## How spacetime constrains the structure of quantum theory

Caroline L. Jones, Stefan L. Ludescher, Albert Aloy, Andrew J. P. Garner, Oscar C. O. Dahlsten, Markus P. Müller IQOQI Vienna \& Perimeter Institute


Der Wissenschaftsfonds.

## Overview

1. Motivations: QG and device-independent QIT
2. Relativity of simultaneity and the qubit
3. Randomness generation via rotational symmetry
4. Conclusions

## Overview

1. Motivations: QG and device-independent QIT
2. Relativity of simultaneity and the qubit
3. Randomness generation via rotational symmetry
4. Conclusions

## Quantum gravity: an analogy

## Quantum gravity: an analogy

## Wanted: a complete theory of evolution.

## Quantum gravity: an analogy

## Wanted: a complete theory of evolution.

## biological traits $\longleftrightarrow$ environment

## Quantum gravity: an analogy

Wanted: a complete theory of evolution.

## biological traits $\longleftrightarrow$ environment


confined to desert, no fossils: sparse empirical evidence.

## Quantum gravity: an analogy

Wanted: a complete theory of evolution.

## biological traits $\longleftrightarrow$ environment


confined to desert, no fossils: sparse empirical evidence.

- Option 1: try to develop a full-blown theory directly.
- Option 2: first, study how the environment constrains biological traits.


## Quantum gravity: an analogy

Wanted: a complete theory of evolution.

## biological traits $\longleftrightarrow$ environment


confined to desert, no fossils: sparse empirical evidence.

- Option 1: try to develop a full-blown theory directly.
- Option 2: first, study how the environment constrains biological traits. Needs imagination of how biology could be different.


## Quantum gravity: an analogy

## Quantum gravity: an analogy

Wanted: a complete theory of quantum gravity.

## Quantum gravity: an analogy

Wanted: a complete theory of quantum gravity. sparse empirical evidence.

## Quantum gravity: an analogy

Wanted: a complete theory of quantum gravity. sparse empirical evidence.

- Option 1: try to develop a full-blown theory directly.
- Option 2: first, study how spacetime constrains quantum theory.


## Quantum gravity: an analogy

Wanted: a complete theory of quantum gravity. sparse empirical evidence.

- Option 1: try to develop a full-blown theory directly.
- Option 2: first, study how spacetime constrains quantum theory. Needs (mathematical) imagination of how the universe's probabilistic theory could be different.


## Quantum gravity: an analogy

Wanted: a complete theory of quantum gravity. sparse empirical evidence.

- Option 1: try to develop a full-blown theory directly.

- Option 2: first, study how spacetime constrains quantum theory. Needs (mathematical) imagination of how the universe's probabilistic theory could be different.

superstrong nonlocality?

higher-order interference?


## Further motivation: (semi-)device-independent QIT

Goal: Generate certified random bits.

## Further motivation: (semi-)device-independent QIT

## Goal: Generate certified random bits.

Why not just send single photons on a half-silvered mirror?


## Further motivation: (semi-)device-independent QIT

Goal: Generate certified random bits.
Why not just send single photons on a half-silvered mirror?


## Further motivation: (semi-)device-independent QIT

Goal: Generate certified random bits.
Why not just send single photons on a half-silvered mirror?


Device-independent randomness expansion:
Violation of Bell inequality $\Rightarrow$ outcomes uncorrelated with rest of the world
See e.g.: A. Acín, Randomness and quantum non-locality, QCRYPT 2012 talk. V. Scarani, Bell nonlocality, Oxford Graduate Texts (2019).

Semi-device-independent (SDI): allow communication, add assumption.

## Further motivation: (semi-)device-independent QIT

Semi-device-independent (SDI): allow communication, add assumption.


## Further motivation: (semi-)device-independent QIT

Semi-device-independent (SDI): allow communication, add assumption.


Observed correlations $p(a \mid x, y)$ imply $H(A \mid X, Y, \Lambda) \gg 0$.

## Further motivation: (semi-)device-independent QIT

Semi-device-independent (SDI): allow communication, add assumption.


Observed correlations $p(r \cdot \mid x, y)$ imply $H(A \mid X, Y, \Lambda) \gg 0$.

Drawback assumption not physically well-motivated \& requires QT.

## Further motivation: (semi-)device-independent QIT

Semi-device-independent (SDI): allow communication, add assumption.


Observation: in many experiments, settings are spatiotemporal quantities.

## Further motivation: (semi-)device-independent QIT

Semi-device-independent (SDI): allow communication, add assumption.


Observed correlations $p(a \mid x, y)$ imply $H(A \mid X, Y, \Lambda) \gg 0$.

Drawback: assumption not physically well-motivated \& requires QT.
Observation: in many experiments, settings are spatiotemporal quantities.
Idea: reformulate in terms of spacetime symmetries, w/o assuming QT. Can quantum phenomenology / functionality be reproduced?

## Overview

1. Motivations: QG and device-independent QIT
2. Relativity of simultaneity and the qubit
3. Randomness generation via rotational symmetry
4. Conclusions

## Overview

1. Motivations: QG and device-independent QIT
2. Relativity of simultaneity and the qubit
3. Randomness generation via rotational symmetry
4. Conclusions

Imagine what the quantum bit could be instead...


## Imagine what the quantum bit could be instead...


P. Jordan, J. von Neumann, E. Wigner, On an algebraic generalization of the quantum mechanical formalism, Annals of Mathematics 35, 29-64 (1934).

## Imagine what the quantum bit could be instead...



## Imagine what the quantum bit could be instead...


P. Jordan, J. von Neumann, E. Wigner, On an algebraic generalization of the quantum mechanical formalism, Annals of Mathematics 35, 29-64 (1934).


## Imagine what the quantum bit could be instead...


P. Jordan, J. von Neumann, E. Wigner, On an algebraic generalization of the quantum mechanical formalism, Annals of Mathematics 35, 29-64 (1934).

bit
$\mathbb{R}$-qubit

Take care:
qutrit etc.
not a ball!

Does spacetime constrain d?
"Why" d=3 ?

## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).


## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).


## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).


## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).


North-pole state: particle definitely in upper branch.

## Constraints from relativity

```
A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).
```



South-pole state: particle definitely in lower branch.

## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).

d-dim. "Bloch sphere"

State on equator $z=0$ : probability $1 / 2$ for each.

## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).

d-dim. "Bloch sphere"

State on equator $z=0$ : probability $1 / 2$ for each.
$p(u p)=\frac{1}{2}(z+1)$

## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).

d-dim. "Bloch sphere"

What transformations $T$ can we perform locally in one arm...
... reversibly, i.e. without any information loss?

## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).

d-dim. "Bloch sphere"
$T$ must be a rotation of the Bloch ball (reversible+linear)...
... and must preserve $p$ (up), i.e. preserve the $z$-axis.

## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).

d-dim. "Bloch sphere"
$T$ must be a rotation of the Bloch ball (reversible+linear)...
... and must preserve $p$ (up), i.e. preserve the $z$-axis.

## Constraints from relativity

A. Garner, MM, O. C. O. Dahlsten, Proc. R. Soc. A 473, 20170596 (2017).

d-dim. "Bloch sphere"
$T$ must be a rotation of the Bloch ball (reversible+linear)...
... and must preserve $p$ (up), i.e. preserve the $z$-axis.

$$
\mathcal{G}_{A}=\mathcal{G}_{B} \simeq \operatorname{SO}(d-1) .
$$

## Constraints from relativity



Relativity: there's a frame of reference in which $T_{A}$ happens before $T_{B \ldots}$

$$
\mathcal{G}_{A}=\mathcal{G}_{B} \simeq \operatorname{SO}(d-1)
$$

## Constraints from relativity



Relativity: there's a frame of reference in which $T_{A}$ happens before $T_{B} \ldots$ ... and another frame where it's the other way around.

$$
\mathcal{G}_{A}=\mathcal{G}_{B} \simeq \operatorname{SO}(d-1)
$$

## Constraints from relativity

$$
\Rightarrow T_{A} T_{B}=T_{B} T_{A} \text { for all } T_{A}, T_{B} \in \mathrm{SO}(d-1)
$$

d-dim. "Bloch sphere"

Relativity: there's a frame of reference in which $T_{A}$ happens before $T_{B \ldots}$ ... and another frame where it's the other way around.

$$
\mathcal{G}_{A}=\mathcal{G}_{B} \simeq \operatorname{SO}(d-1) .
$$

## Constraints from relativity

$\Rightarrow T_{A} T_{B}=T_{B} T_{A}$ for all $T_{A}, T_{B} \in \mathrm{SO}(d-1)$.
$\Rightarrow d \leq 3$.

d-dim. "Bloch sphere"

Relativity: there's a frame of reference in which $T_{A}$ happens before $T_{B \ldots}$ ... and another frame where it's the other way around.

$$
\mathcal{G}_{A}=\mathcal{G}_{B} \simeq \operatorname{SO}(d-1) .
$$



So far, we assumed: $\mathcal{G}_{A}=\mathcal{G}_{B}$. Assumption of relationality.



So far, we assumed: $\mathcal{G}_{A}=\mathcal{G}_{B}$. Assumption of relationality. Whatever happens in one arm can be undone in the other arm.

## Constraints from relativity

Let's relax this assumption to $\mathcal{G}_{A} \simeq \mathcal{G}_{B}$.
$\Rightarrow d=5$. Quaternionic QM survives.


So far, we assumed: $\mathcal{G}_{A}=\mathcal{G}_{B}$. Assumption of relationality.
Whatever happens in one arm can be undone in the other arm.

## Classification of possibilities

A1) Beam splitter can prepare any upper-branch probability $p$.
A2) Every pure state with the same $p$ can be prepared by reversible operations applied locally on the two arms.
$A 3$ ) The groups of operations of $A$ and $B$ are isomorphic.

## Classification of possibilities

A1) Beam splitter can prepare any upper-branch probability $p$.
A2) Every pure state with the same $p$ can be prepared by reversible operations applied locally on the two arms.
$A 3$ ) The groups of operations of $A$ and $B$ are isomorphic.

Theorem 6.2. Under the assumptions A1, A2, A3, relativity of simultaneity (REL) allows for the following possibilities and not more:
$-d=1$ (the classical bit), with $\mathcal{G}_{\mathrm{A}}=\mathcal{G}_{\mathrm{B}}=\{\mathbf{1}\}$ (i.e. without any non-trivial local transformations),
$-d=2$ (the quantum bit over the real numbers), with $\mathcal{G}_{\mathrm{A}}=\mathcal{G}_{\mathrm{B}}=\mathbb{Z}_{2}$,
$-d=3$ (the standard quantum bit over the complex numbers), with $\mathcal{G}_{\mathrm{A}}=\mathcal{G}_{\mathrm{B}}=\mathrm{SO}(2)=\mathrm{U}(1)$,
$-d=5$ (the quaternionic quantum bit), with $\mathcal{G}_{\mathrm{AB}}=\mathrm{SO}(4), \mathcal{G}_{\mathrm{A}}$ the left- and $\mathcal{G}_{\mathrm{B}}$ the right-isoclinic rotations in $\mathrm{SO}(4)$ (or vice versa) which are both isomorphic to $\mathrm{SU}(2)$, and $\mathcal{G}_{\mathrm{A}} \cap \mathcal{G}_{\mathrm{B}}=\{+\mathbb{I},-\mathbb{I}\}$.

## Classification of possibilities

A1) Beam splitter can prepare any upper-branch probability $p$.
A2) Every pure state with the same $p$ can be prepared by reversible operations applied locally on the two arms.
$A 3$ ) The groups of operations of $A$ and $B$ are isomorphic.

Theorem 6.2. Under the assumptions A1, A2, A3, relativity of simultaneity (REL) allows for the following possibilities and not more:
$-d=1$ (the classical bit), with $\mathcal{G}_{\mathrm{A}}=\mathcal{G}_{\mathrm{B}}=\{\mathbf{1}\}$ (i.e. without any non-trivial local transformations),
$-d=2$ (the quantum bit over the real numbers), with $\mathcal{G}_{\mathrm{A}}=\mathcal{G}_{\mathrm{B}}=\mathbb{Z}_{2}$,
$-d=3$ (the standard quantum bit over the complex numbers), with $\mathcal{G}_{\mathrm{A}}=\mathcal{G}_{\mathrm{B}}=\mathrm{SO}(2)=\mathrm{U}(1)$,
$-d=5$ (the quaternionic quantum bit), with $\mathcal{G}_{\mathrm{AB}}=\mathrm{SO}(4), \mathcal{G}_{\mathrm{A}}$ the left- and $\mathcal{G}_{\mathrm{B}}$ the right-isoclinic rotations in $\mathrm{SO}(4)$ (or vice versa) which are both isomorphic to $\mathrm{SU}(2)$, and $\mathcal{G}_{\mathrm{A}} \cap \mathcal{G}_{\mathrm{B}}=\{+\mathbb{I},-\mathbb{I}\}$.

Relativity of simultaneity singles out the associative division algebras.

## Overview

1. Motivations: QG and device-independent QIT
2. Relativity of simultaneity and the qubit
3. Randomness generation via rotational symmetry
4. Conclusions

## Overview

1. Motivations: QG and device-independent QIT
2. Relativity of simultaneity and the qubit
3. Randomness generation via rotational symmetry
4. Conclusions

## Randomness generation: quantum analysis



## Randomness generation: quantum analysis



## Randomness generation: quantum analysis



If input is $x=1$ : do nothing to preparation device; if $x=2$ : rotate it (relative to measurement device) by angle $\alpha$.

## Randomness generation: quantum analysis



If input is $x=1$ : do nothing to preparation device; if $x=2$ : rotate it (relative to measurement device) by angle $\alpha$.

SDI assumption: "spin" of system $\leq \mathbf{J}$
No further assumptions on devices / system.

## Randomness generation: quantum analysis



If input is $x=1$ : do nothing to preparation device; if $x=2$ : rotate it (relative to measurement device) by angle $\alpha$.

SDI assumption: "spin" of system $\leq \mathbf{J}$
No further assumptions on devices / system.
Rotation described by (projective) unitary representation of SO(2):

$$
U_{\alpha}=\bigoplus_{j=-J}^{J} n_{j} e^{i j \alpha}, \quad P(b \mid \alpha)=\operatorname{tr}\left(M_{b} U_{\alpha} \rho U_{\alpha}^{\dagger}\right)
$$

## Randomness generation: quantum analysis



## Randomness generation: quantum analysis

fixed $\alpha$


$$
E_{x}=P(+1 \mid x)-P(-1 \mid x)
$$



## Randomness generation: quantum analysis

$$
E_{x}=P(+1 \mid x)-P(-1 \mid x)
$$

fixed $\alpha$


$$
b \in\{ \pm 1\} \downarrow
$$

## - "Boring" deterministic correlations: outcome $b$ independent of $x$



## Randomness generation: quantum analysis



- "Boring" deterministic correlations: outcome $b$ independent of $x$
- "Interesting" deterministic correlations: outcome $b$ is a function of $x$



## Randomness generation: quantum analysis



- "Boring" deterministic correlations: outcome $b$ independent of $x$
- "Interesting" deterministic correlations: outcome $b$ is a function of $x$

Suppose ( $E_{1}, E_{2}$ ) observed. Looks random. But:


## Randomness generation: quantum analysis



- "Boring" deterministic correlations: outcome $b$ independent of $x$
- "Interesting" deterministic correlations: outcome $b$ is a function of $x$

Suppose ( $E_{1}, E_{2}$ ) observed. Looks random. But:

$$
\left(E_{1}, E_{2}\right)=\sum_{\lambda} p(\lambda)\left(E_{1}^{(\lambda)}, E_{2}^{(\lambda)}\right)_{\operatorname{det}}
$$



## Randomness generation: quantum analysis

$$
E_{x}=P(+1 \mid x)-P(-1 \mid x)
$$

fixed $\alpha$

Can predict outcome!

- "Boring" deterministic correlations: outcome $b$ independent of $x$
- "Interesting" deterministic correlations: outcome $b$ is a function of $x$

Suppose ( $E_{1}, E_{2}$ ) observed. Looks random. But:

$$
\left(E_{1}, E_{2}\right)=\sum_{\lambda} p(\lambda)\left(E_{1}^{(\lambda)}, E_{2}^{(\lambda)}\right)_{\operatorname{det}}
$$



## Randomness generation: quantum analysis



Which correlations are possible?

## Randomness generation: quantum analysis



Which correlations are possible? Theorem: exactly those:
$\frac{1}{2}\left(\sqrt{1+E_{1}} \sqrt{1+E_{2}}+\sqrt{1-E_{1}} \sqrt{1-E_{2}}\right) \geq\left\{\begin{array}{cl}\cos (J \alpha) & \text { if }|J \alpha|<\frac{\pi}{2} \\ 0 & \text { if }|J \alpha| \geq \frac{\pi}{2}\end{array}\right.$
C. L. Jones, S. L. Ludescher, A. Aloy, MM, arXiv:2210.14811 using results of
T. Van Himbeeck, E. Woodhead, N. J. Cerf, R. García-Patrón, S. Pironio, Quantum 1, 33 (2017).

## Randomness generation: quantum analysis



Which correlations are possible? Theorem: exactly those:

$$
\frac{1}{2}\left(\sqrt{1+E_{1}} \sqrt{1+E_{2}}+\sqrt{1-E_{1}} \sqrt{1-E_{2}}\right) \geq\left\{\begin{array}{cl}
\cos (J \alpha) & \text { if }|J \alpha|<\frac{\pi}{2} \\
0 & \text { if }|J \alpha| \geq \frac{\pi}{2}
\end{array}\right.
$$



Angle $\alpha \geq \pi /(2 J)$ :
no certifiable randomness.

## Randomness generation: quantum analysis



Which correlations are possible? Theorem: exactly those:
$\frac{1}{2}\left(\sqrt{1+E_{1}} \sqrt{1+E_{2}}+\sqrt{1-E_{1}} \sqrt{1-E_{2}}\right) \geq\left\{\begin{array}{c}\cos (J \alpha) \\ \text { if }|J \alpha|<\frac{\pi}{2} \\ 0 \\ \text { if }|J \alpha| \geq \frac{\pi}{2}\end{array}\right.$


Blue curved set of correlations.
If observed correlation away from red line: certifiable private randomness.

## Quantum theory is actually not needed



## Quantum theory is actually not needed



- Can we formulate our SDI assumption without quantum terminology?
- Can we use the protocol to certify random numbers without QT?


## Quantum theory is actually not needed



- Can we formulate our SDI assumption without quantum terminology?
- Can we use the protocol to certify random numbers without QT?
- Can we understand the curved boundary of correlations from spatial symmetry alone, without assuming QT?



## Quantum theory is actually not needed



- Can we formulate our SDI assumption without quantum terminology?
- Can we use the protocol to certify random numbers without QT?
- Can we understand the curved boundary of correlations from spatial symmetry alone, without assuming QT?

Yes we can!


## Rotation boxes beyond quantum theory

- Definition of quantum spin-J boxes:


## Rotation boxes beyond quantum theory

- Definition of quantum spin-J boxes: $\mathcal{Q}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha) \mid p(b \mid \alpha)=\operatorname{tr}\left(E_{b} U_{\alpha} \rho U_{\alpha}^{\dagger}\right)\right\}$, $\left\{E_{b}\right\}$ some POVM, $\rho$ some density matrix, $U_{\alpha}=\bigoplus_{j=-J}^{J} n_{j} e^{i j \alpha}$, with arbitrary multiplicities $n_{j}$.


## Rotation boxes beyond quantum theory

- Definition of quantum spin-J boxes:
$\mathcal{Q}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha) \mid p(b \mid \alpha)=\operatorname{tr}\left(E_{b} U_{\alpha} \rho U_{\alpha}^{\dagger}\right)\right\}$,
$\left\{E_{b}\right\}$ some POVM, $\rho$ some density matrix,
$U_{\alpha}=\bigoplus_{j=-J}^{J} n_{j} e^{i j \alpha}$, with arbitrary multiplicities $n_{j}$.
Consequence: every $p$ is a trigonometric polynomial of degree $\mathbf{2 J}$

$$
\text { (e.g. } p(+\mid \alpha)=\frac{1}{2}+\frac{1}{2} \cos \alpha \quad \text { for } J=\frac{1}{2} \text { ). }
$$

- Definition of quantum spin-J boxes:
$\mathcal{Q}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha) \mid p(b \mid \alpha)=\operatorname{tr}\left(E_{b} U_{\alpha} \rho U_{\alpha}^{\dagger}\right)\right\}$, $\left\{E_{b}\right\}$ some POVM, $\rho$ some density matrix,
$U_{\alpha}=\bigoplus_{j=-J}^{J} n_{j} e^{i j \alpha}$, with arbitrary multiplicities $n_{j}$.
Consequence: every $p$ is a trigonometric polynomial of degree $2 \mathbf{J}$

$$
\text { (e.g. } p(+\mid \alpha)=\frac{1}{2}+\frac{1}{2} \cos \alpha \quad \text { for } J=\frac{1}{2} \text { ). }
$$

- Definition of (general) spin-J rotation boxes:

$$
\mathcal{R}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha)=c_{0}+\sum_{j=1}^{2 J} c_{j} \cos (j \alpha)+s_{j} \sin (j \alpha)\right\}
$$

$$
0 \leq p(+1 \mid \alpha) \leq 1 \quad \text { for all } \alpha
$$

## Rotation boxes beyond quantum theory

- Definition of quantum spin-J boxes:

$$
\mathcal{Q}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha) \mid p(b \mid \alpha)=\operatorname{tr}\left(E_{b} U_{\alpha} \rho U_{\alpha}^{\dagger}\right)\right\},
$$

- Definition of (general) spin-J rotation boxes:

$$
\mathcal{R}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha)=c_{0}+\sum_{j=1}^{2 J} c_{j} \cos (j \alpha)+s_{j} \sin (j \alpha)\right\}
$$

## Rotation boxes beyond quantum theory

- Definition of quantum spin-J boxes:

$$
\mathcal{Q}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha) \mid p(b \mid \alpha)=\operatorname{tr}\left(E_{b} U_{\alpha} \rho U_{\alpha}^{\dagger}\right)\right\},
$$

- Definition of (general) spin-J rotation boxes:

$$
\mathcal{R}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha)=c_{0}+\sum_{j=1}^{2 J} c_{j} \cos (j \alpha)+s_{j} \sin (j \alpha)\right\}
$$

## Rotation boxes beyond quantum theory

- Definition of quantum spin-J boxes:

$$
\mathcal{Q}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha) \mid p(b \mid \alpha)=\operatorname{tr}\left(E_{b} U_{\alpha} \rho U_{\alpha}^{\dagger}\right)\right\},
$$

- Definition of (general) spin-J rotation boxes:

$$
\mathcal{R}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha)=c_{0}+\sum_{j=1}^{2 J} c_{j} \cos (j \alpha)+s_{j} \sin (j \alpha)\right\}
$$

Clearly $\mathcal{Q}_{J} \subseteq \mathcal{R}_{J}$.

- Definition of quantum spin-J boxes:

$$
\mathcal{Q}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha) \mid p(b \mid \alpha)=\operatorname{tr}\left(E_{b} U_{\alpha} \rho U_{\alpha}^{\dagger}\right)\right\},
$$

- Definition of (general) spin-J rotation boxes:

$$
\mathcal{R}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha)=c_{0}+\sum_{j=1}^{2 J} c_{j} \cos (j \alpha)+s_{j} \sin (j \alpha)\right\}
$$

Clearly $\mathcal{Q}_{J} \subseteq \mathcal{R}_{J}$.
It can be shown directly that $\mathcal{Q}_{0}=\mathcal{R}_{0}$ and $\mathcal{Q}_{1 / 2}=\mathcal{R}_{1 / 2}$. However, for some larger J, we have $\mathcal{Q}_{J} \subsetneq \mathcal{R}_{J}$, details here:
A. Aloy, T. Galley, C. L. Jones, S. L. Ludescher, MM, upcoming (2023).

## Rotation boxes beyond quantum theory

- Definition of quantum spin-J boxes:

$$
\mathcal{Q}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha) \mid p(b \mid \alpha)=\operatorname{tr}\left(E_{b} U_{\alpha} \rho U_{\alpha}^{\dagger}\right)\right\},
$$

- Definition of (general) spin-J rotation boxes:

$$
\mathcal{R}_{J}:=\left\{\alpha \mapsto p(+1 \mid \alpha)=c_{0}+\sum_{j=1}^{2 J} c_{j} \cos (j \alpha)+s_{j} \sin (j \alpha)\right\}
$$

Clearly $\mathcal{Q}_{J} \subseteq \mathcal{R}_{J}$.
It can be shown directly that $\mathcal{Q}_{0}=\mathcal{R}_{0}$ and $\mathcal{Q}_{1 / 2}=\mathcal{R}_{1 / 2}$. However, for some larger J, we have $\mathcal{Q}_{J} \subsetneq \mathcal{R}_{J}$, details here:
A. Aloy, T. Galley, C. L. Jones, S. L. Ludescher, MM, upcoming (2023).
$\mathcal{R}_{J}$ from rep. of $\mathrm{SO}(2)$ on (non-quantum) "orbitope" state spaces

## Boxes for only two input angles

$$
\frac{1}{2}\left(\sqrt{1+E_{1}} \sqrt{1+E_{2}}+\sqrt{1-E_{1}} \sqrt{1-E_{2}}\right) \geq\left\{\begin{array}{cc}
\cos (J \alpha) & \text { if }|J \alpha|<\frac{\pi}{2} \\
0 & \text { if }|J \alpha| \geq \frac{\pi}{2}
\end{array}\right.
$$


$\mathcal{Q}_{J, \alpha}$ quantum correlations (for $\mathbf{2}$ angles)

## Boxes for only two input angles

$$
\frac{1}{2}\left(\sqrt{1+E_{1}} \sqrt{1+E_{2}}+\sqrt{1-E_{1}} \sqrt{1-E_{2}}\right) \geq\left\{\begin{array}{cl}
\cos (J \alpha) & \text { if }|J \alpha|<\frac{\pi}{2} \\
0 & \text { if }|J \alpha| \geq \frac{\pi}{2}
\end{array}\right.
$$


$\mathcal{Q}_{J, \alpha}$ quantum correlations (for $\mathbf{2}$ angles)
$\mathcal{R}_{J, \alpha}$ : rotation box correlations (for $\mathbf{2}$ angles)

## Boxes for only two input angles

$$
\frac{1}{2}\left(\sqrt{1+E_{1}} \sqrt{1+E_{2}}+\sqrt{1-E_{1}} \sqrt{1-E_{2}}\right) \geq\left\{\begin{array}{cl}
\cos (J \alpha) & \text { if }|J \alpha|<\frac{\pi}{2} \\
0 & \text { if }|J \alpha| \geq \frac{\pi}{2}
\end{array}\right.
$$



## Boxes for only two input angles

$$
\frac{1}{2}\left(\sqrt{1+E_{1}} \sqrt{1+E_{2}}+\sqrt{1-E_{1}} \sqrt{1-E_{2}}\right) \geq\left\{\begin{array}{cc}
\cos (J \alpha) & \text { if }|J \alpha|<\frac{\pi}{2} \\
0 & \text { if }|J \alpha| \geq \frac{\pi}{2}
\end{array}\right.
$$



Theorem: $\mathcal{Q}_{J, \alpha}=\mathcal{R}_{J, \alpha}$. C. L. Jones, S. L. Ludescher, A. Aloy, MM, arXiv:2210.14811

## Boxes for only two input angles

$$
\frac{1}{2}\left(\sqrt{1+E_{1}} \sqrt{1+E_{2}}+\sqrt{1-E_{1}} \sqrt{1-E_{2}}\right) \geq\left\{\begin{array}{cc}
\cos (J \alpha) & \text { if }|J \alpha|<\frac{\pi}{2} \\
0 & \text { if }|J \alpha| \geq \frac{\pi}{2}
\end{array}\right.
$$



Theorem: $\mathcal{Q}_{J, \alpha}=\mathcal{R}_{J, \alpha}$. C. L. Jones, S. L. Ludescher, A. Aloy, MM, arXiv:2210.14811
Can derive set of quantum correlations without assuming QT.

## Boxes for only two input angles

$$
\frac{1}{2}\left(\sqrt{1+E_{1}} \sqrt{1+E_{2}}+\sqrt{1-E_{1}} \sqrt{1-E_{2}}\right) \geq\left\{\begin{array}{cc}
\cos (J \alpha) & \text { if }|J \alpha|<\frac{\pi}{2} \\
0 & \text { if }|J \alpha| \geq \frac{\pi}{2}
\end{array}\right.
$$


$\mathcal{Q}_{J, \alpha}$ quantum correlations (for $\mathbf{2}$ angles)
$\mathcal{R}_{J, \alpha}$ : rotation box correlations (for $\mathbf{2}$ angles)

Clearly $\quad \mathcal{Q}_{J, \alpha} \subseteq \mathcal{R}_{J, \alpha}$.

Theorem: $\mathcal{Q}_{J, \alpha}=\mathcal{R}_{J, \alpha}$. C. L. Jones, S. L. Ludescher, A. Aloy, MM, arXiv:2210.14811
Can derive set of quantum correlations without assuming QT.
Even eavesdropper with classical side information about beyond-quantum physics cannot predict the outcomes.


## Overview

1. Motivations: QG and device-independent QIT
2. Relativity of simultaneity and the qubit
3. Randomness generation via rotational symmetry
4. Conclusions

## Overview

1. Motivations: QG and device-independent QIT
2. Relativity of simultaneity and the qubit
3. Randomness generation via rotational symmetry
4. Conclusions

- Modest approach complementing direct QG approaches: study the constraints of spacetime on QT in simple scenarios.
- Relativity of simultaneity constrains the dimensionality of the qubit.
- Rotational symmetry determines the set of quantum correlations and the security of a SDI randomness generation protocol.
- Goal: theory-agnostic analysis of experiments in space and time.
- Speculation: is this (weak) evidence that QT might be modified in other regimes of space and time?

qubit: arXiv:1412.7112
randomness: arXiv:2210.14811

