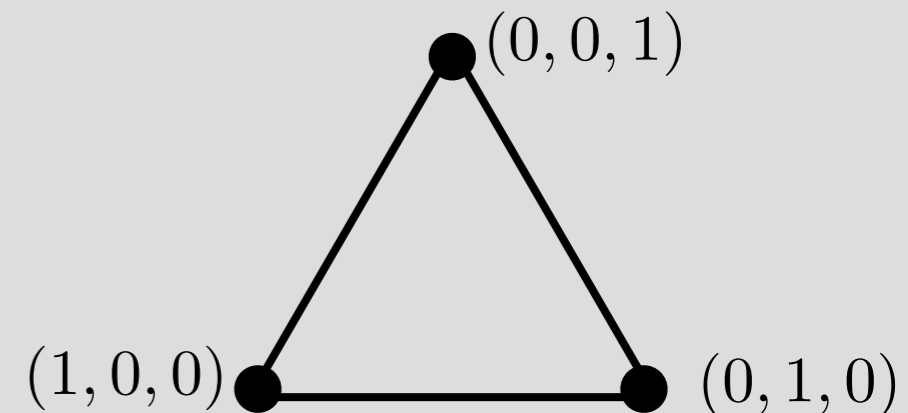
$A^* \simeq A$ via $\langle X, Y \rangle = \text{tr}(XY)$.

Classical probability theory (QT): \mathcal{C}_n

$A = \mathbb{R}^n \simeq A^*$

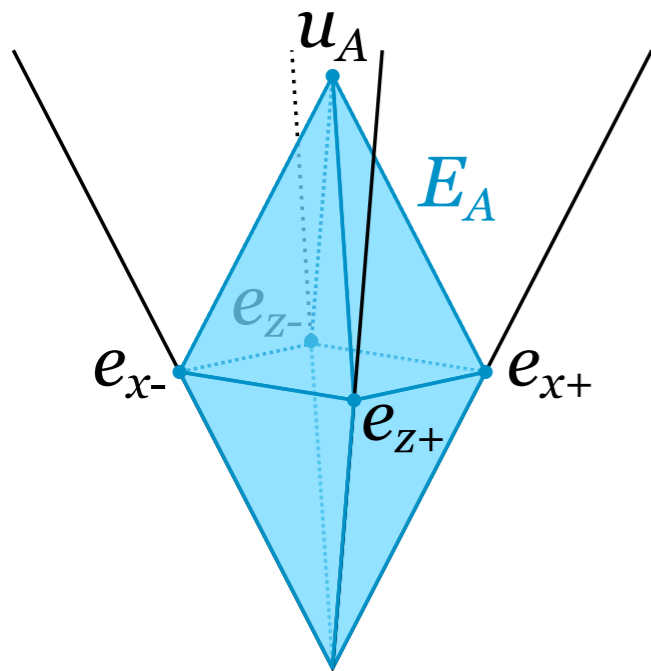
$E_A = \{(e_1, \dots, e_n) \mid 0 \leq e_i \leq 1\}$

$\Omega_A = \left\{ (p_1, \dots, p_n) \mid p_i \geq 0, \sum_i p_i = 1 \right\}.$

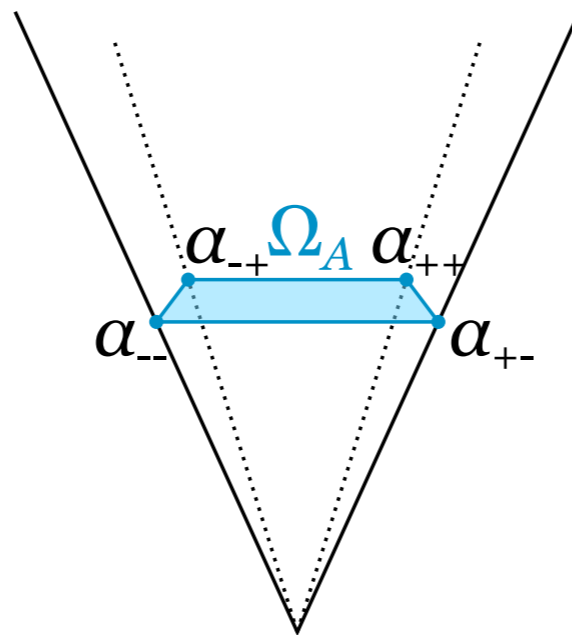


General probabilistic theories

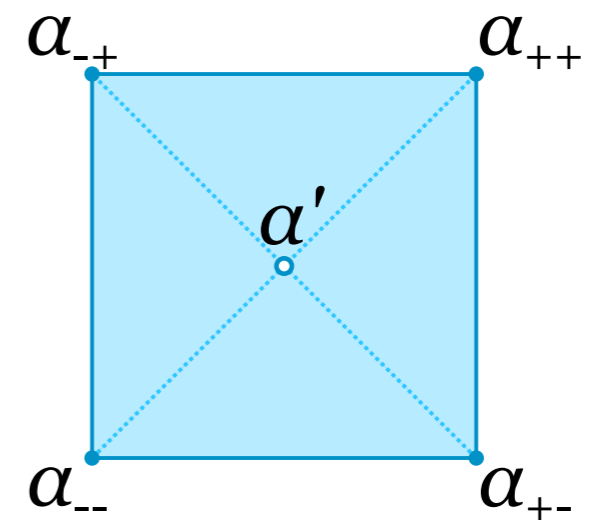
The gbit $\mathcal{A} = (\mathbb{R}^3, \Omega_A, E_A)$



a) Cone of effects A_+



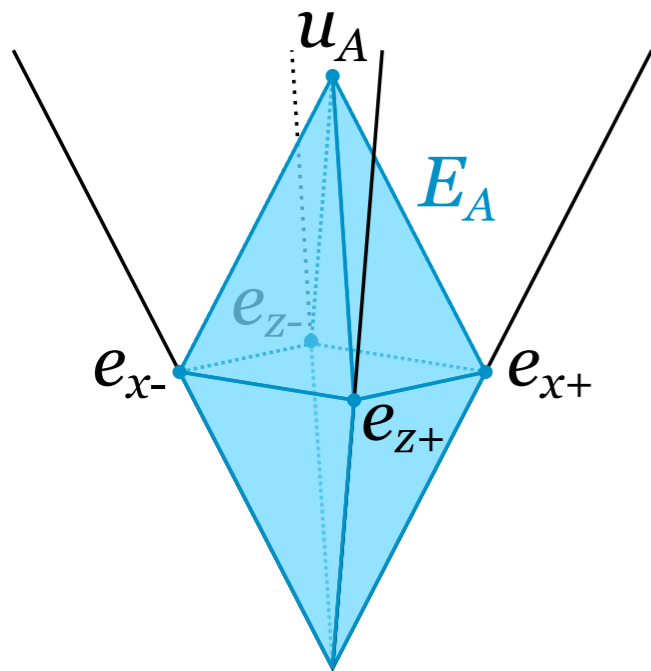
b) Cone of states A_+^*



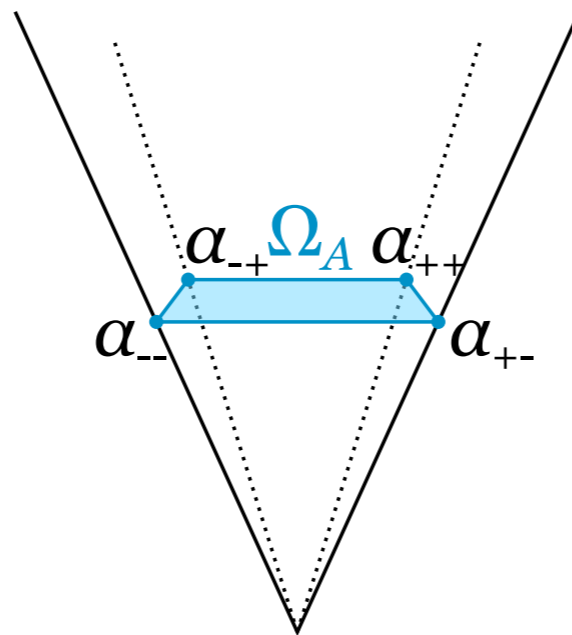
c) Normalized states Ω_A

General probabilistic theories

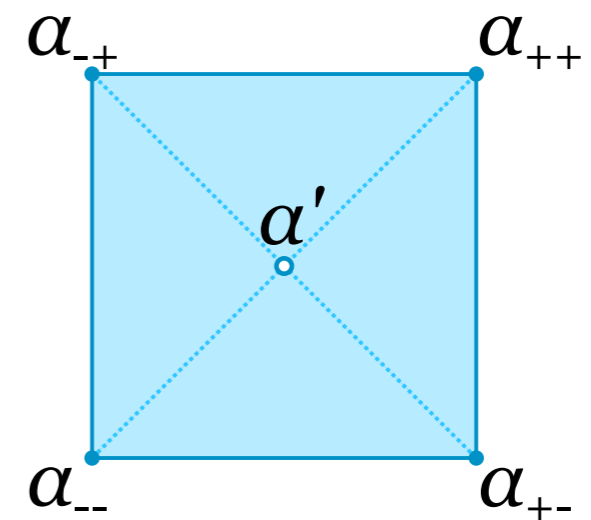
The gbit $\mathcal{A} = (\mathbb{R}^3, \Omega_A, E_A)$



a) Cone of effects A_+



b) Cone of states A_+^*



c) Normalized states Ω_A

The four pure states $\alpha_{\pm, \pm}$ are **pairwise** perfectly distinguishable, but **not jointly** \implies this cannot be a classical or quantum system.

Theory-agnostic tomography

Idea: Identify a **physical system**. Perform as many preparations and measurements as possible; **fit a GPT to the data**; compare with \mathcal{Q}_n .

Theory-agnostic tomography

Idea: Identify a **physical system**. Perform as many preparations and measurements as possible; **fit a GPT to the data**; compare with \mathcal{Q}_n .

[1] M. D. Mazurek, M. F. Pusey, K. J. Resch, and R. W. Spekkens, PRX Quantum **2**, 020302 (2021).

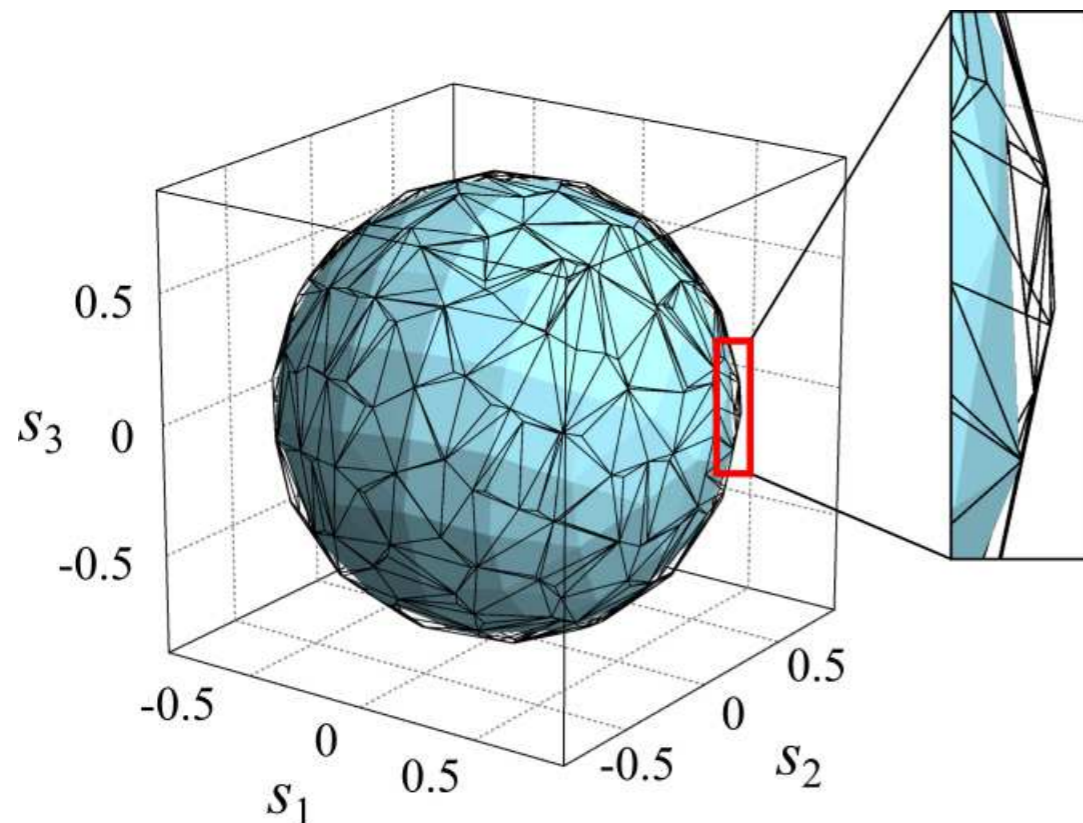
[2] M. Grabowecky, C. Pollack, A. Cameron, R. W. Spekkens, and K. J. Resch, Phys. Rev. A **105**, 032204 (2022).

Theory-agnostic tomography

Idea: Identify a **physical system**. Perform as many preparations and measurements as possible; **fit a GPT to the data**; compare with \mathcal{Q}_n .

[1] M. D. Mazurek, M. F. Pusey, K. J. Resch, and R. W. Spekkens, PRX Quantum **2**, 020302 (2021).

[2] M. Grabowecky, C. Pollack, A. Cameron, R. W. Spekkens, and K. J. Resch, Phys. Rev. A **105**, 032204 (2022).



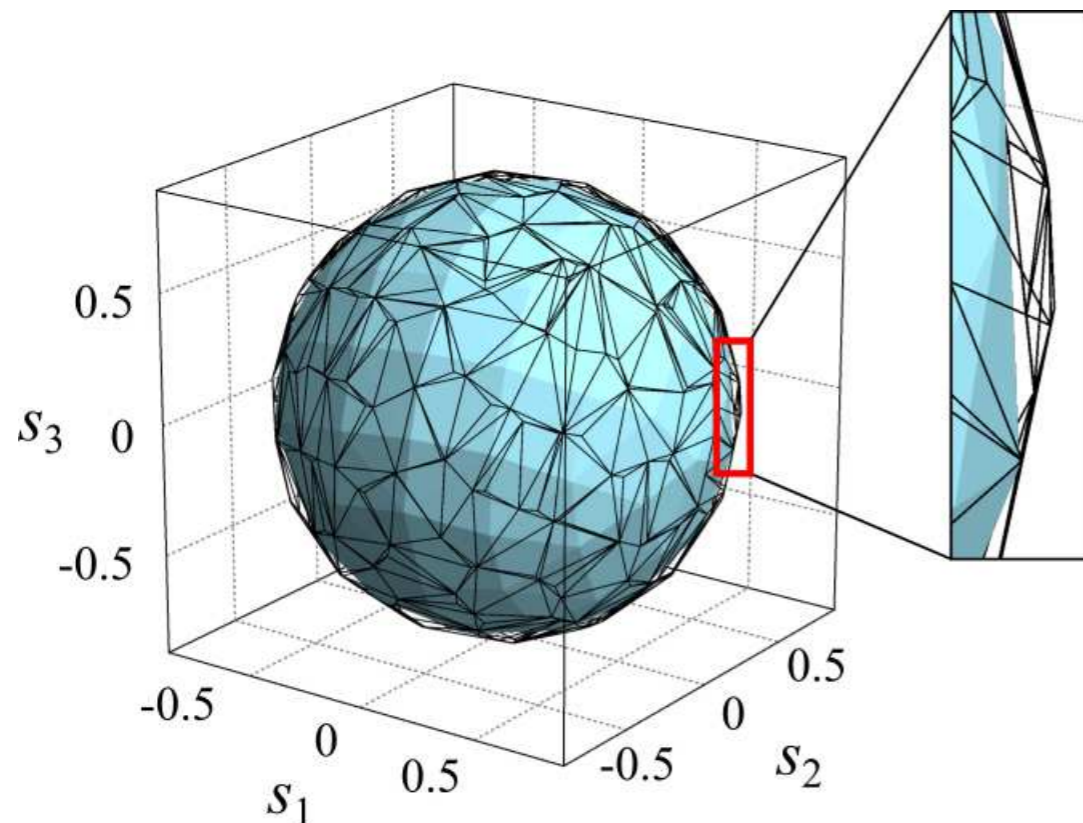
[1]: Polarization degree of freedom of a single photon: “bumpy qubit” $\approx \mathcal{Q}_2$.

Theory-agnostic tomography

Idea: Identify a **physical system**. Perform as many preparations and measurements as possible; **fit a GPT to the data**; compare with \mathcal{Q}_n .

[1] M. D. Mazurek, M. F. Pusey, K. J. Resch, and R. W. Spekkens, PRX Quantum **2**, 020302 (2021).

[2] M. Grabowecky, C. Pollack, A. Cameron, R. W. Spekkens, and K. J. Resch, Phys. Rev. A **105**, 032204 (2022).



[1]: Polarization degree of freedom of a single photon: “bumpy qubit” $\approx \mathcal{Q}_2$.

Tomographic completeness loophole:
can never be sure that we probed the system *completely*.

What if we just see a (low-dimensional) “**shadow**”?

Let’s **drop** the **tomographic completeness** assumption.

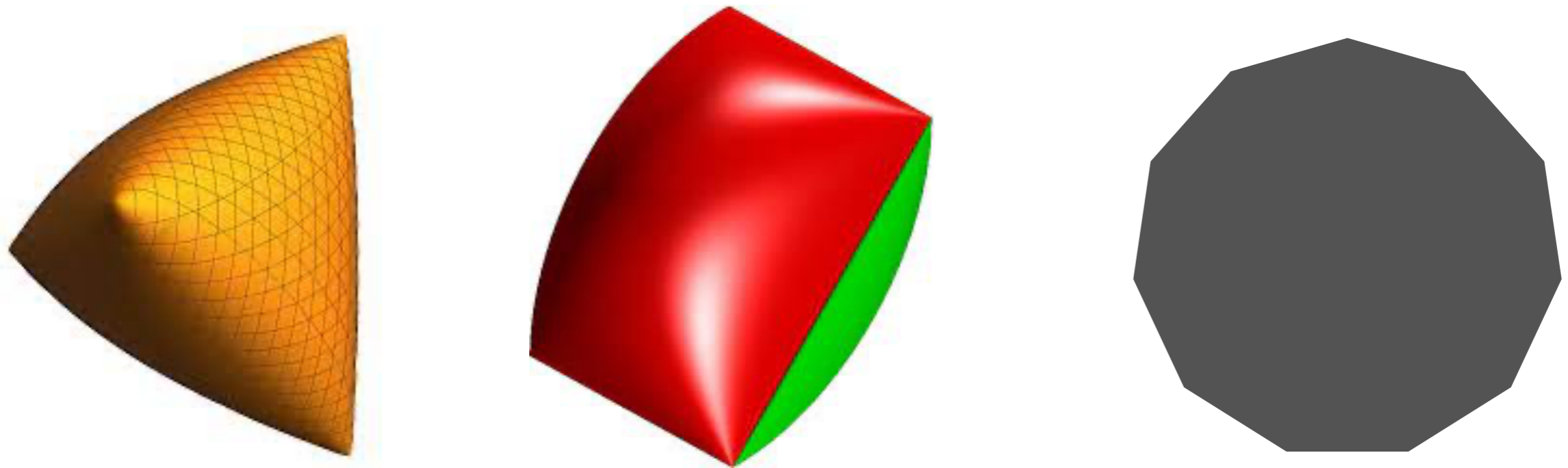
“Effective physical system”: **defined** by a set of accessible procedures.

What if we just see a (low-dimensional) “shadow”?

Let’s **drop** the **tomographic completeness** assumption.

“Effective physical system”: **defined** by a set of accessible procedures.

If we do theory-agnostic tomography on an effective physical system and obtain some weird noisy GPT, is **QT** a **possible/plausible explanation**?

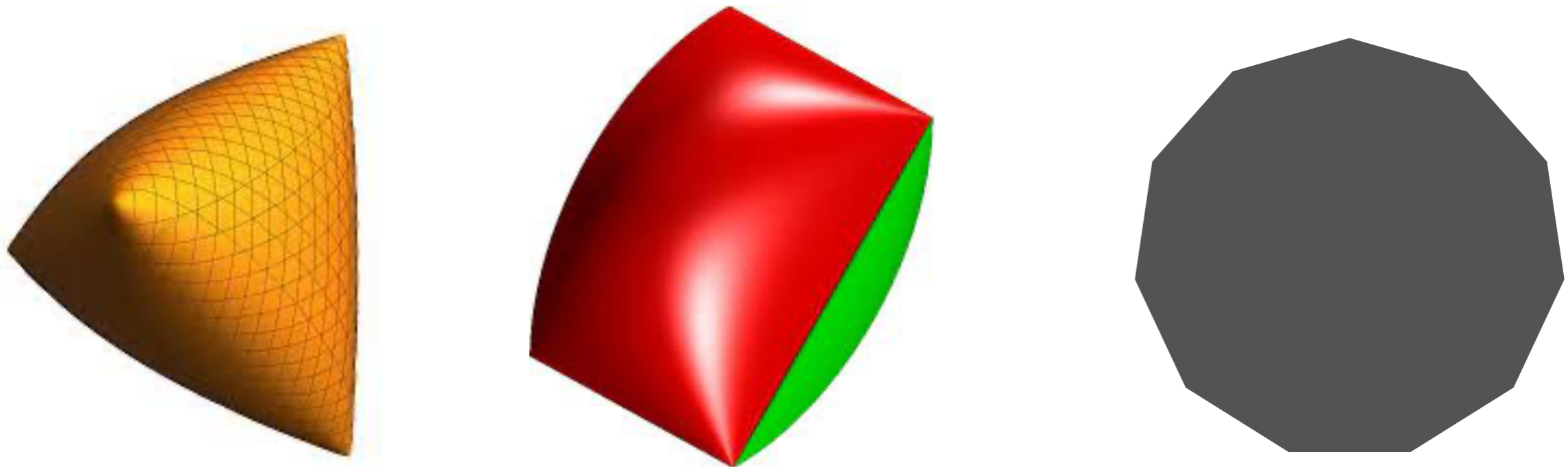


What if we just see a (low-dimensional) “shadow”?

Let’s **drop** the **tomographic completeness** assumption.

“Effective physical system”: **defined** by a set of accessible procedures.

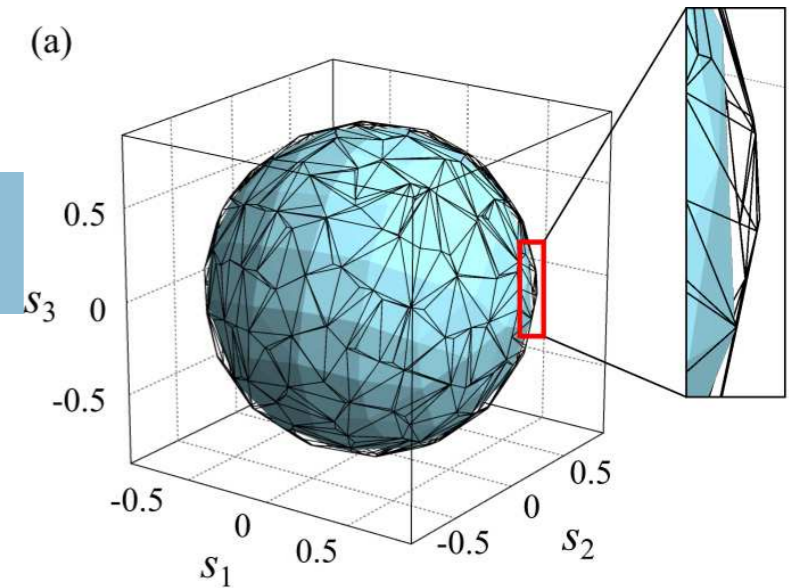
If we do theory-agnostic tomography on an effective physical system and obtain some weird noisy GPT, is **QT** a **possible/plausible explanation**?



Is **fundamental QT** a plausible explanation of a given **effective GPT**?

Overview

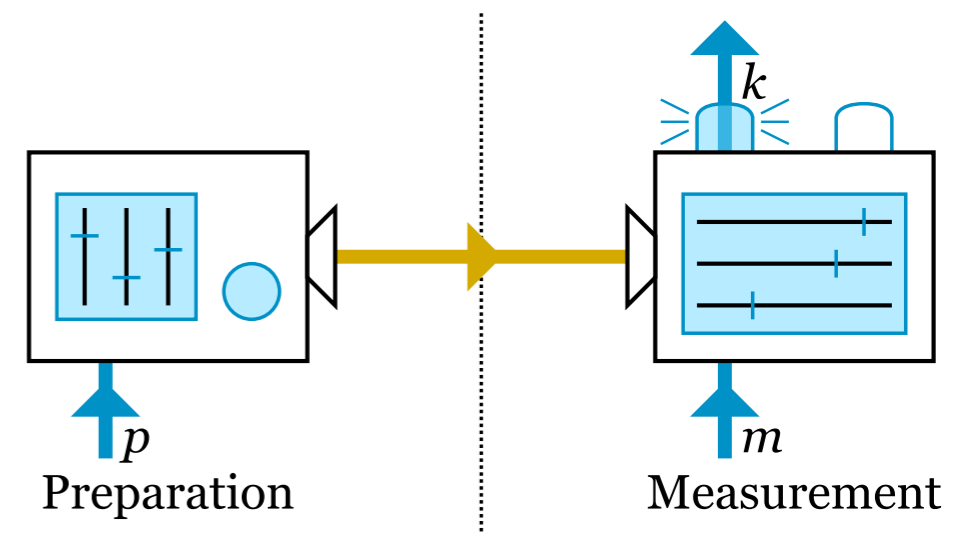
1. GPTs and theory-agnostic tomography



2. Contextuality, simulations, and embeddings

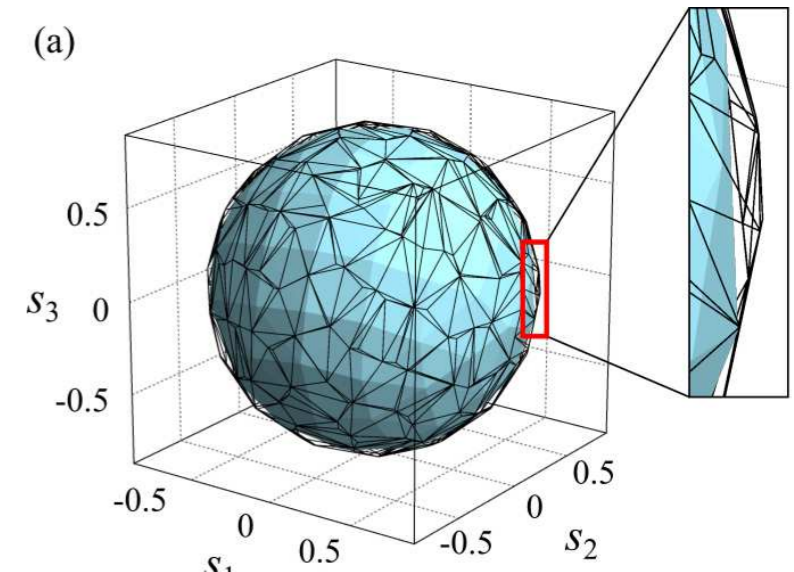
3. Exact embeddings into quantum theory

4. An experimental test of QT



Overview

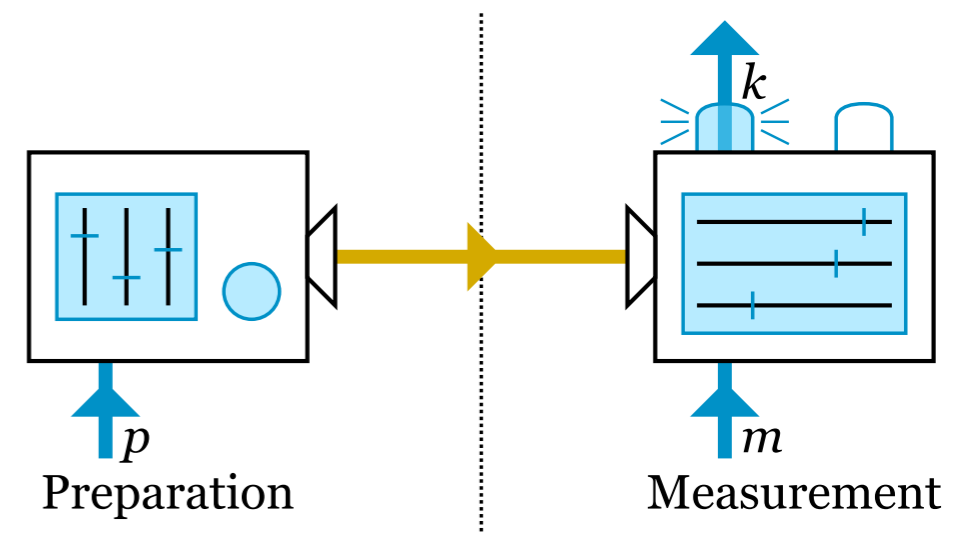
1. GPTs and theory-agnostic tomography



2. Contextuality, simulations, and embeddings

3. Exact embeddings into quantum theory

4. An experimental test of QT



Spekkens' notion of noncontextuality: quick recap

Contextuality for preparations, transformations and unsharp measurements

R. W. Spekkens*

Perimeter Institute for Theoretical Physics, 35 King St. North, Waterloo, Ontario N2J 2W9, Canada

(Dated: Feb. 25, 2005)



Spekkens' notion of noncontextuality: quick recap

Contextuality for preparations, transformations and unsharp measurements

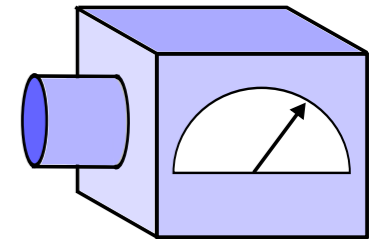
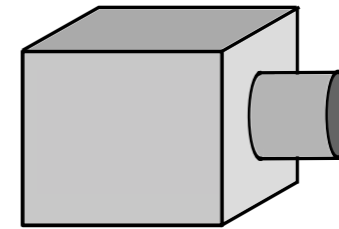
R. W. Spekkens*

Perimeter Institute for Theoretical Physics, 35 King St. North, Waterloo, Ontario N2J 2W9, Canada

(Dated: Feb. 25, 2005)



Recall the notion of an **operational theory**.



Spekkens' notion of noncontextuality: quick recap

Contextuality for preparations, transformations and unsharp measurements

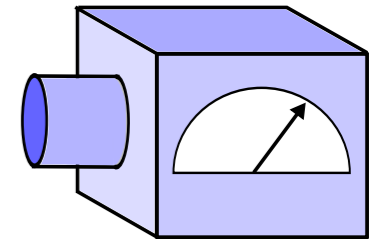
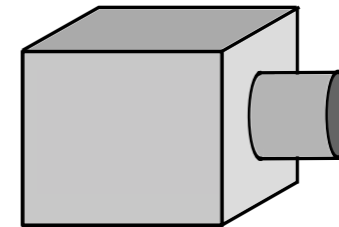
R. W. Spekkens*

Perimeter Institute for Theoretical Physics, 35 King St. North, Waterloo, Ontario N2J 2W9, Canada

(Dated: Feb. 25, 2005)



Recall the notion of an **operational theory**.



Ontological model of a system (e.g. of a qubit):

A set of classical variables Λ .

Spekkens' notion of noncontextuality: quick recap

Contextuality for preparations, transformations and unsharp measurements

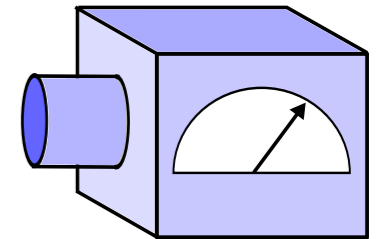
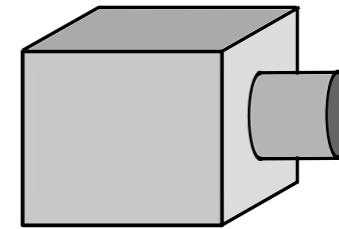
R. W. Spekkens*

Perimeter Institute for Theoretical Physics, 35 King St. North, Waterloo, Ontario N2J 2W9, Canada

(Dated: Feb. 25, 2005)



Recall the notion of an **operational theory**.



Ontological model of a system (e.g. of a qubit):

A set of classical variables Λ .

Preparation procedure $P \longleftrightarrow$ distribution $\mu_P(\lambda)$

Spekkens' notion of noncontextuality: quick recap

Contextuality for preparations, transformations and unsharp measurements

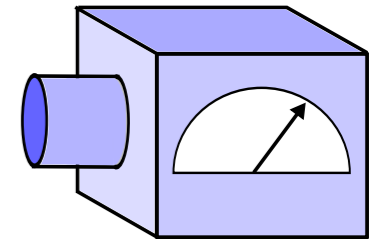
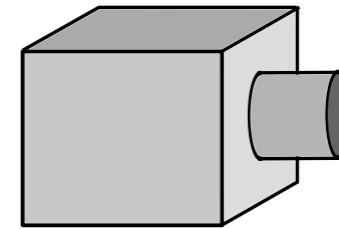
R. W. Spekkens*

Perimeter Institute for Theoretical Physics, 35 King St. North, Waterloo, Ontario N2J 2W9, Canada

(Dated: Feb. 25, 2005)



Recall the notion of an **operational theory**.



Ontological model of a system (e.g. of a qubit):

A set of classical variables Λ .

Preparation procedure $P \longleftrightarrow$ distribution $\mu_P(\lambda)$

Outcome k of measurement $M \longleftrightarrow$ response function $\chi_{M,k}(\lambda)$

Spekkens' notion of noncontextuality: quick recap

Contextuality for preparations, transformations and unsharp measurements

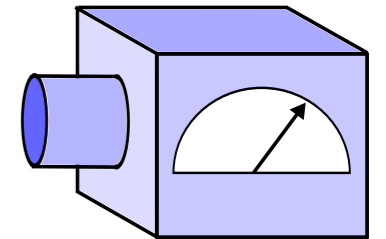
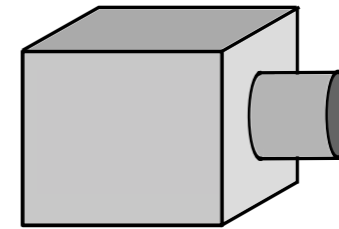
R. W. Spekkens*

Perimeter Institute for Theoretical Physics, 35 King St. North, Waterloo, Ontario N2J 2W9, Canada

(Dated: Feb. 25, 2005)



Recall the notion of an **operational theory**.



Ontological model of a system (e.g. of a qubit):

A set of classical variables Λ .

Preparation procedure $P \longleftrightarrow$ distribution $\mu_P(\lambda)$

Outcome k of measurement $M \longleftrightarrow$ response function $\chi_{M,k}(\lambda)$

such that

$$\text{Prob}(k|P, M) = \int_{\Lambda} d\lambda \mu_P(\lambda) \chi_{M,k}(\lambda).$$

Spekkens' notion of noncontextuality: quick recap

$$\text{Prob}(k|P, M) = \int_{\Lambda} d\lambda \mu_P(\lambda) \chi_{M,k}(\lambda)$$

Spekkens' notion of noncontextuality: quick recap

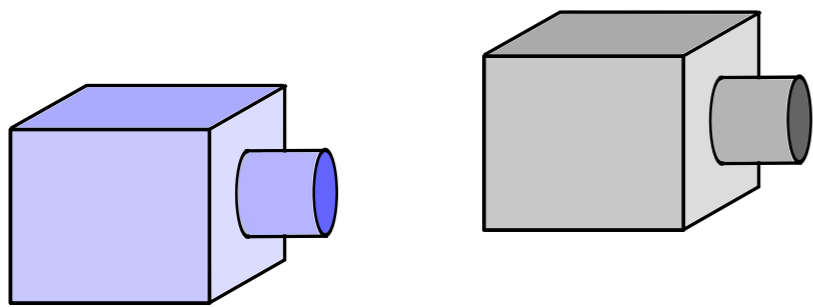
$$\text{Prob}(k|P, M) = \int_{\Lambda} d\lambda \mu_P(\lambda) \chi_{M,k}(\lambda)$$

The ontological model is **preparation-noncontextual** if $P \sim P' \Rightarrow \mu_P = \mu_{P'}$.

Spekkens' notion of noncontextuality: quick recap

$$\text{Prob}(k|P, M) = \int_{\Lambda} d\lambda \mu_P(\lambda) \chi_{M,k}(\lambda)$$

The ontological model is **preparation-noncontextual** if $P \sim P' \Rightarrow \mu_P = \mu_{P'}$.

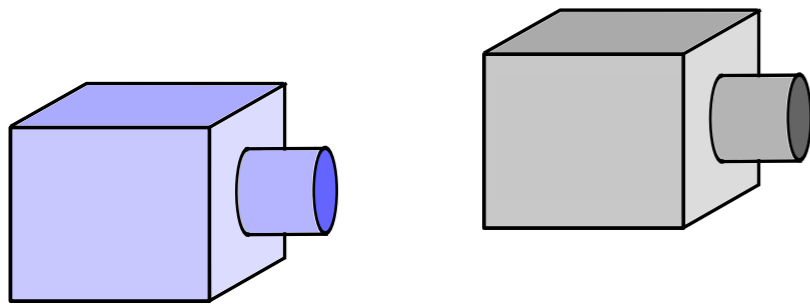


Intuition: preparation procedures are statistically indistinguishable **because** they prepare the same distribution over Λ .

Spekkens' notion of noncontextuality: quick recap

$$\text{Prob}(k|P, M) = \int_{\Lambda} d\lambda \mu_P(\lambda) \chi_{M,k}(\lambda)$$

The ontological model is **preparation-noncontextual** if $P \sim P' \Rightarrow \mu_P = \mu_{P'}$.



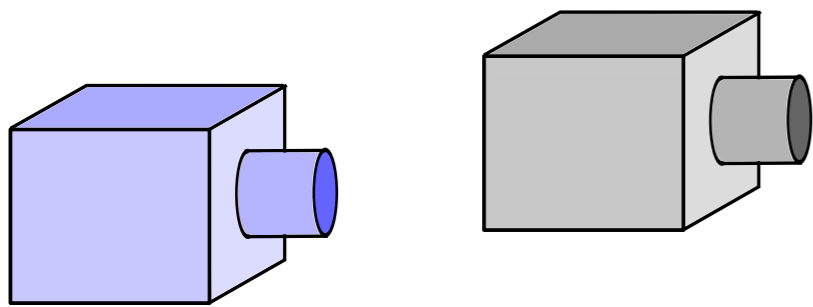
Intuition: preparation procedures are statistically indistinguishable **because** they prepare the same distribution over Λ .

Measurement-noncontextuality: $(k, M) \sim (k', M') \Rightarrow \chi_{k,M}(\lambda) = \chi_{k',M'}(\lambda)$

Spekkens' notion of noncontextuality: quick recap

$$\text{Prob}(k|P, M) = \int_{\Lambda} d\lambda \mu_P(\lambda) \chi_{M,k}(\lambda)$$

The ontological model is **preparation-noncontextual** if $P \sim P' \Rightarrow \mu_P = \mu_{P'}$.



Intuition: preparation procedures are statistically indistinguishable **because** they prepare the same distribution over Λ .

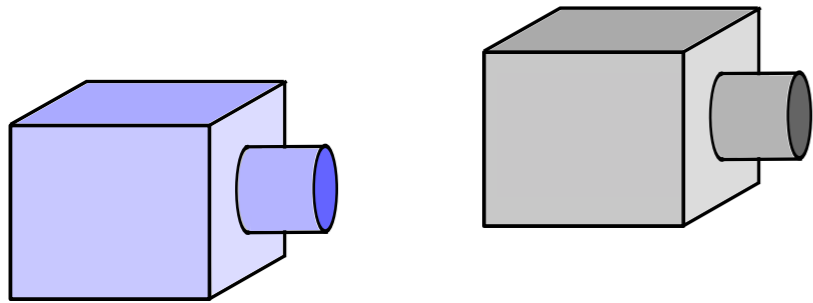
Measurement-noncontextuality: $(k, M) \sim (k', M') \Rightarrow \chi_{k,M}(\lambda) = \chi_{k',M'}(\lambda)$

Theorem: Ontological models of QM-systems must be preparation-contextual (and, assuming outcome-determinism for sharp meas., measurement-contextual).

Spekkens' notion of noncontextuality: quick recap

$$\text{Prob}(k|P, M) = \int_{\Lambda} d\lambda \mu_P(\lambda) \chi_{M,k}(\lambda)$$

The ontological model is **preparation-noncontextual** if $P \sim P' \Rightarrow \mu_P = \mu_{P'}$.



Intuition: preparation procedures are statistically indistinguishable **because** they prepare the same distribution over Λ .

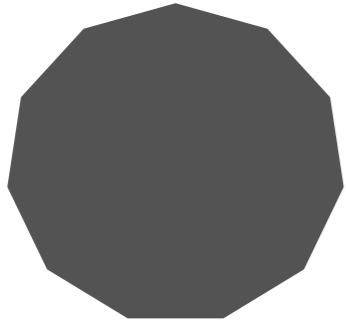
Measurement-noncontextuality: $(k, M) \sim (k', M') \Rightarrow \chi_{k,M}(\lambda) = \chi_{k',M'}(\lambda)$

Theorem: Ontological models of QM-systems must be preparation-contextual (and, assuming outcome-determinism for sharp meas., measurement-contextual).

Intuition: Contextual models are implausible because they are **fine-tuned:** operationally, $P \sim P'$, but ontologically, $\mu_P \neq \mu_{P'}$.

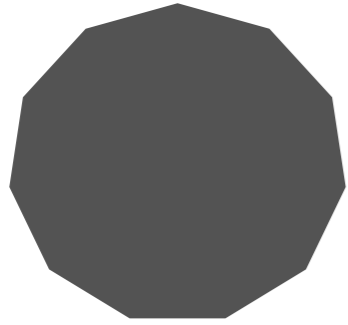
An instance of Leibniz' principle of the "identity of the indiscernibles".

Simulations and embeddings



Effective GPT $\mathcal{A} = (A, \Omega_A, E_A)$ found in the lab

Simulations and embeddings

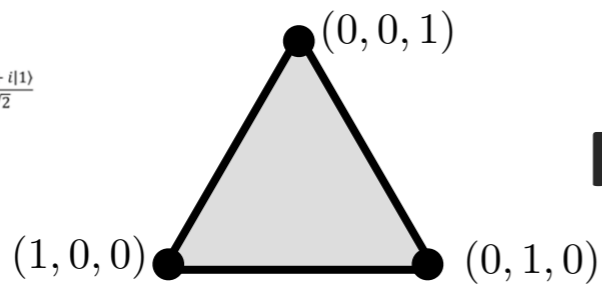
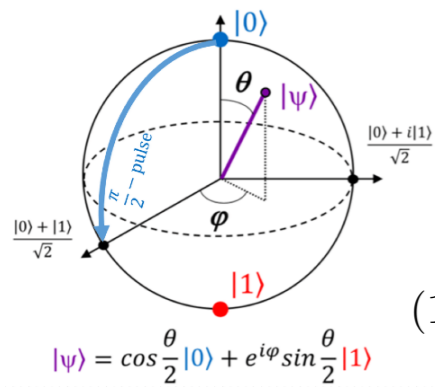


Effective GPT $\mathcal{A} = (A, \Omega_A, E_A)$ found in the lab

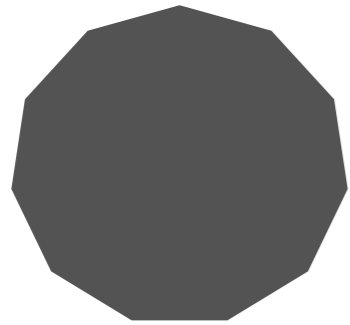
...simulated by...



Fundamental GPT $\mathcal{B} = (B, \Omega_B, E_B)$



Simulations and embeddings

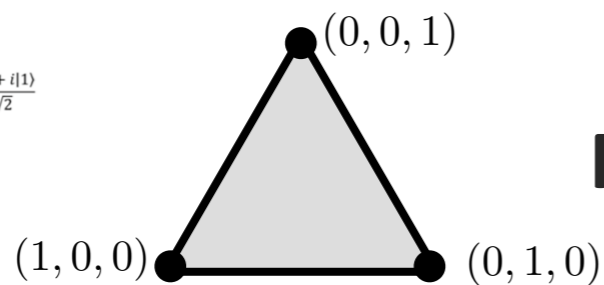
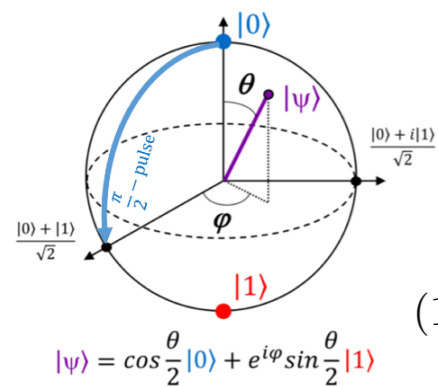


Effective GPT $\mathcal{A} = (A, \Omega_A, E_A)$ found in the lab

...simulated by...



Fundamental GPT $\mathcal{B} = (B, \Omega_B, E_B)$



Effectively preparing state ω_A means **fundamentally** preparing some ω_B , but ω_B may depend on the preparation *procedure*, i.e. the *context*. Collect all those states into a set $\Omega_B(\omega_A) := \{\omega_B\}$.

Simulations and embeddings

Definition. An ε -simulation of effective GPT \mathcal{A} by fundamental GPT \mathcal{B} :

Effective state ω_A \mapsto set of simulating states $\Omega_B(\omega_A)$,

effective effect e_A \mapsto set of simulating effects $E_B(e_A)$,

Simulations and embeddings

Definition. An ε -simulation of effective GPT \mathcal{A} by fundamental GPT \mathcal{B} :

Effective state ω_A \mapsto set of simulating states $\Omega_B(\omega_A)$,

effective effect e_A \mapsto set of simulating effects $E_B(e_A)$,

such that all outcome probabilities are reproduced up to ε :

$$|(\omega_A, e_A) - (\omega_B, e_B)| \leq \varepsilon \text{ for all } \omega_B \in \Omega_B(\omega_A), e_B \in E_B(e_A)$$

and, essentially, mixtures are valid simulations of mixtures (see paper).

Simulations and embeddings

Definition. An ε -**simulation** of effective GPT \mathcal{A} by fundamental GPT \mathcal{B} :

Effective state ω_A \mapsto set of simulating states $\Omega_B(\omega_A)$,

effective effect e_A \mapsto set of simulating effects $E_B(e_A)$,

such that all outcome probabilities are reproduced up to ε :

$$|(\omega_A, e_A) - (\omega_B, e_B)| \leq \varepsilon \text{ for all } \omega_B \in \Omega_B(\omega_A), e_B \in E_B(e_A)$$

and, essentially, mixtures are valid simulations of mixtures (see paper).

Simulation is **univalent** if all $\Omega_B(\omega_A)$, $E_B(e_A)$ contain **one** element.

Simulations and embeddings

Definition. An ε -**simulation** of effective GPT \mathcal{A} by fundamental GPT \mathcal{B} :

Effective state ω_A \mapsto set of simulating states $\Omega_B(\omega_A)$,

effective effect e_A \mapsto set of simulating effects $E_B(e_A)$,

such that all outcome probabilities are reproduced up to ε :

$$|(\omega_A, e_A) - (\omega_B, e_B)| \leq \varepsilon \text{ for all } \omega_B \in \Omega_B(\omega_A), e_B \in E_B(e_A)$$

and, essentially, mixtures are valid simulations of mixtures (see paper).

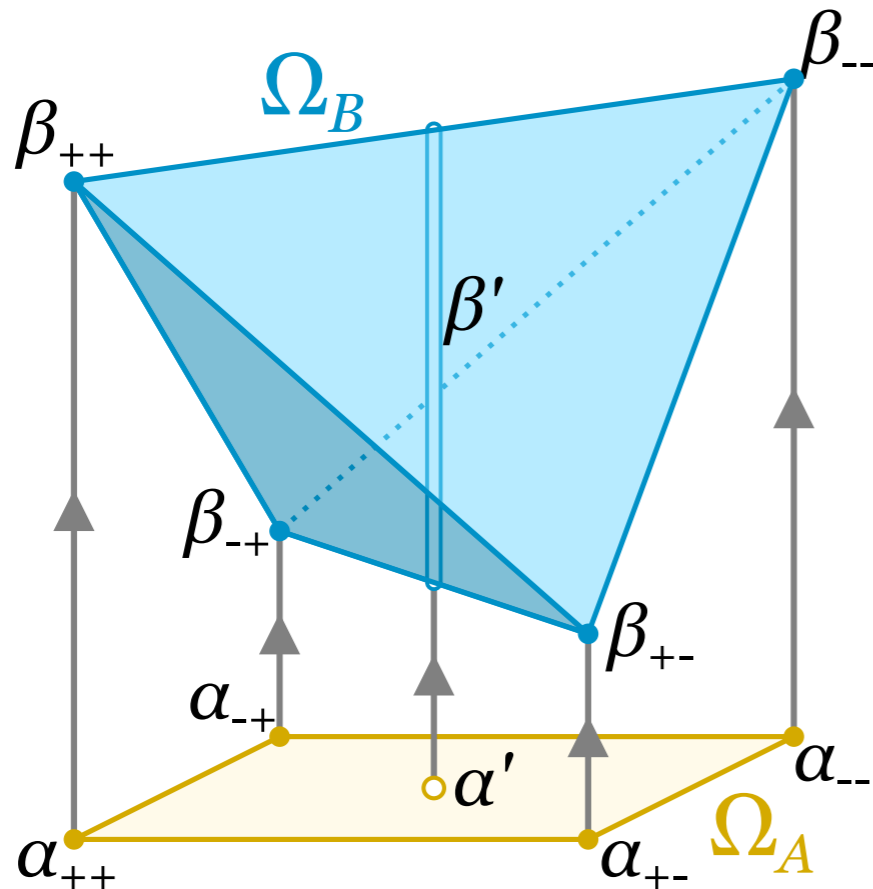
Simulation is **univalent** if all $\Omega_B(\omega_A)$, $E_B(e_A)$ contain **one** element.

Special case $\mathcal{A} = \text{QT}$, $\mathcal{B} = \text{classical probability theory}$:

Simulations are **ontological models**, and univalence = **noncontextuality**.

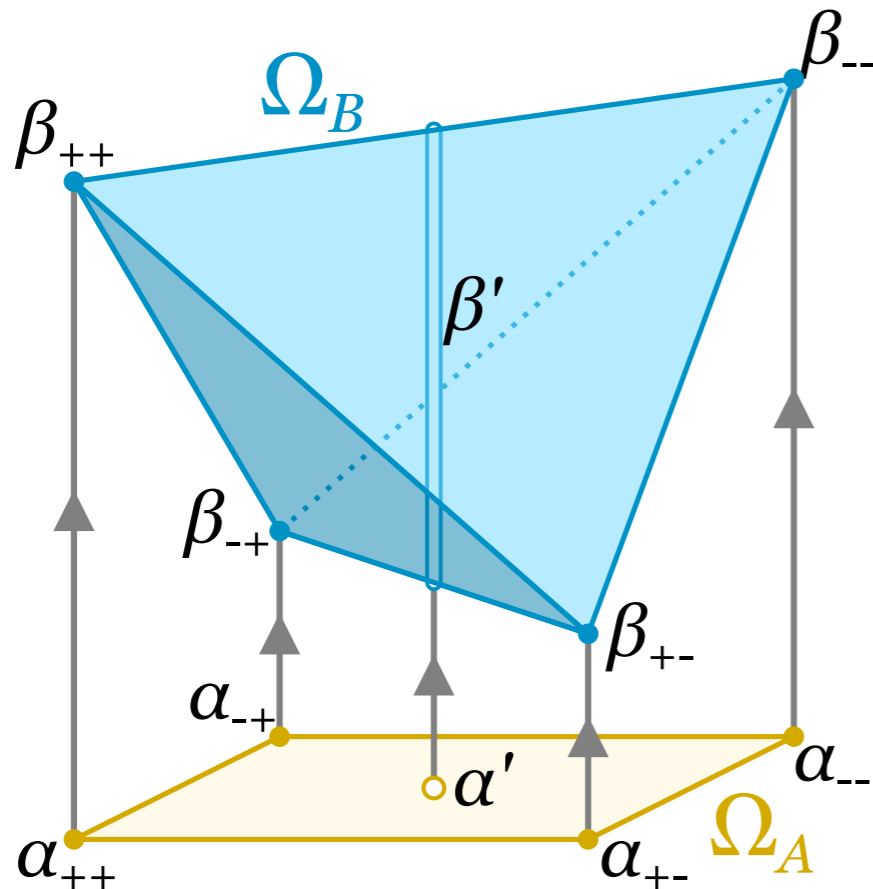
Simulations and embeddings

Example (“Holevo projection”): simulating the gbit $\mathcal{A} = (\mathbb{R}^3, \Omega_A, E_A)$ with a classical 4-level system \mathcal{C}_4 .



Simulations and embeddings

Example (“Holevo projection”): simulating the gbit $\mathcal{A} = (\mathbb{R}^3, \Omega_A, E_A)$ with a classical 4-level system \mathcal{C}_4 .

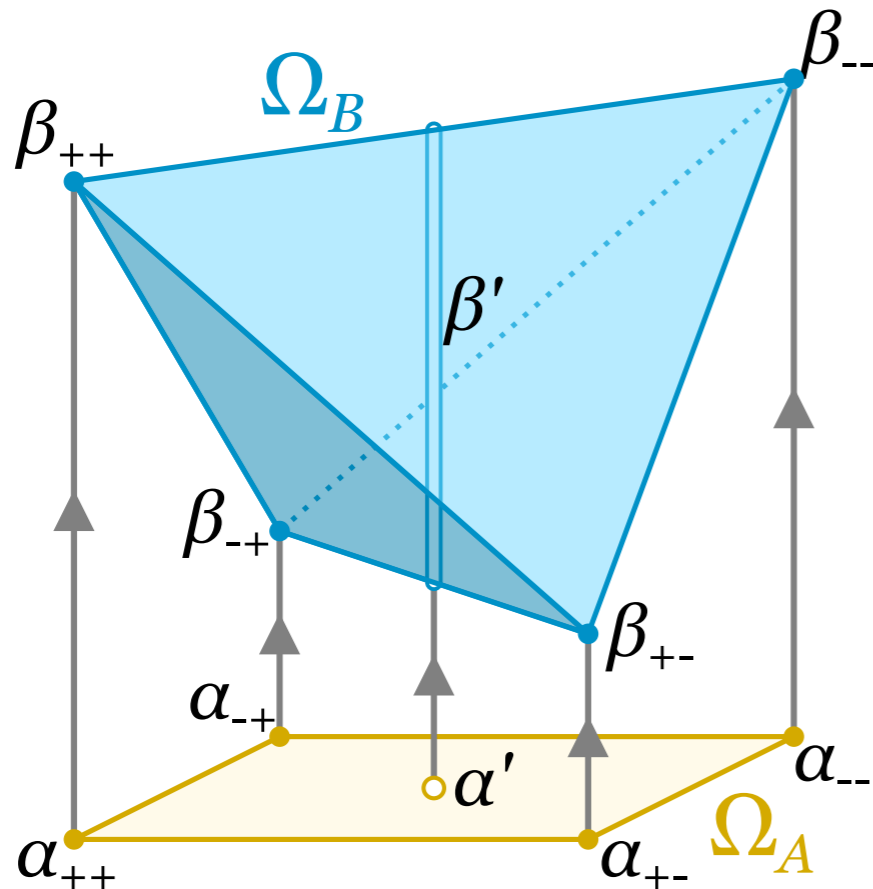


$$\Omega_B(\alpha_{\pm\pm}) = \{\beta_{\pm\pm}\},$$

but $\Omega_B(\alpha')$ = {states β' on blue line}.

Simulations and embeddings

Example (“Holevo projection”): simulating the gbit $\mathcal{A} = (\mathbb{R}^3, \Omega_A, E_A)$ with a classical 4-level system \mathcal{C}_4 .



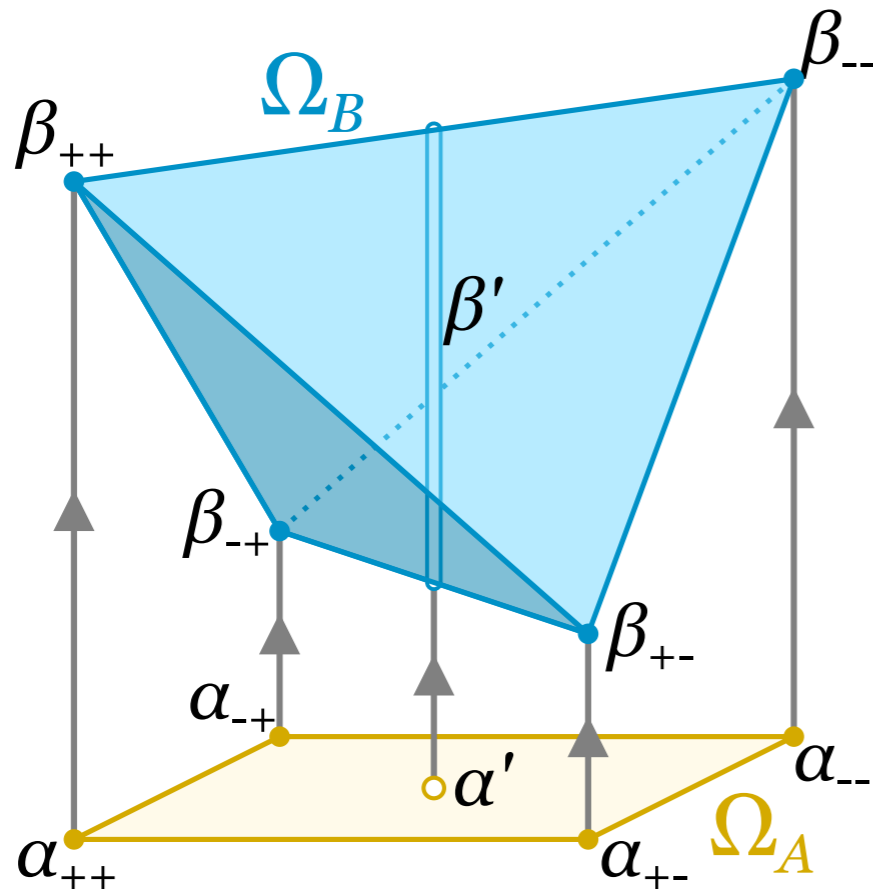
$$\Omega_B(\alpha_{\pm\pm}) = \{\beta_{\pm\pm}\},$$

but $\Omega_B(\alpha')$ = {states β' on blue line}.

(Preparation) contextuality = multivalence:
the fundamental state β' does not only depend on α' , but *must* also depend on the way it has been prepared.

Simulations and embeddings

Example (“Holevo projection”): simulating the gbit $\mathcal{A} = (\mathbb{R}^3, \Omega_A, E_A)$ with a classical 4-level system \mathcal{C}_4 .



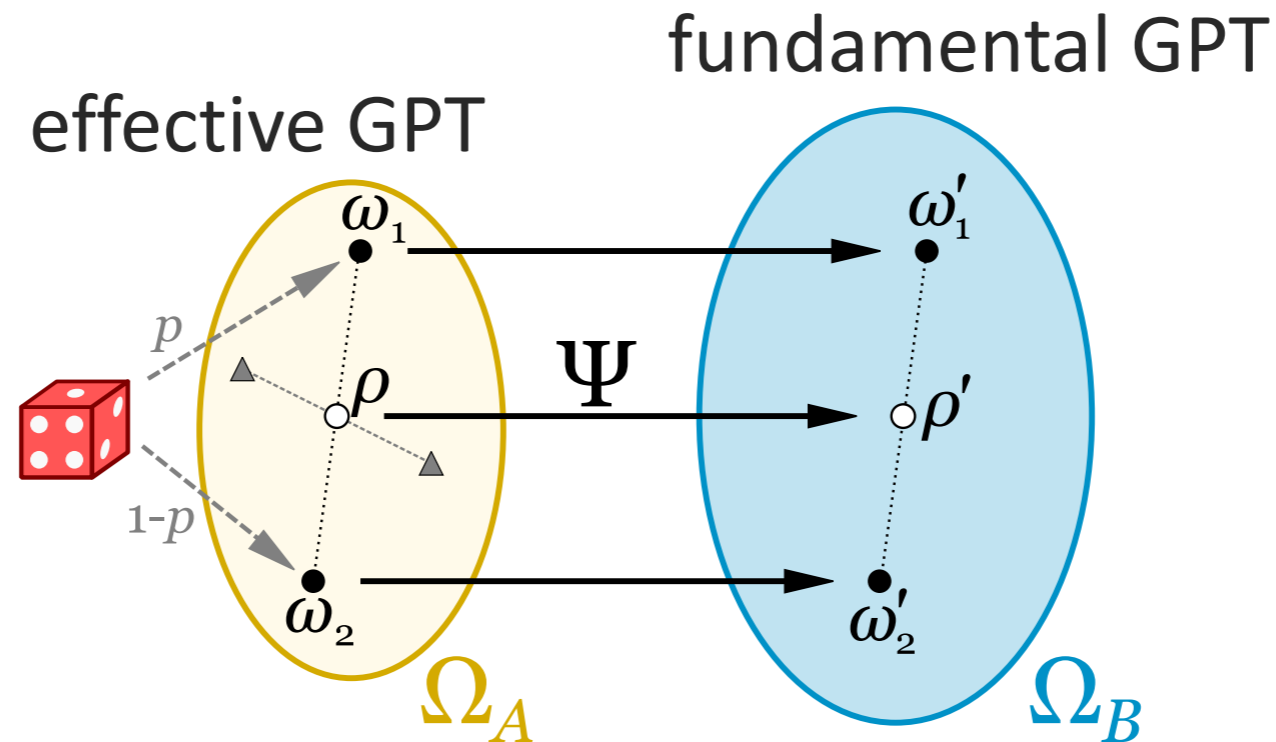
$$\Omega_B(\alpha_{\pm\pm}) = \{\beta_{\pm\pm}\},$$

but $\Omega_B(\alpha') = \{\text{states } \beta' \text{ on blue line}\}.$

(Preparation) contextuality = multivalence:
the fundamental state β' does not only depend on α' , but *must* also depend on the way it has been prepared.

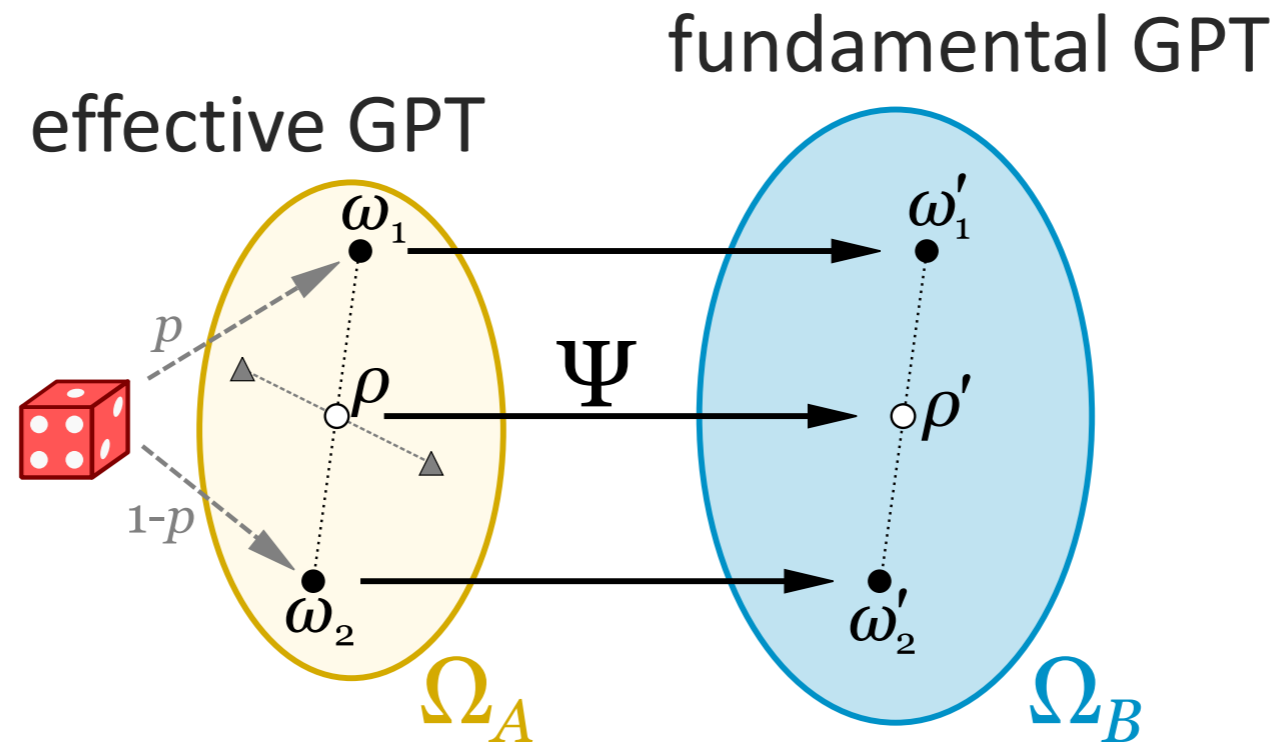
This is an instance of implausible fine-tuning:
the statistical differences among the fundamental states
are miraculously *exactly* “washed out” on the effective level.

Univalent simulations are **embeddings**



Lemma 2. Every univalent ε -simulation of \mathcal{A} by \mathcal{B} defines an ε -embedding of \mathcal{A} into \mathcal{B} , and vice versa.

Univalent simulations are **embeddings**



Lemma 2. Every univalent ε -simulation of \mathcal{A} by \mathcal{B} defines an ε -embedding of \mathcal{A} into \mathcal{B} , and vice versa.

An ε -**embedding** consists of two linear maps Ψ and Φ such that

- Ψ maps the normalized states of \mathcal{A} into those of \mathcal{B} ,
- Φ maps the effects of \mathcal{A} into those of \mathcal{B} ,
- outcome probabilities are preserved up to ε .

Summary of this part

Multivalent simulations (that cannot be made univalent) are **implausible** because they are **fine-tuned**, cf. Holevo projection.

Univalent simulation (of A by B) = **embedding** (of A into B).

Embeddable into CPT (a classical probability simplex) \mathcal{C}_n
= univalently simulatable by fundamental CPT
= **noncontextual** in the sense of **Spekkens**
= plausibly “classical”.

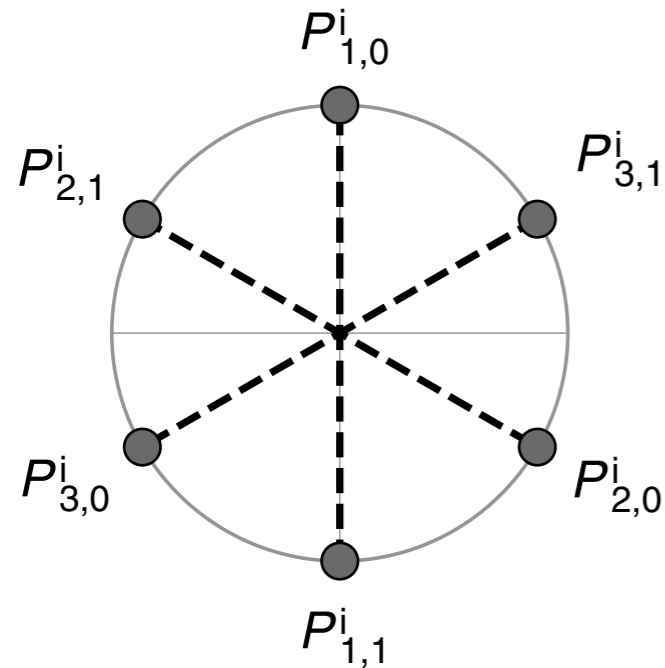
Embeddable into QT (a positive semidefinite cone) \mathcal{Q}_n
= univalently simulatable by fundamental QT
= plausibly “quantum”.

Noncontextual inequalities and approximate embeddings

- [4] M. D. Mazurek et al., *An experimental test of noncontextuality without unphysical idealizations*, Nat. Comm. **7**, 11780 (2016).

Noncontextual inequalities and approximate embeddings

- [4] M. D. Mazurek et al., *An experimental test of noncontextuality without unphysical idealizations*, Nat. Comm. **7**, 11780 (2016).



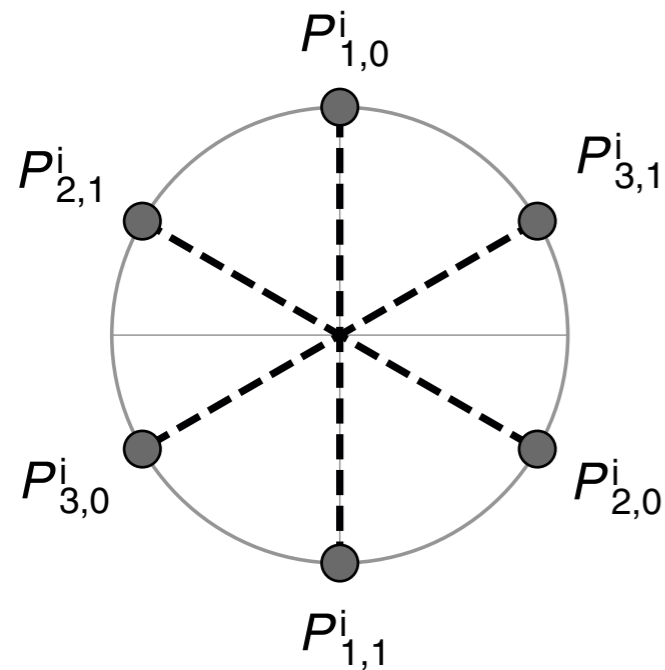
The qubit (actually, rebit) does not have a noncontextual ontological model.

Quantitative statement:

$$A := \frac{1}{6} \sum_{t \in \{1,2,3\}} \sum_{b \in \{0,1\}} P(b | p_{t,b}, m_t) \leq \frac{5}{6}.$$

Noncontextual inequalities and approximate embeddings

[4] M. D. Mazurek et al., *An experimental test of noncontextuality without unphysical idealizations*, Nat. Comm. **7**, 11780 (2016).



The qubit (actually, rebit) does not have a noncontextual ontological model.

Quantitative statement:

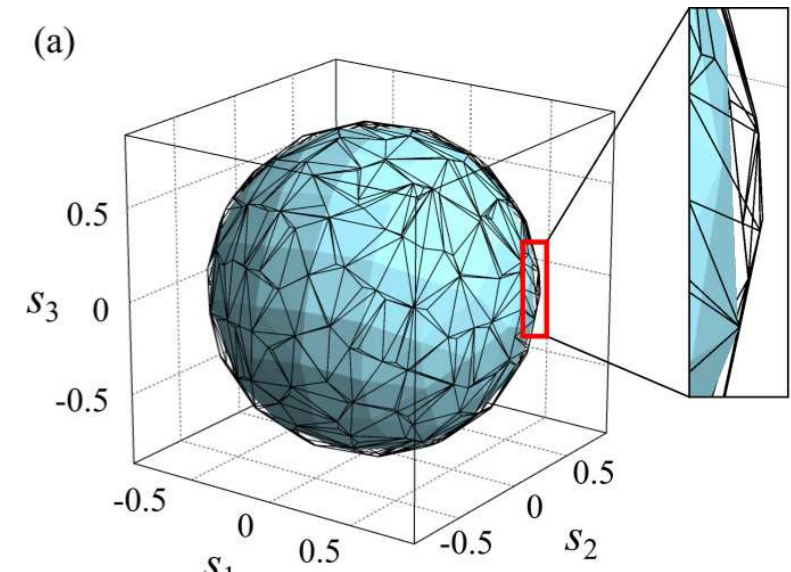
$$A := \frac{1}{6} \sum_{t \in \{1,2,3\}} \sum_{b \in \{0,1\}} P(b | p_{t,b}, m_t) \leq \frac{5}{6}.$$

These imply bounds on the approximate embeddability into classical:

Example 1. *Let $\varepsilon < \frac{1}{6}$. Then the rebit (and thus, also the qubit) cannot be ε -embedded into any \mathcal{C}_n .*

Overview

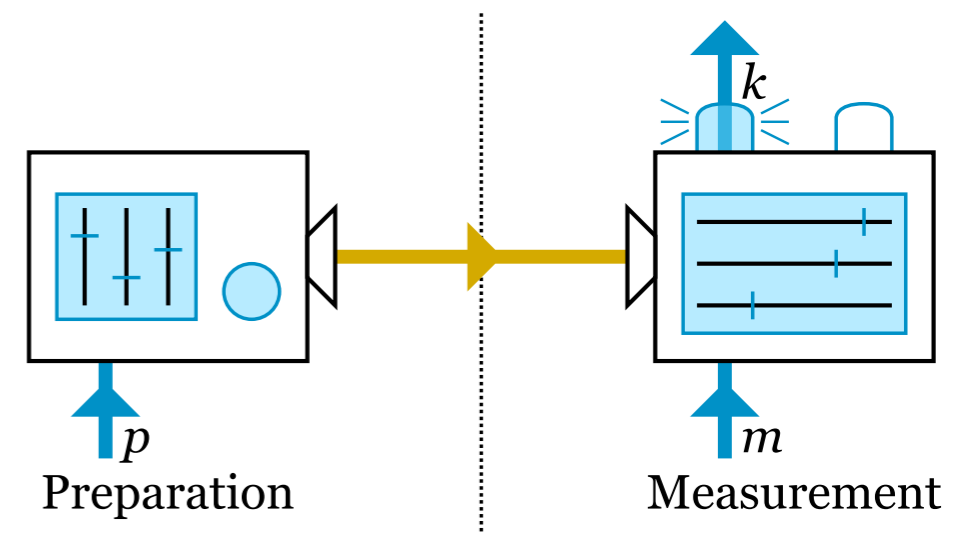
1. GPTs and theory-agnostic tomography



2. Contextuality, simulations, and embeddings

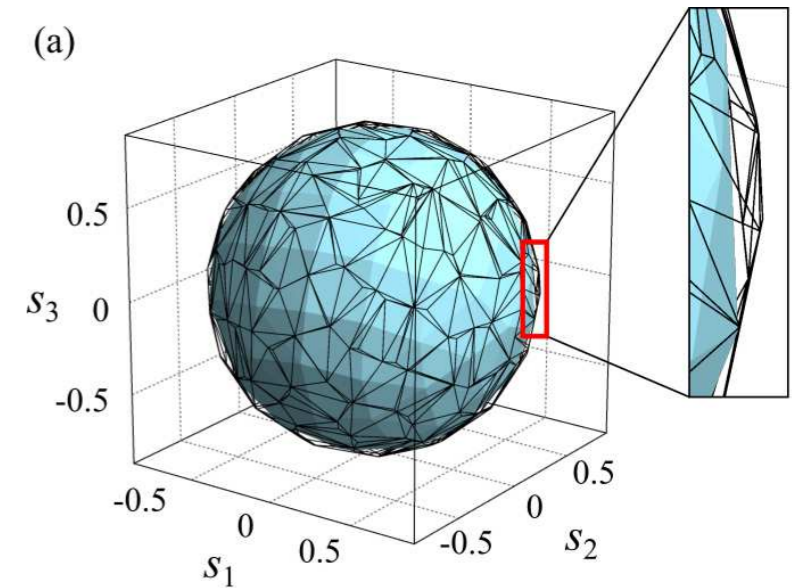
3. Exact embeddings into quantum theory

4. An experimental test of QT



Overview

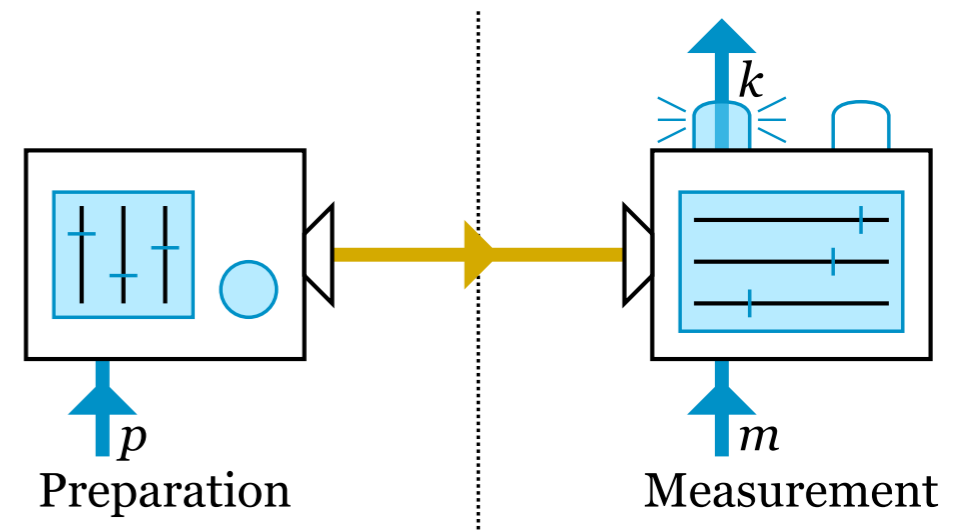
1. GPTs and theory-agnostic tomography



2. Contextuality, simulations, and embeddings

3. Exact embeddings into quantum theory

4. An experimental test of QT



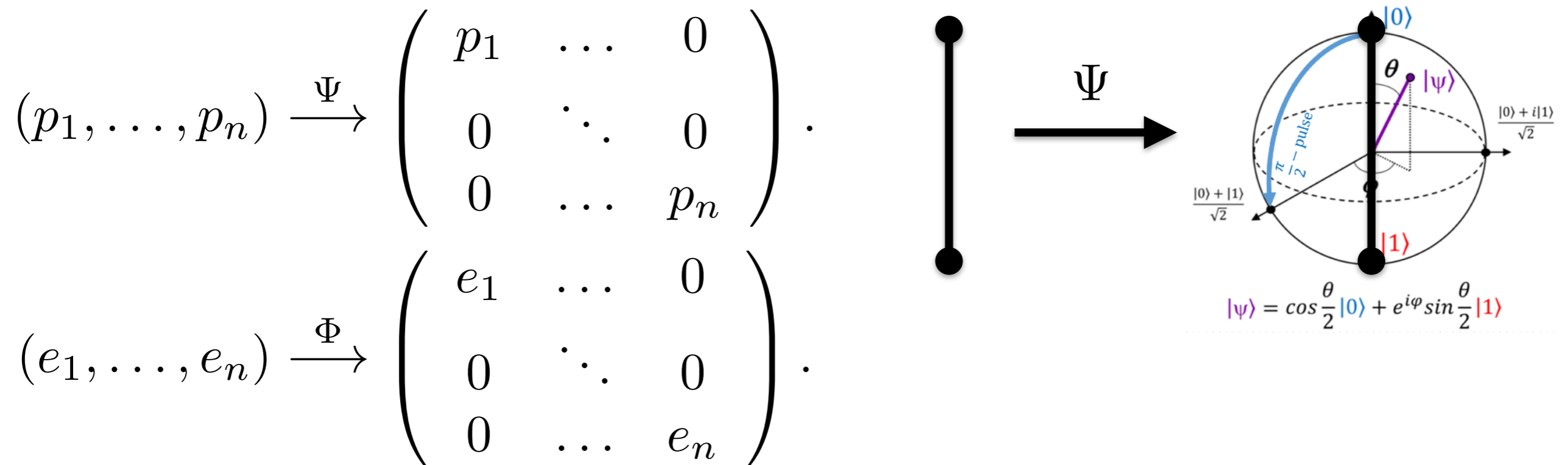
3. Exact embeddings into quantum theory

Which GPTs admit of an **univalent** (“noncontextual”) **simulation by QT**, i.e. can be embedded into QT \mathcal{Q}_n (say, exactly)?

3. Exact embeddings into quantum theory

Which GPTs admit of an **univalent** (“noncontextual”) simulation by QT, i.e. can be embedded into QT \mathcal{Q}_n (say, exactly)?

Example: Classical PT can be embedded into QT.



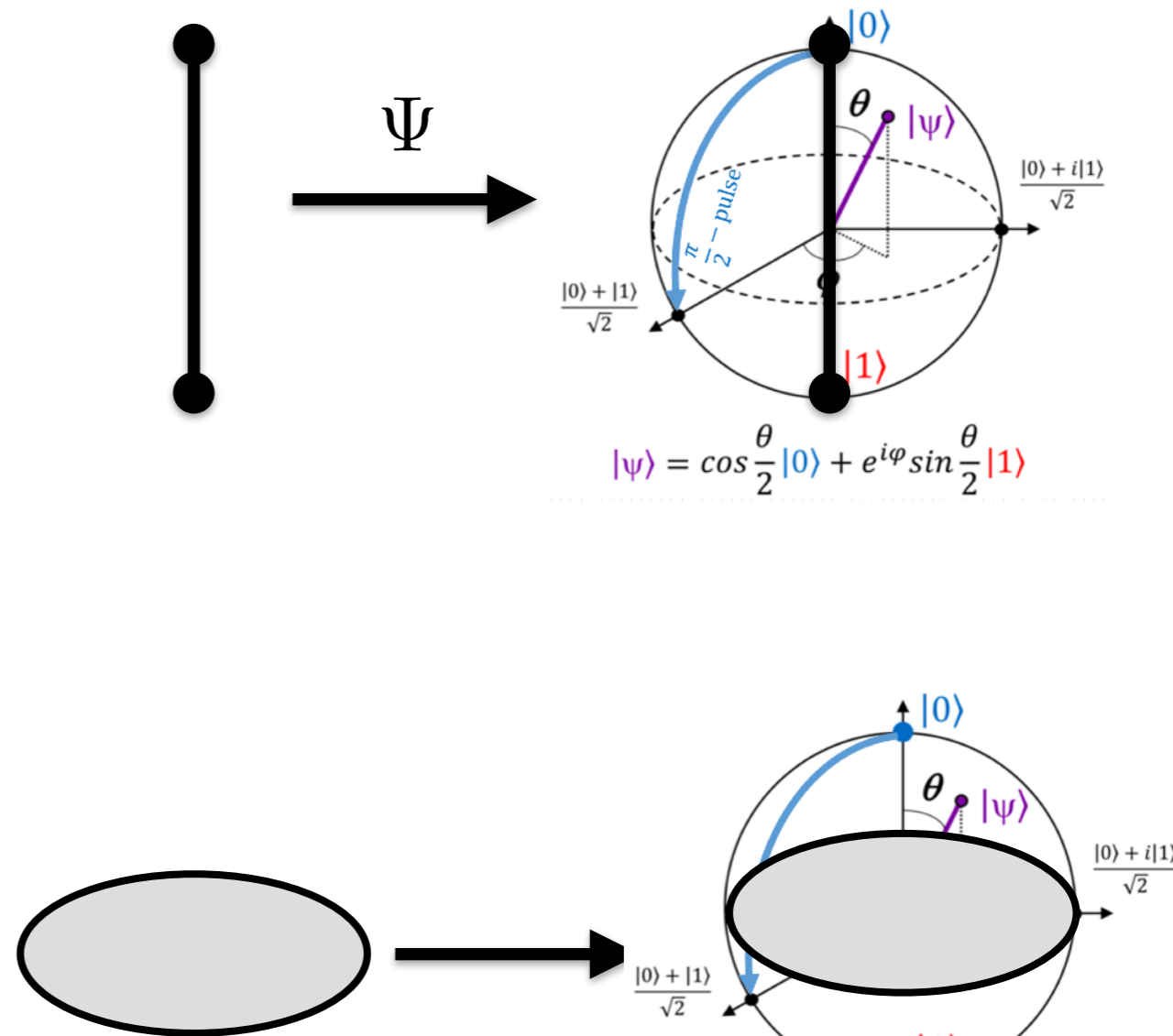
3. Exact embeddings into quantum theory

Which GPTs admit of an **univalent** (“noncontextual”) simulation by QT, i.e. can be embedded into QT \mathcal{Q}_n (say, exactly)?

Example: Classical PT can be embedded into QT.

$$(p_1, \dots, p_n) \xrightarrow{\Psi} \begin{pmatrix} p_1 & \dots & 0 \\ 0 & \ddots & 0 \\ 0 & \dots & p_n \end{pmatrix} \cdot$$

$$(e_1, \dots, e_n) \xrightarrow{\Phi} \begin{pmatrix} e_1 & \dots & 0 \\ 0 & \ddots & 0 \\ 0 & \dots & e_n \end{pmatrix} \cdot$$



Similarly, **QT over the real numbers** can be embedded into QT.

3. Exact embeddings into quantum theory

Focus on the “**unrestricted** GPTs” where all vectors yielding valid probabilities on all states are effects: $\mathcal{A} = (A, \Omega_A, E_A)$

$$E_A = \{e \in A \mid 0 \leq \langle \omega, e \rangle \leq 1 \text{ for all } \omega \in \Omega_A\}.$$

3. Exact embeddings into quantum theory

Focus on the “**unrestricted** GPTs” where all vectors yielding valid probabilities on all states are effects: $\mathcal{A} = (A, \Omega_A, E_A)$

$$E_A = \{e \in A \mid 0 \leq \langle \omega, e \rangle \leq 1 \text{ for all } \omega \in \Omega_A\}.$$

Theorem 2. *An unrestricted GPT can be exactly embedded into finite-dimensional quantum theory if and only if it corresponds to a special Euclidean Jordan algebra.*

3. Exact embeddings into quantum theory

Focus on the “**unrestricted** GPTs” where all vectors yielding valid probabilities on all states are effects: $\mathcal{A} = (A, \Omega_A, E_A)$

$$E_A = \{e \in A \mid 0 \leq \langle \omega, e \rangle \leq 1 \text{ for all } \omega \in \Omega_A\}.$$

Theorem 2. *An unrestricted GPT can be exactly embedded into finite-dimensional quantum theory if and only if it corresponds to a special Euclidean Jordan algebra.*

- QT over real numbers \mathbb{R} , complex numbers \mathbb{C} , **quaternions** \mathbb{H} ,
- d -dimensional **Bloch ball** state spaces,
- direct sums of those, including **CPT** and QT with **superselection rules**.

3. Exact embeddings into quantum theory

Focus on the “**unrestricted** GPTs” where all vectors yielding valid probabilities on all states are effects: $\mathcal{A} = (A, \Omega_A, E_A)$

$$E_A = \{e \in A \mid 0 \leq \langle \omega, e \rangle \leq 1 \text{ for all } \omega \in \Omega_A\}.$$

Theorem 2. *An unrestricted GPT can be exactly embedded into finite-dimensional quantum theory if and only if it corresponds to a special Euclidean Jordan algebra.*

- QT over real numbers \mathbb{R} , complex numbers \mathbb{C} , **quaternions** \mathbb{H} ,
- d -dimensional **Bloch ball** state spaces,
- direct sums of those, including **CPT** and QT with **superselection rules**.

These are the **only unrestricted GPTs** that are “plausibly quantum”.

3. Exact embeddings into quantum theory

Focus on the “**unrestricted** GPTs” where all vectors yielding valid probabilities on all states are effects: $\mathcal{A} = (A, \Omega_A, E_A)$

$$E_A = \{e \in A \mid 0 \leq \langle \omega, e \rangle \leq 1 \text{ for all } \omega \in \Omega_A\}.$$

Theorem 2. *An unrestricted GPT can be exactly embedded into finite-dimensional quantum theory only if it corresponds to a special Euclidean space.*

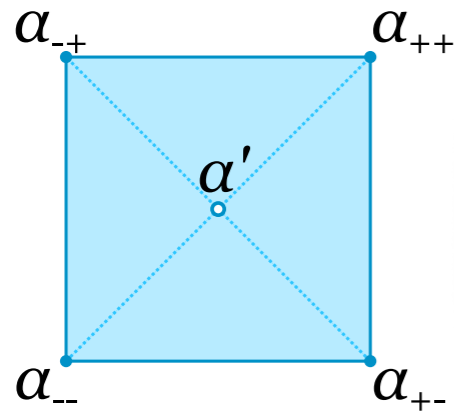
Needs 2^d -dim.
Hilbert space for
simulation!

- QT over real numbers \mathbb{R} , complex numbers \mathbb{C} , **quaternions** \mathbb{H} ,
- **d -dimensional Bloch ball state spaces,**
- direct sums of those, including **CPT** and QT with **superselection rules.**

These are the **only unrestricted GPTs** that are “plausibly quantum”.

Non-exact embeddings into quantum theory

Example: the gbit

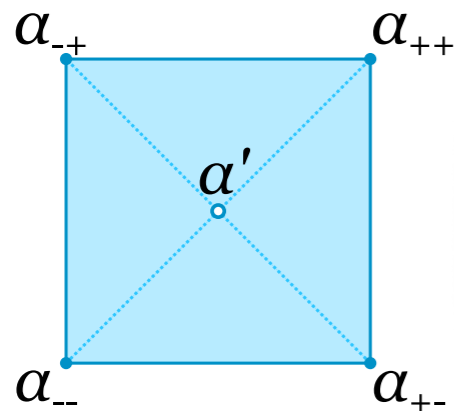


Example 2. *Let $\varepsilon \leq 0.1014$. Then the gbit cannot be ε -embedded into any \mathcal{Q}_n or \mathcal{Q}_∞ .*

There is no better-than-10% univalent (“noncontextual”) simulation by QT.

Non-exact embeddings into quantum theory

Example: the gbit

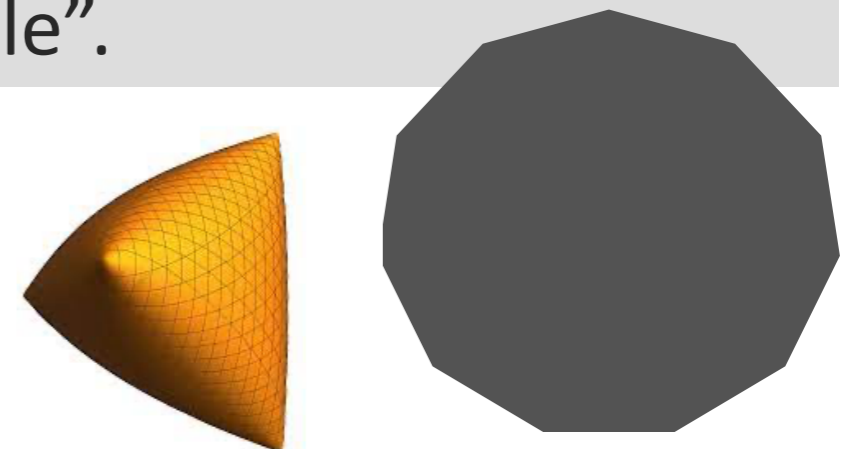


Example 2. Let $\varepsilon \leq 0.1014$. Then the gbit cannot be ε -embedded into any \mathcal{Q}_n or \mathcal{Q}_∞ .

There is no better-than-10% univalent (“noncontextual”) simulation by QT.

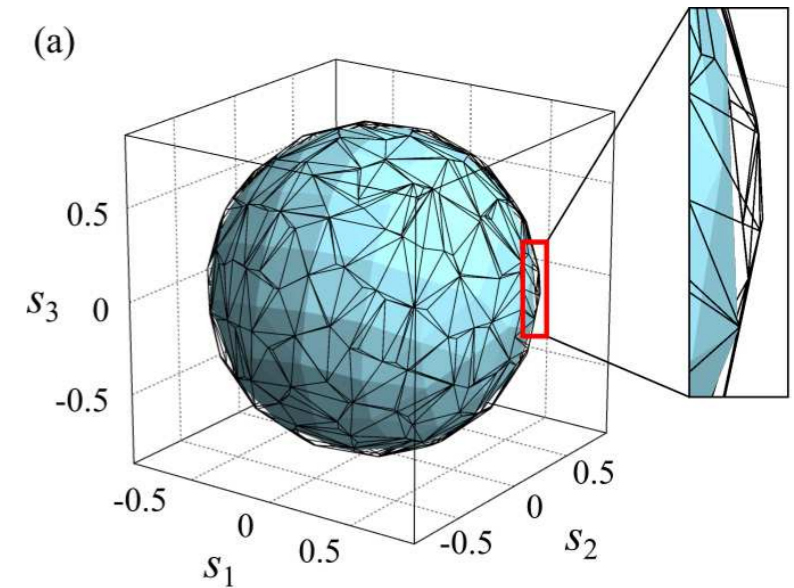
Also shown in our paper:

can **use known results on Bell inequalities** to certify nonembeddability. Impractical and inefficient, but “proof of principle”.



Overview

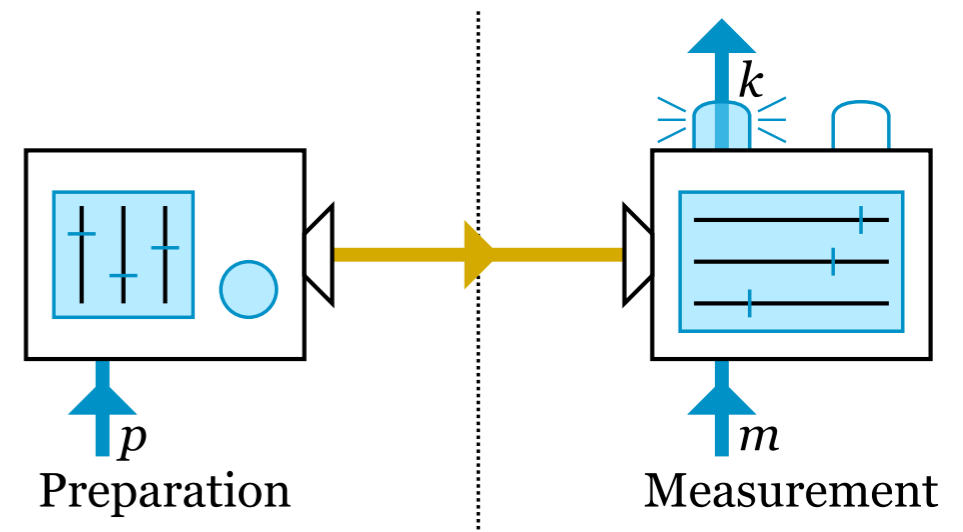
1. GPTs and theory-agnostic tomography



2. Contextuality, simulations, and embeddings

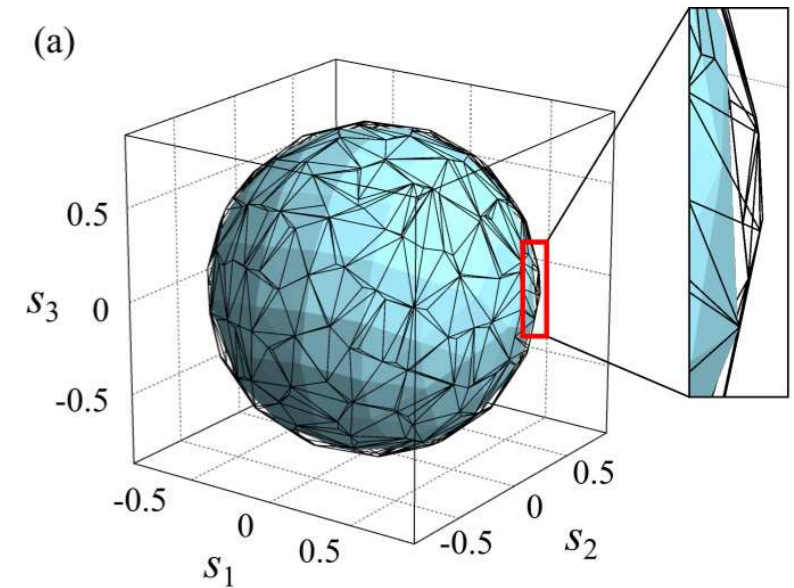
3. Exact embeddings into quantum theory

4. An experimental test of QT



Overview

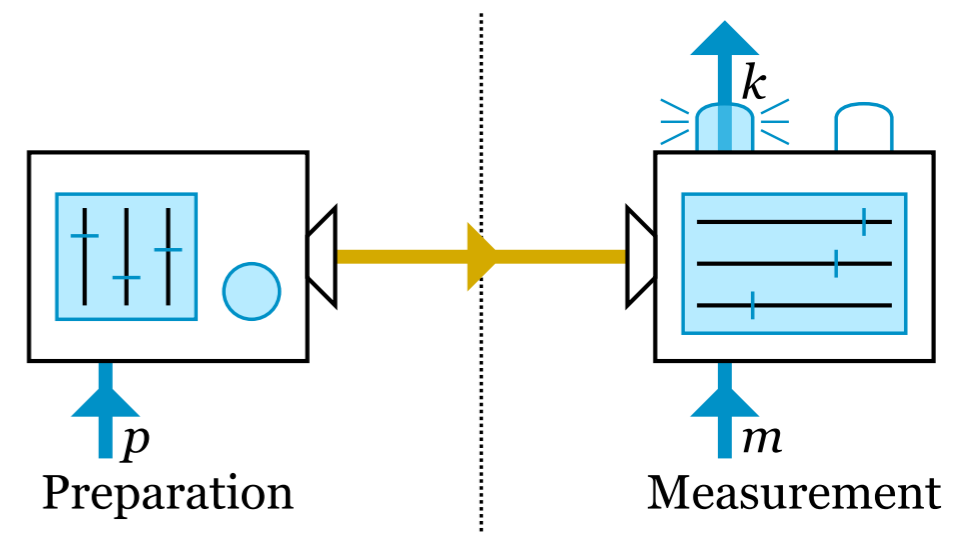
1. GPTs and theory-agnostic tomography



2. Contextuality, simulations, and embeddings

3. Exact embeddings into quantum theory

4. An experimental test of QT



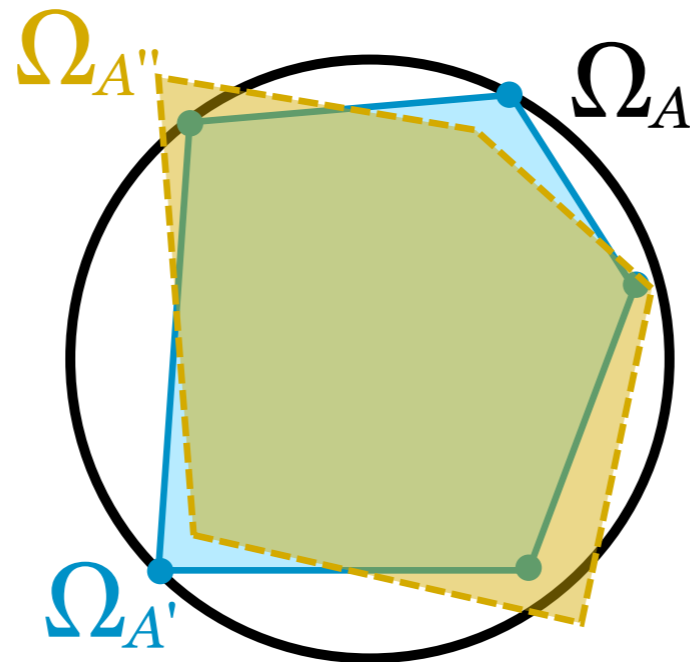
Suggestion: experimental test of QT

Suggestion: experimental test of QT

Perform theory-agnostic tomography on an **effective physical system** in your laboratory.

Test whether the resulting effective GPT is ε -**embeddable into QT**, where ε is a function of the experimental uncertainty.

If, surprisingly, “no”, then this challenges QT.

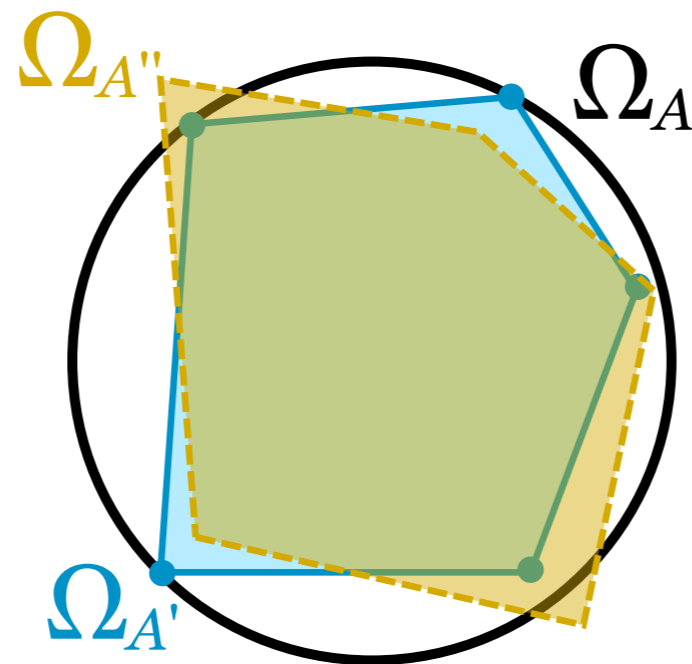


Suggestion: experimental test of QT

Perform theory-agnostic tomography on an **effective physical system** in your laboratory.

Test whether the resulting effective GPT is ε -**embeddable into QT**, where ε is a function of the experimental uncertainty.

If, surprisingly, “no”, then this challenges QT.



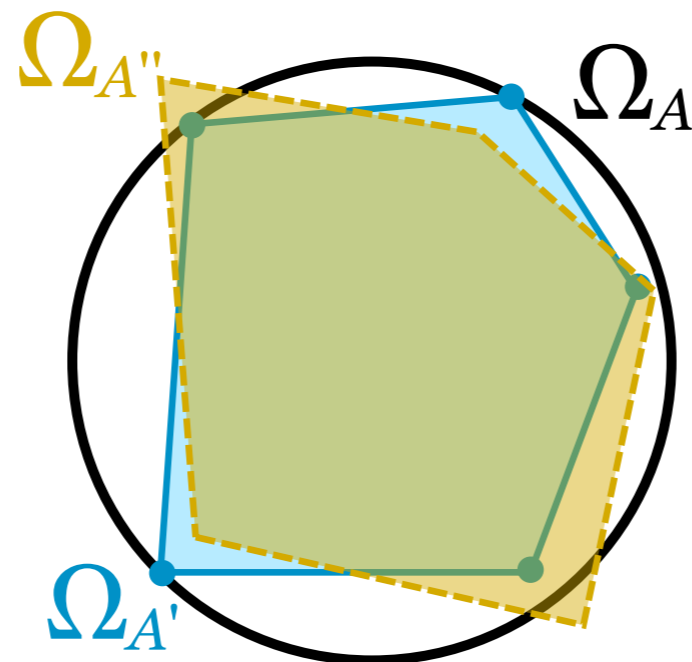
Typically a
restricted GPT.

Suggestion: experimental test of QT

Perform theory-agnostic tomography on an **effective physical system** in your laboratory.

Test whether the resulting effective GPT is ε -**embeddable into QT**, where ε is a function of the experimental uncertainty.

If, surprisingly, “no”, then this challenges QT.



Typically a
restricted GPT.

A quantum explanation of the result is then **similarly implausible** as a classical (contextual) explanation of the quantum state space.

Summary

- Have generalized Spekkens' notion of generalized noncontextuality:
“Processes that are statistically indistinguishable in an effective theory should not require explanation by multiple distinguishable processes in a more fundamental theory.”
- Results: Several structural insights, a new experimental test of QT, Jordan algebras are the only **unrestricted** GPTs embeddable into QT.
- Note: the experiments will **not** just test QT “against other probabilistic theories / GPTs”, **but against arbitrary** modifications impacting prepare-and-measure-statistics. We use GPTs only as a tool to analyze the latter.

arXiv:2112.09719 (update soon),
to appear in Physical Review X.

Thank you!