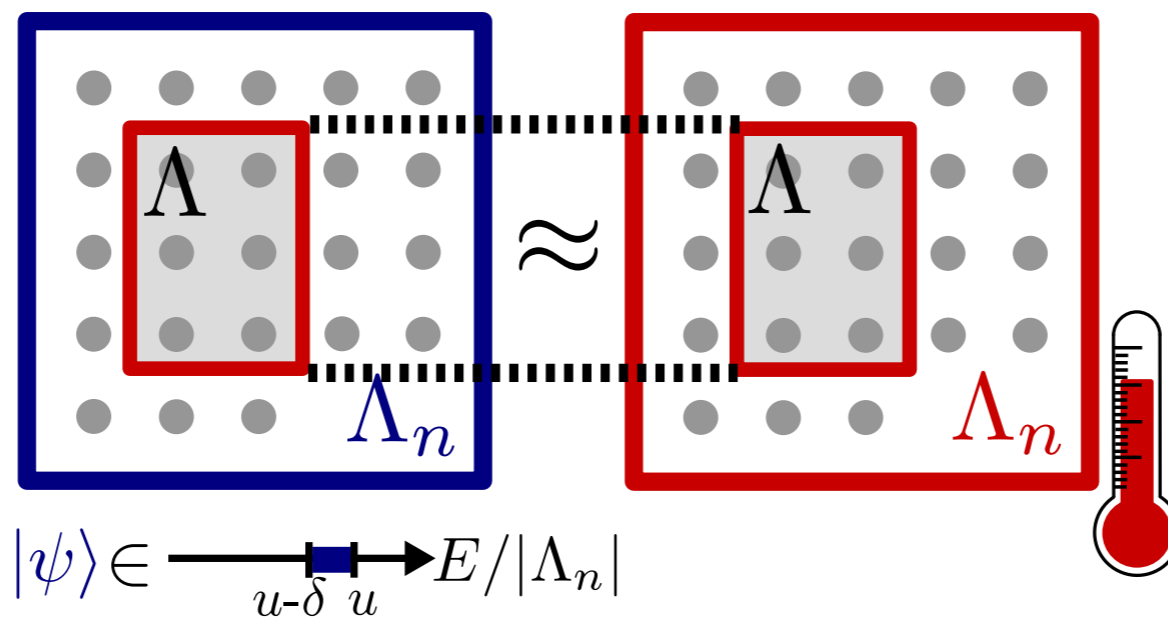


# Thermalization and canonical typicality in translation-invariant quantum lattice systems

Markus P. Müller

Institute für Theoretische Physik, Universität Heidelberg

$$\text{Tr}_{\Lambda_n \setminus \Lambda} |\psi\rangle\langle\psi| \approx \text{Tr}_{\Lambda_n \setminus \Lambda} \frac{\exp(-\beta H_{\Lambda_n}^p)}{Z}$$



joint work with Emily Adlam, Lluís Masanes, and Nathan Wiebe.

# Outline

## 1. The setup

Quantum lattice systems with interaction.

2. Equivalence  
of ensembles

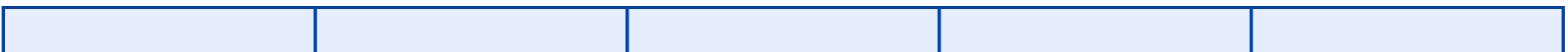
```
graph TD; A(2. Equivalence of ensembles) --> B(3. Canonical typicality); A --> C(4. Dynamical thermalization);
```

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4. Dynamical  
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## 5. Eigenstate thermalization hypothesis

A weak version that holds for all models.



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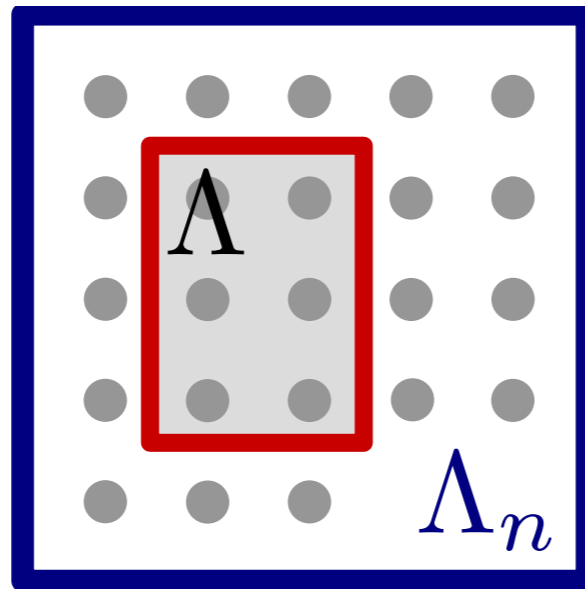
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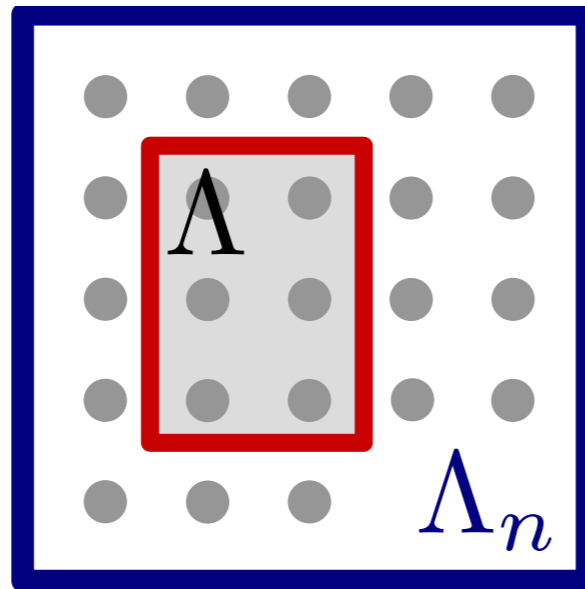
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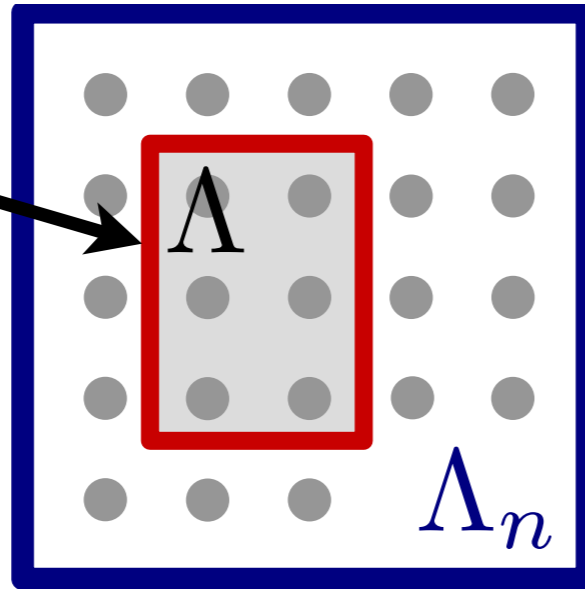


For example  $\Lambda_n = [-n, n]^\nu$ .



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Small sub-  
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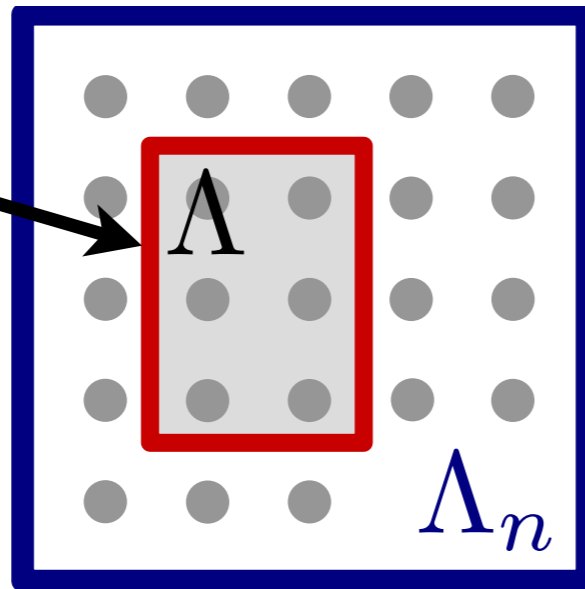


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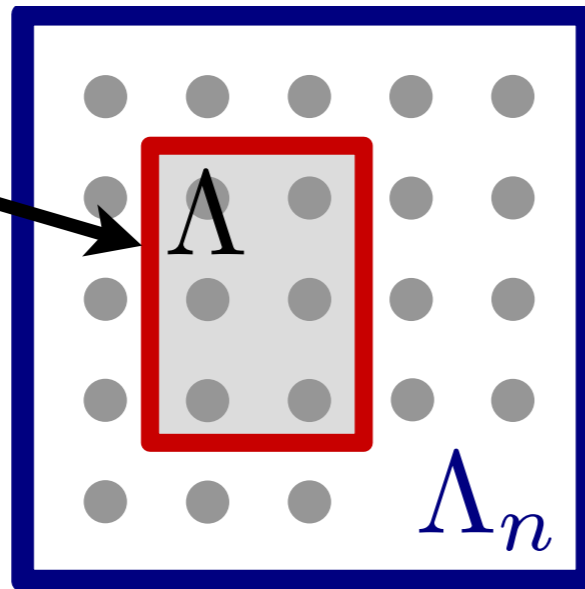
$$H_{\Lambda_n} = -J \sum_{i=1}^n (X_i X_{i+1} + Y_i Y_{i+1}) - h \sum_{i=1}^n Z_i.$$

Finite-range, translation-invariant; otherwise arbitrary.



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$$H_{\Lambda_n}^p = -J \sum_{i=1}^n (X_i X_{i+1} + Y_i Y_{i+1}) - h \sum_{i=1}^n Z_i - JX_n X_1 - JY_n Y_1.$$

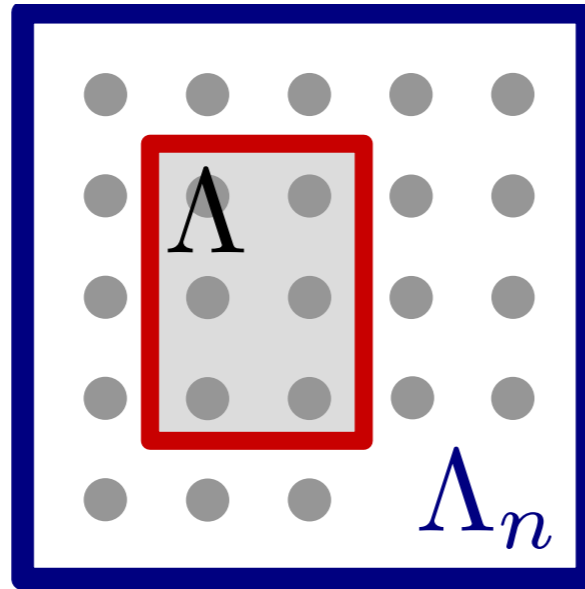
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In this talk: periodic boundary conditions.

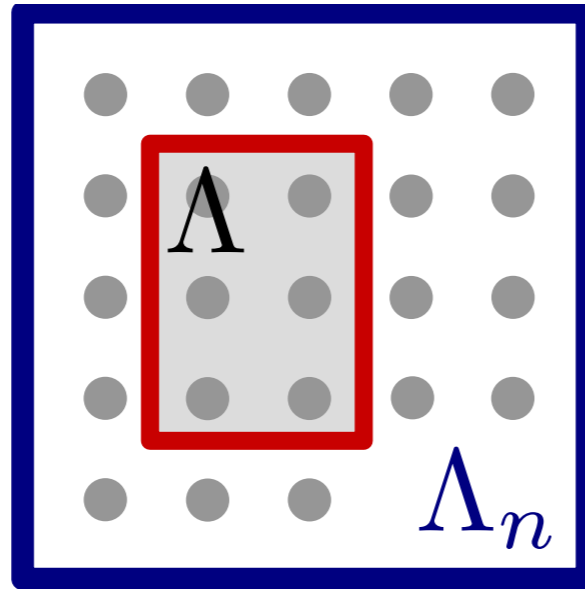




# Why translation-invariant lattice systems?



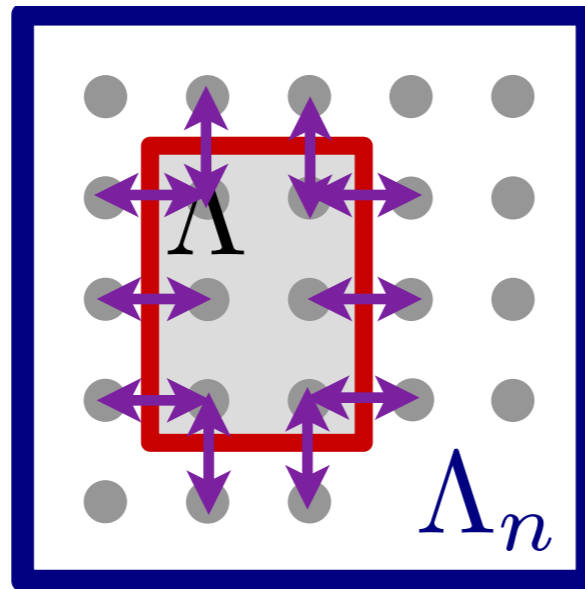
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- Goldstein et al.; Popescu et al.; Reimann; Short et al., ...:  
Subsystems of closed quantum systems equilibrate;  
but **equilibrium state is not in general thermal (Gibbs)!**



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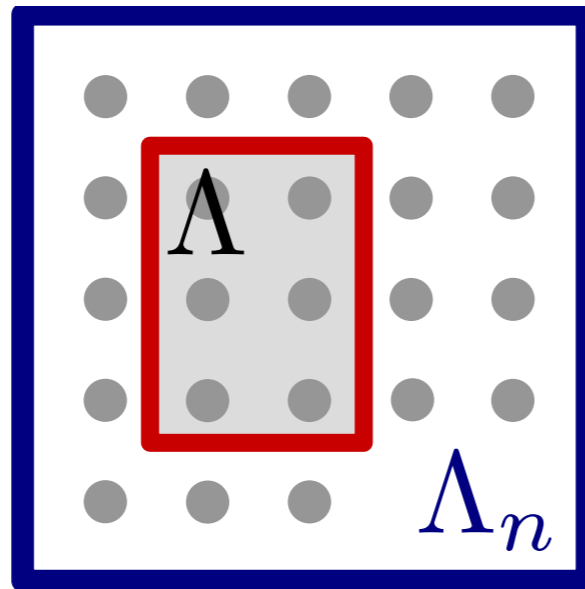


$$\|H_{\text{int}}\| \ll k_B T$$

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- We show: **T.I. + F.R.  $\Rightarrow$  thermality, also for strong interaction.**



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Every sequence of states with asymptotically minimal free energy density is equivalent to the canonical ensemble.





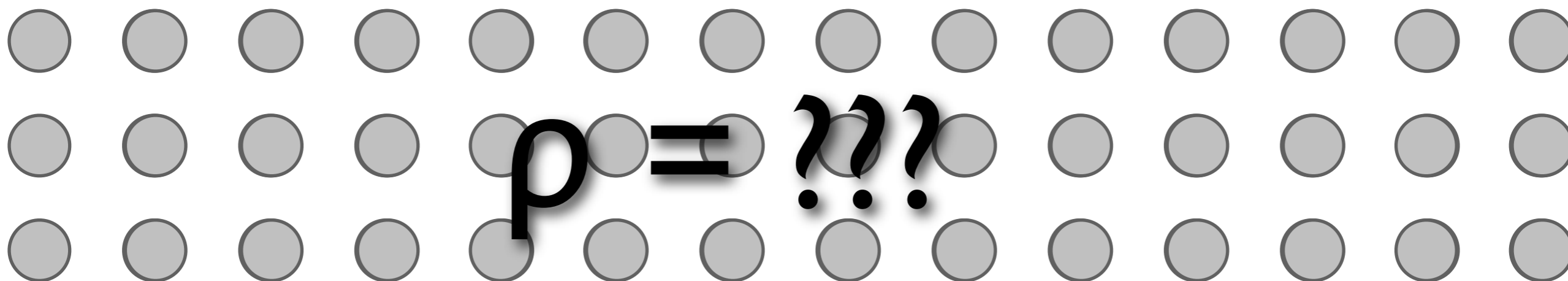
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Proof goes via infinite-lattice Gibbs states.



Definition: (cf. Barry Simon, Stat. Mech. of Lattice Gases)

A state  $\omega$  on the infinite lattice is a family of density matrices

$$(\omega_\Lambda)_{\Lambda \subset \mathbb{Z}^\nu \text{ finite}}$$

satisfying  $\Lambda' \subset \Lambda \Rightarrow \omega_{\Lambda'} = \text{Tr}_{\Lambda \setminus \Lambda'} \omega_\Lambda$ .



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Local Gibbs states  $\rho_{\Lambda_n}(\beta) := \exp(-\beta H_{\Lambda_n}^p) / Z$

minimize the free energy functional

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Use analogous definition on infinite lattice via densities

$$u(\omega) := \lim_{n \rightarrow \infty} \frac{1}{|\Lambda_n|} \text{tr}(\omega_{\Lambda_n} H_{\Lambda_n}^p),$$

$$s(\omega) := \lim_{n \rightarrow \infty} \frac{1}{|\Lambda_n|} S(\omega_{\Lambda_n}).$$



Definition: A translation-invariant state  $\omega$  on the infinite lattice is a **Gibbs state at inv.temp.  $\beta$**  if it minimizes the F.E. density, i.e.

$$u(\omega) - s(\omega)/\beta \leq f_{\text{th}}(\beta),$$

where  $f_{\text{th}}(\beta)$  is the limiting F.E. density of local Gibbs states.

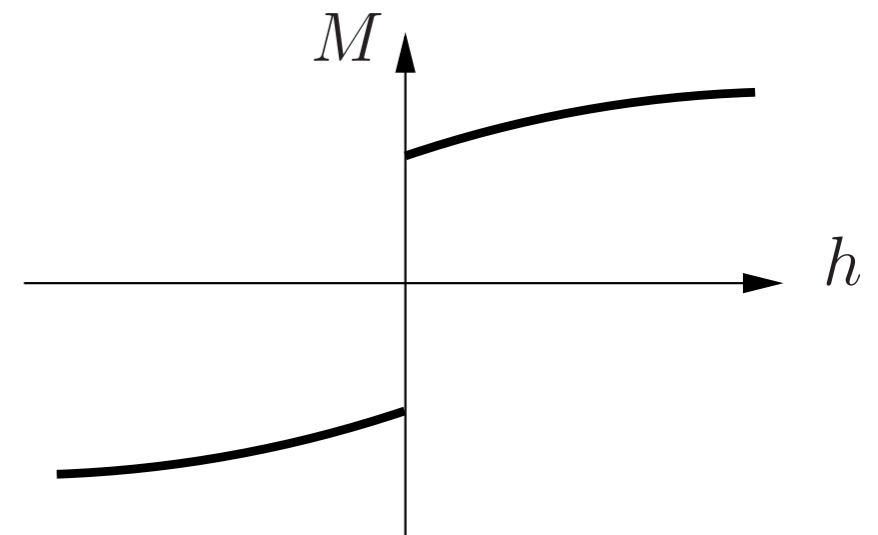


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- ! If  $\nu \geq 2$  then there can be **more than one Gibbs state**  
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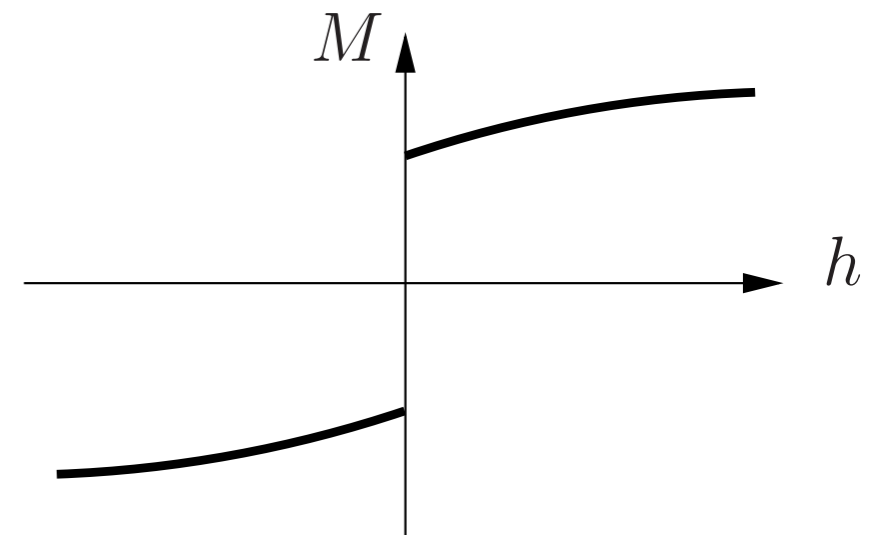


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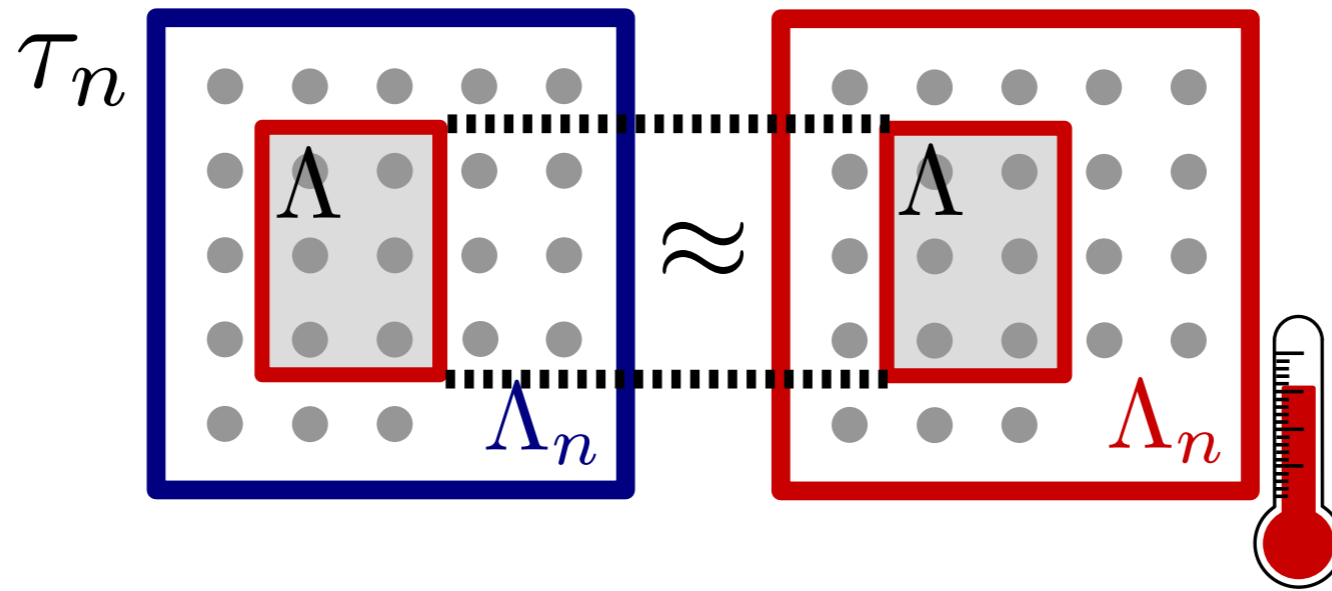
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Now show:

Every sequence of states with asymptotically minimal free energy density is equivalent to the canonical ensemble.

$$\text{Tr}_{\Lambda_n \setminus \Lambda} \mathcal{T}_n \approx \text{Tr}_{\Lambda_n \setminus \Lambda} \frac{\exp(-\beta H_{\Lambda_n}^p)}{Z}$$





$$\mathrm{Tr}_{\Lambda_n \setminus \Lambda} \tau_n \approx \mathrm{Tr}_{\Lambda_n \setminus \Lambda} \frac{\exp(-\beta H_{\Lambda_n}^p)}{Z}$$

**Theorem 2.** Suppose that  $(\tau_n)_{n \in \mathbb{N}}$  is any sequence of  $\Lambda_n$ -translation-invariant states on  $\Lambda_n$ , and  $\beta > 0$  such that there is a unique phase around inverse temperature  $\beta$ . If

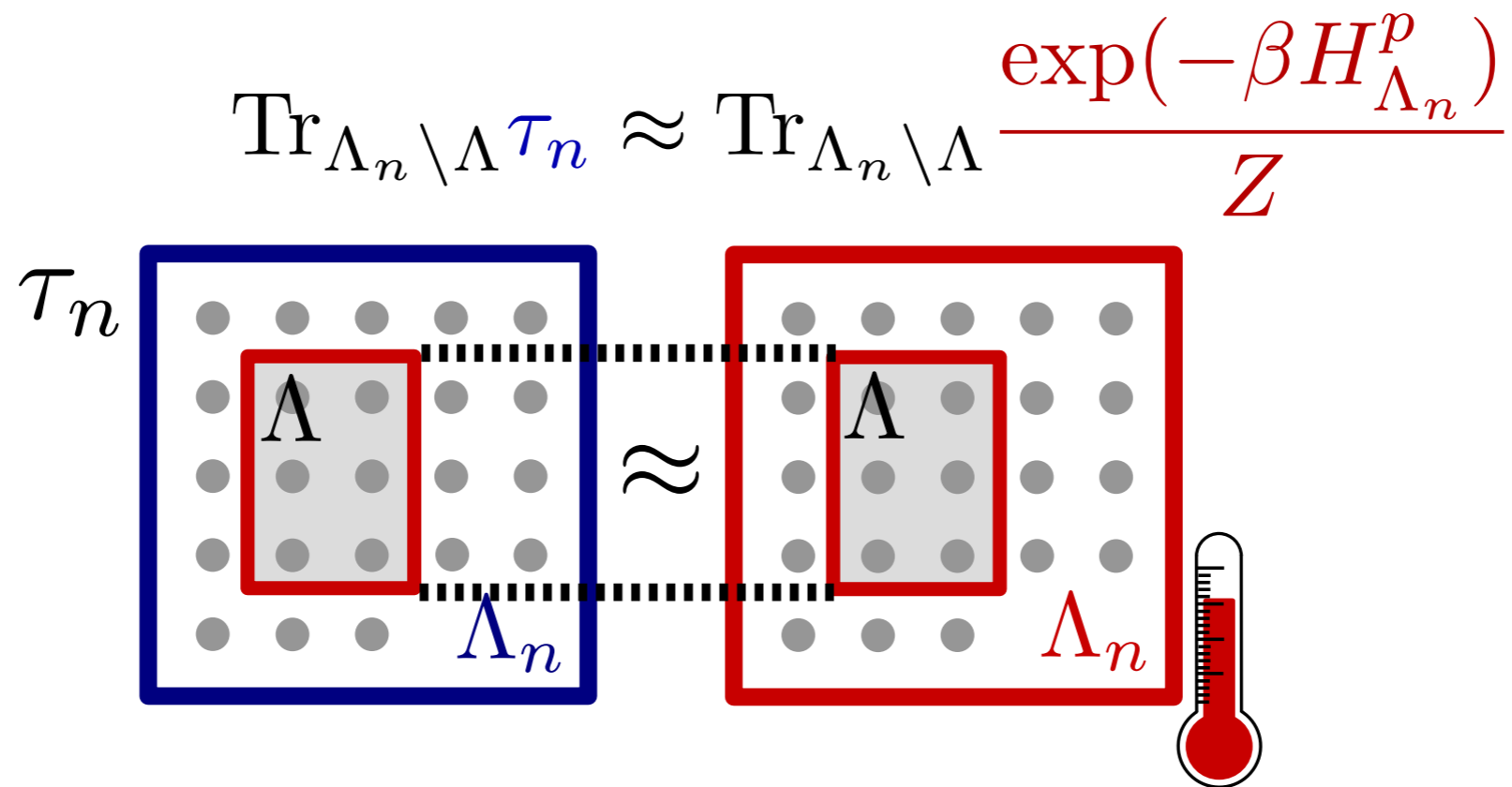
$$\limsup_{n \rightarrow \infty} \frac{1}{|\Lambda_n|} (\mathrm{tr}(\tau_n H_{\Lambda_n}^{BC}) - S(\tau_n)/\beta) \leq f_{\mathrm{th}}(\beta) \quad (2)$$

for some choice of boundary conditions  $BC$ , then

$$\lim_{n \rightarrow \infty} \left\| \mathrm{Tr}_{\Lambda_n \setminus \Lambda} \tau_n - \mathrm{Tr}_{\Lambda_n \setminus \Lambda} \frac{\exp(-\beta_n H_{\Lambda_n}^p)}{Z_n} \right\|_1 = 0, \quad (3)$$

where we may set  $\beta_n$  either equal to the fixed value  $\beta$ , or equal to the solution of  $\mathrm{tr}(H_{\Lambda_n}^p \rho_{\Lambda_n}^p(\beta_n))/|\Lambda_n| = u(\beta)$ .





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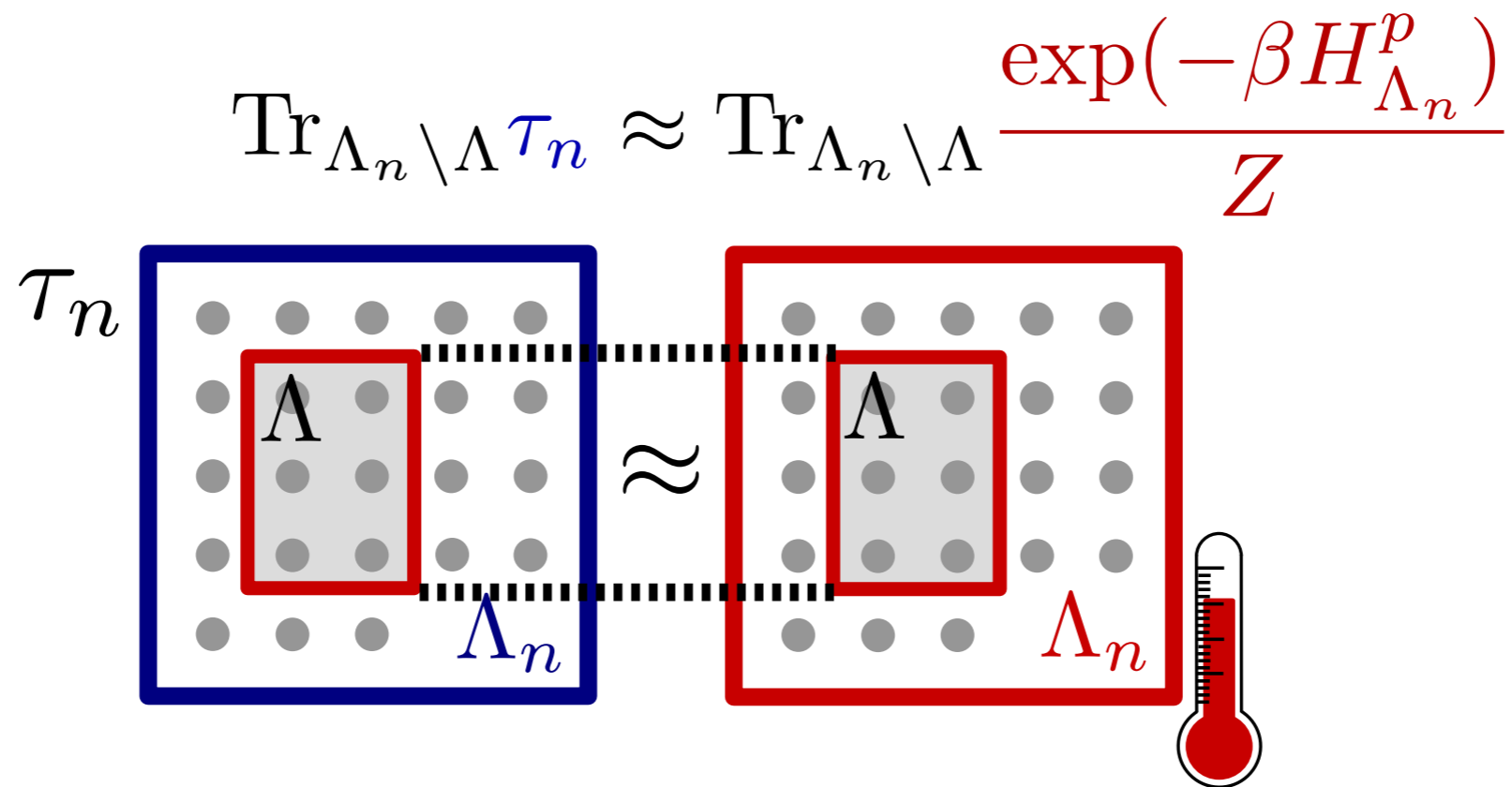
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**Example:**  $\tau_n =$  mixture on  $\mathrm{span} \{ |E\rangle \mid u - \delta \leq E/|\Lambda_n| \leq u \}$   
**Microcanonical ensemble!**





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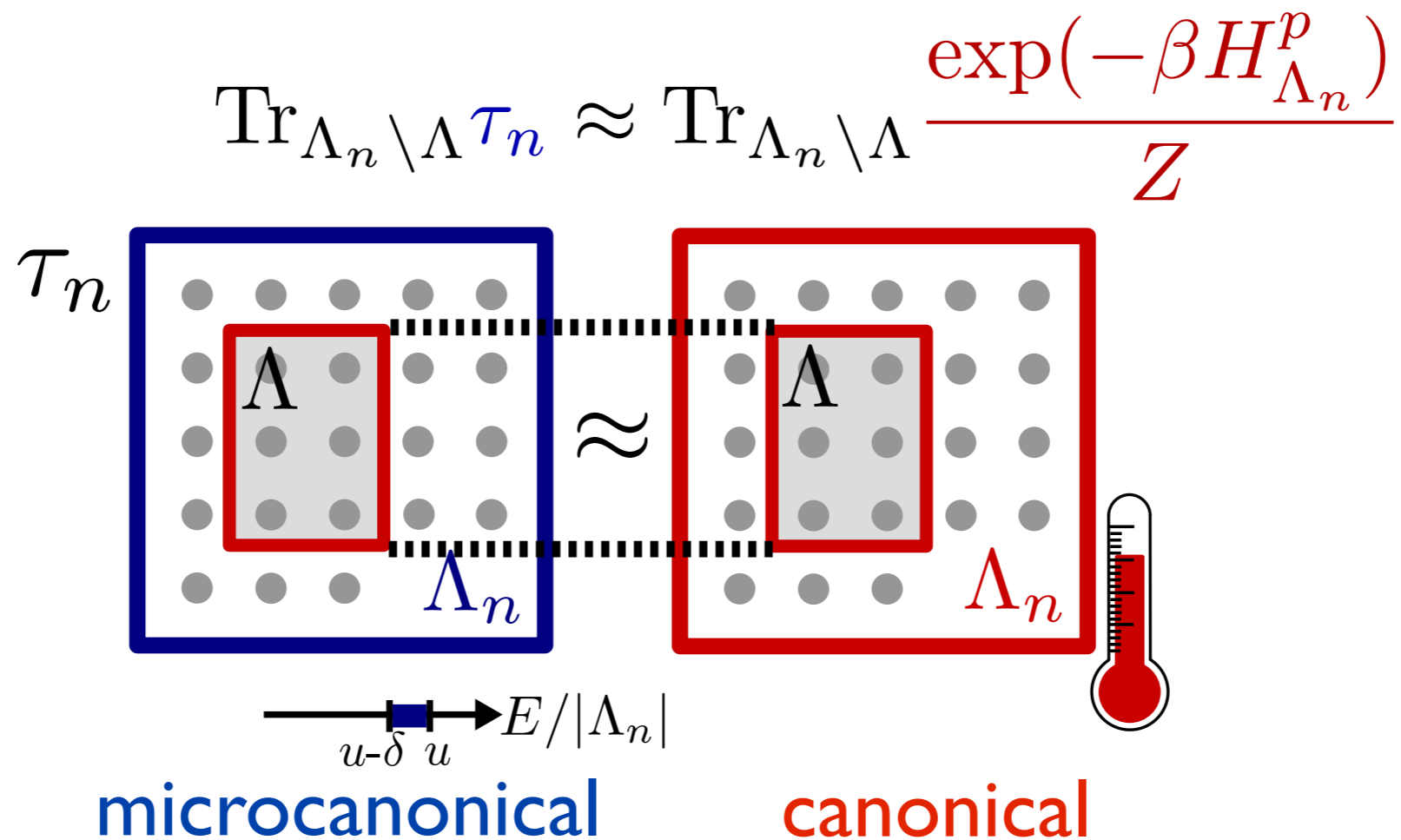
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**Proof:** they both converge locally to the unique infinite-volume Gibbs state.





For local qubits, without interaction,  $\delta=0$  (de Finetti Theorem):

$$\left\| \text{Tr}_{\Lambda_n \setminus \Lambda} \tau_n - \text{Tr}_{\Lambda_n \setminus \Lambda} \frac{\exp(-\beta H_{\Lambda_n}^p)}{Z} \right\|_1 \leq \frac{4|\Lambda|}{|\Lambda_n|}$$

Similar analytic bound for  $\delta>0$ , and numerically with interaction.



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A weak version that holds for all models.



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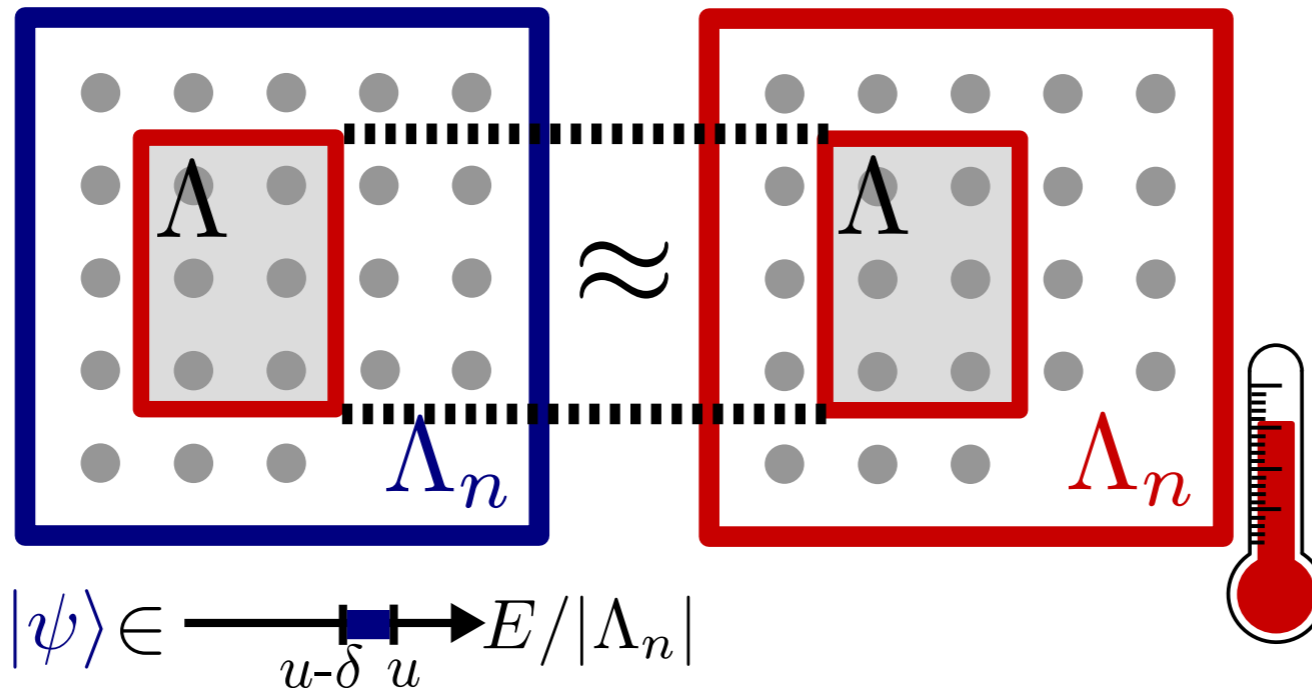
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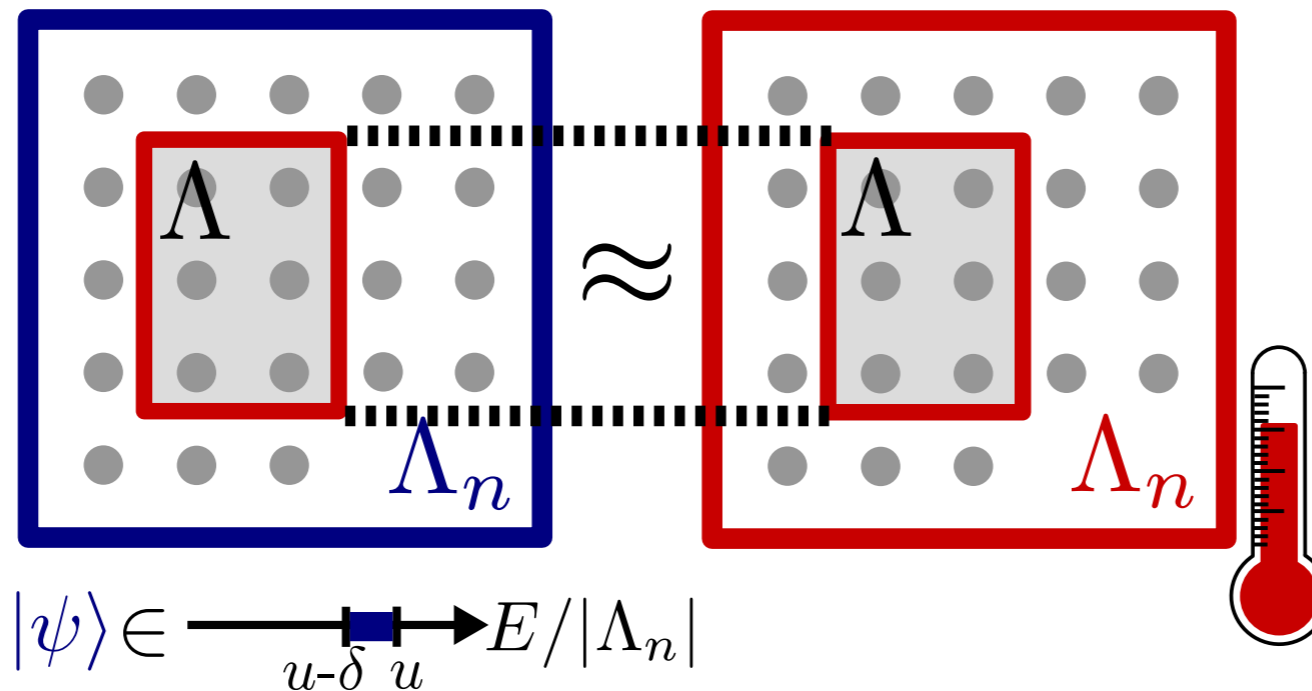
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Draw a pure state  $|\psi\rangle$   
at random from the  
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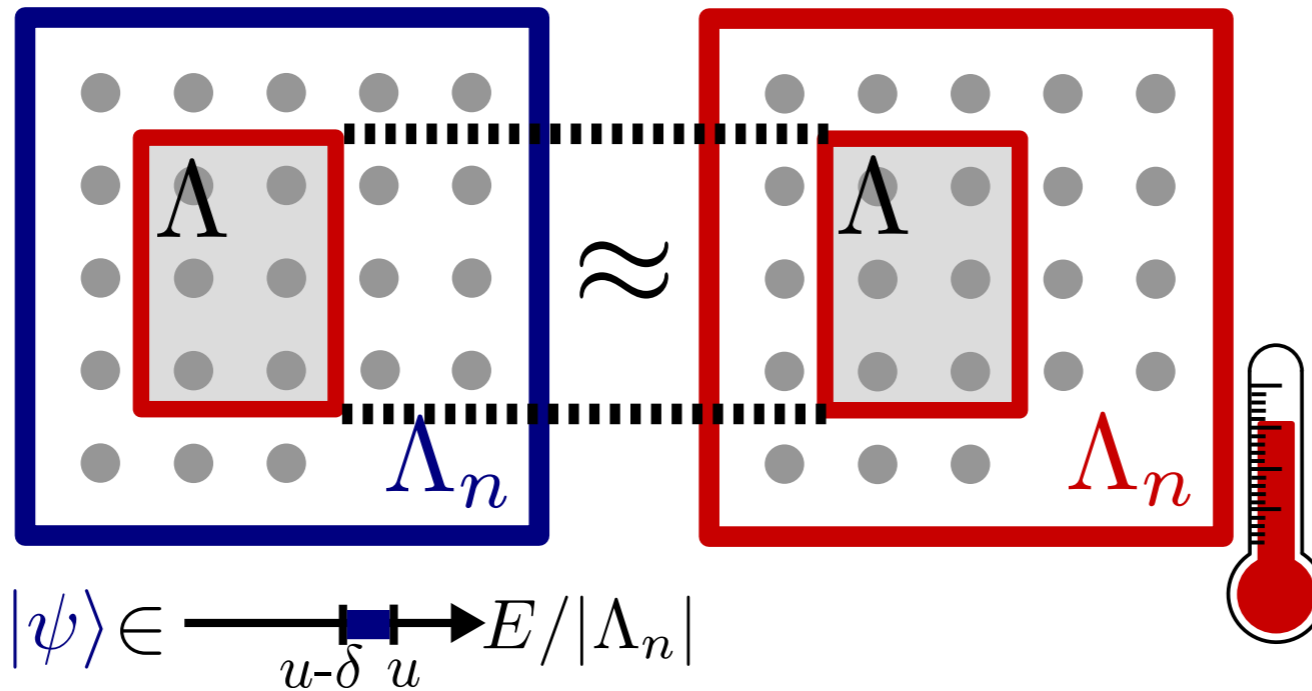
$$\text{span} \{ |E\rangle \mid u - \delta \leq E/|\Lambda_n| \leq u \}$$

If there is a unique  
phase around  $\beta(u)$  then:





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**Theorem 1.** For any  $\varepsilon \geq 0$ , the probability  $p$  that a state  $|\psi\rangle \in T_n^p$  sampled according to the unitarily invariant measure satisfies

$$\left\| \text{Tr}_{\Lambda_n \setminus \Lambda} |\psi\rangle\langle\psi| - \text{Tr}_{\Lambda_n \setminus \Lambda} \frac{\exp(-\beta H_{\Lambda_n}^p)}{Z} \right\|_1 \geq \varepsilon + \Delta_{n,\Lambda}$$

is doubly-exponentially small in the lattice size  $|\Lambda_n|$ ; that is,  $p \leq \exp(-\varepsilon^2 \exp(|\Lambda_n|s + o(|\Lambda_n|)))$ , where  $s = s(\omega_\beta)$  is the entropy density of the corresponding Gibbs state, and  $\Delta_{n,\Lambda}$  is a sequence of positive real numbers with  $\lim_{n \rightarrow \infty} \Delta_{n,\Lambda} = 0$  for every fixed  $\Lambda$ . Here,  $\beta$  can either be set equal to  $\beta(u)$  as defined above, or equal to the solution of  $\text{tr}(H_{\Lambda_n}^p \rho_{\Lambda_n}^p(\beta))/|\Lambda_n| = u$  (which depends on  $n$ ).



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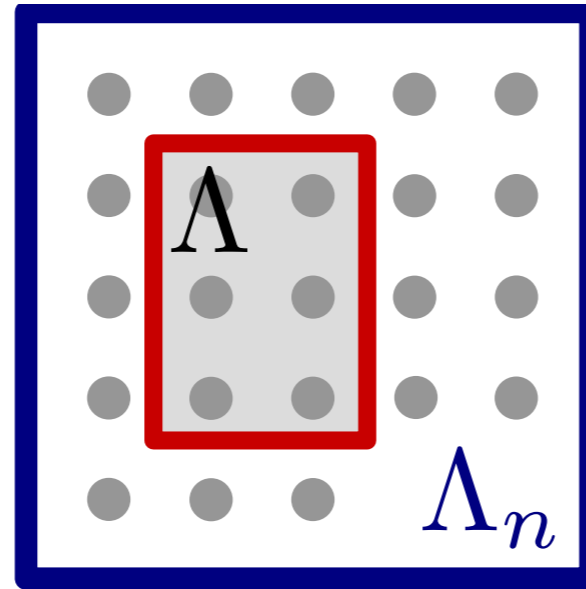
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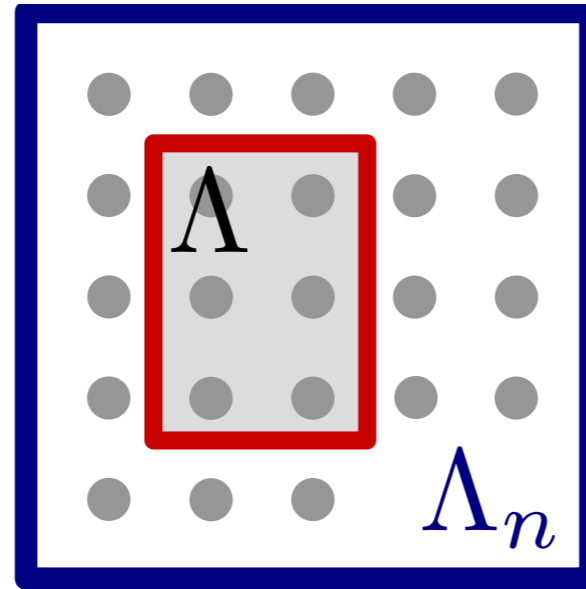
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Full lattice  $\Lambda_n$  evolves unitarily:  $|\psi(t)\rangle = \exp(-itH_{\Lambda_n}^p)|\psi(0)\rangle$



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We show: if Shannon entropy of initial occupation numbers

$$p_n = |\langle E_n | \psi(0) \rangle|^2$$

is close to maximal, in first order in  $|\Lambda_n|$ ,

then

$$\text{Tr}_{\Lambda_n \setminus \Lambda} |\psi(t)\rangle \langle \psi(t)| \approx \text{Tr}_{\Lambda_n \setminus \Lambda} \frac{\exp(-\beta H_{\Lambda_n}^p)}{Z}$$

for most times  $t$ .

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ETH (Deutsch '91, Srednicki '94):

Results could be significantly sharpened if

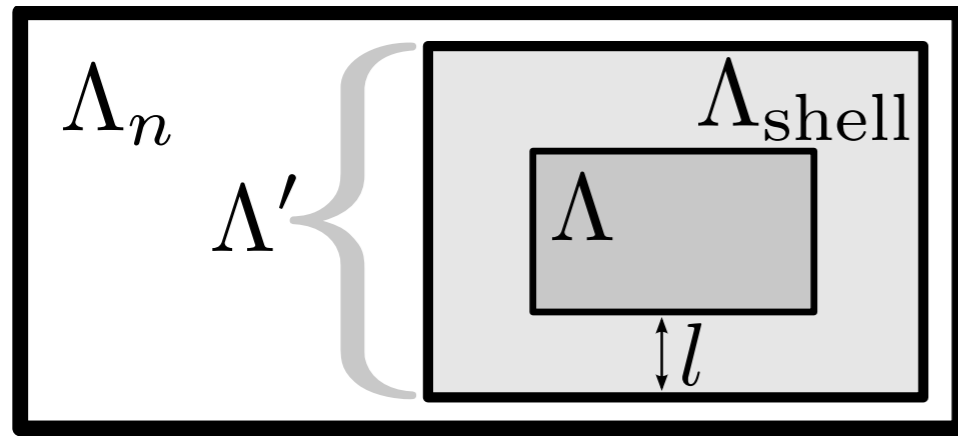
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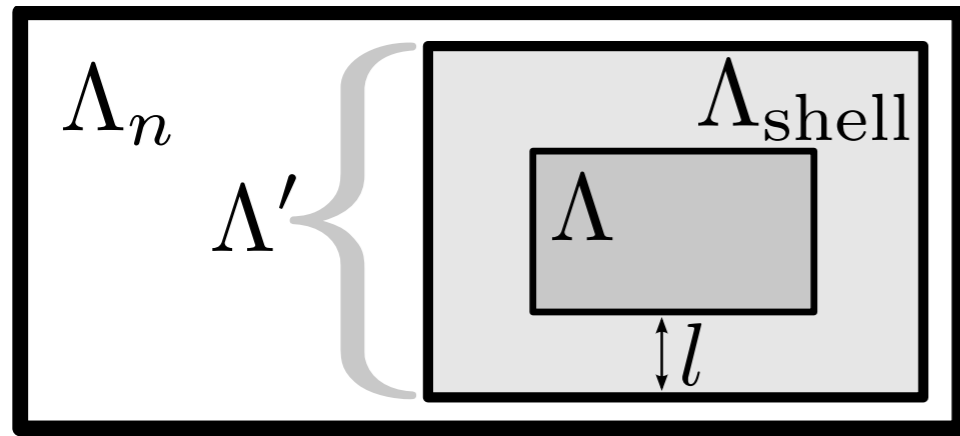
Natural conjecture: If  $|E\rangle$  is an energy eigenstate on  $\Lambda_n$ , then

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**Theorem 4.** *There is a state  $\omega_E$  on  $\Lambda'$  such that*

$$\|\text{Tr}_{\Lambda_{\text{shell}}}(\omega_E) - \text{Tr}_{\Lambda_n \setminus \Lambda} |E\rangle\langle E|\|_1 \leq \kappa \cdot e^{-c(l-r)/2}$$

*which is weakly diagonal in the eigenbasis  $\{|e\rangle\}$  of  $H_{\Lambda'}$ , i.e.*

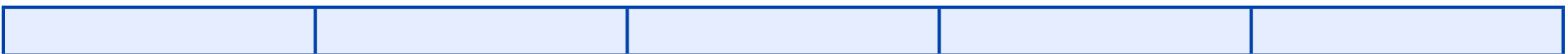
$$|\langle e_1 | \omega_E | e_2 \rangle| \leq e^{-(l-r)(e_1 - e_2)^2 / (8cv^2)},$$

*where  $\kappa, c, v > 0$  are constants, and  $r$  the range of interaction.*



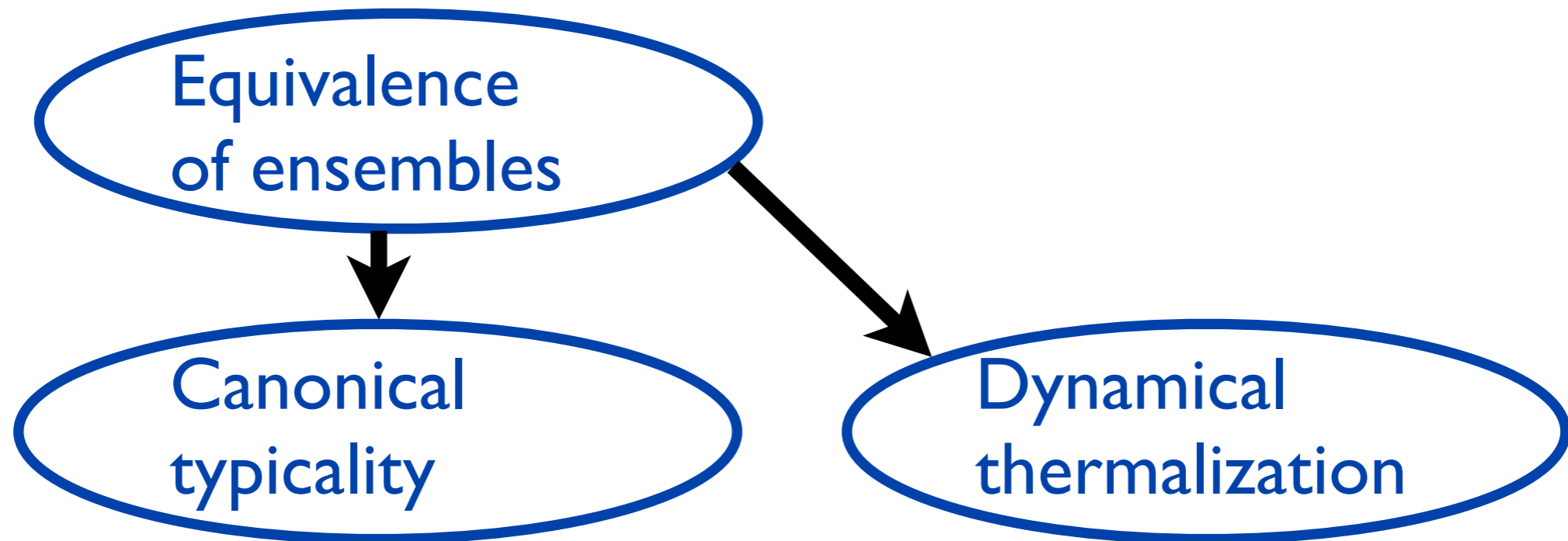
# Conclusions

The Gibbs state emerges naturally in translation-invariant quantum lattice systems.



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Open problem: Can we prove a strong version of the ETH? Integrability?

[arXiv:1312.7420](https://arxiv.org/abs/1312.7420)

