A resource-theoretic approach to thermodynamics

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Drawing by Lidia del Rio, http://www.itp.phys.ethz.ch/itp/itp/people/delriol

Outline

1. Motivation

What is a resource theory? Why and how thermo as a RT?



- 2. The resource theory of athermality Definition, results, surprises.
- 3. Extracting work from absence of correlations How to do more by knowing less...



A resource-theoretic approach to thermodynamics





1. Motivation

A resource-theoretic approach to thermodynamics





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A resource-theoretic approach to thermodynamics



Alice and Bob are very far apart, so

- they can only act locally on the particle pair,
- can only produce new states of the form $|\varphi_A\rangle|\psi_B\rangle$,
- can talk on the telephone.



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For example, can they transform $|\varphi\rangle \longrightarrow |\psi\rangle$?



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Free states / objects: $|\varphi_A\rangle|\psi_B\rangle$ (all other states are resources). Allowed transformations: local operations + classical communication



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Free states / objects: $|\varphi_A\rangle|\psi_B\rangle$ (all other states are resources). Allowed transformations: local operations + classical communication

Questions: can we transform $|\psi_1\rangle$ to $|\psi_2\rangle$? How many copies of $|\psi_1\rangle$ do we need to obtain *n* copies of



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Resource theory of cooking





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Free operations: cutting, mixing etc. Free objects: water, electricity,...?



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Resource theory of chemistry



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Resource theory of chemistry



Category-theoretic formulation: B. Coecke, T. Fritz, R. W. Spekkens, A mathematical theory of resources,

arXiv:1409.5531



1. Motivation

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Thought experiment: single particle in a box





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State of particle:

$$\left(\begin{array}{c} \operatorname{Prob}(L) \\ \operatorname{Prob}(R) \end{array}\right)$$



1. Motivation

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Thought experiment: single particle in a box



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 $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$ Know for sure the particle is left → can extract energy by expansion



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Thought experiment: single particle in a box



Erasing 1 bit of information costs work.

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This suggests that

- the objects of the theory are the observer's states of knowledge, i.e. probability distributions (on microstates),
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Something's missing: $\begin{pmatrix} p_1 \\ \vdots \\ p_N \end{pmatrix}$ and energies $\begin{pmatrix} E_1 \\ \vdots \\ E_N \end{pmatrix}$.



1. Motivation

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This suggests that

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the thermal states (Gibbs states)

$$\frac{1}{Z} \left(\begin{array}{c} e^{-E_1/(k_B T)} \\ \vdots \\ e^{-E_N/(k_B T)} \end{array} \right)$$



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Why?





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Why?



The inverse process is not dynamically impossible, but hard for us to implement.

→ Thermodynamics derives from limited knowledge about the physical system (thus, lack of control), together with energy conservation and microscopic reversibility.



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Why?



R.T.: maximal generality; clear mathematical "rules of the game"

Reproduces and refines results of standard non-equilibrium thermodynamics.



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2. Resource theory of athermality

2. Resource theory of athermality: precise definition

Remember:

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the thermal states (Gibbs states)

$$=\frac{1}{Z}\left(\begin{array}{c}e^{-E_1/(k_BT)}\\\vdots\\e^{-E_N/(k_BT)}\end{array}\right)$$

The resulting transformations are called thermal operations:



2. Resource theory of athermality

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 $\beta := 1/(k_B T)$ is now fixed.

Definition. A map Φ on a physical system S is a *thermal operation* if there is another (ancilla) system A with energy levels (E_1, \ldots, E_N) and thus Gibbs state $\gamma_A = (e^{-\beta E_1}, \ldots, e^{-\beta E_n})/Z$, and an energy-preserving permutation π on SA such that

 $\Phi(p_S) = (\pi[p_S \otimes \gamma_A])_S.$



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Slightly more general: catalytic thermal operation.



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Completely general: $p_S = \begin{pmatrix} 1/2 \\ 1/2 \end{pmatrix}$ could mean



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2. Resource theory of athermality





One can also allow quantum coherences, $\alpha |0\rangle + \beta |1\rangle$

→ (quantum) thermodynamics "at the nano scale"

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One can also allow quantum coherences, $\alpha |0\rangle + \beta |1\rangle$

→ (quantum) thermodynamics, not necessarily "at the nano scale"

2. Resource theory of athermality

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Work extraction: what is the largest possible W such that



within the rules of the resource theory?

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Work extraction: what is the largest possible W such that



by a catalytic thermal operation?

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2. Resource theory of athermality



by a catalytic thermal operation, if we allow a small probability $\varepsilon > 0$ of error?

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Work cost: what is the smallest possible W such that



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Theorem: The extractable work and work cost are

$$W_{\text{extr}} = k_B T \left(F_0^{\varepsilon}(p_S) - F(\gamma_S) \right),$$

$$W_{\text{cost}} = k_B T \left(F_{\infty}^{\varepsilon}(p_S) - F(\gamma_S) \right),$$

where F_{α} is the Rényi α -free energy:

$$F_{\alpha}(p_S) = k_B T \left(\frac{\operatorname{sgn} \alpha}{\alpha - 1} \log \sum_{i} p_i^{\alpha} \exp\left(\frac{-E_i(1 - \alpha)}{k_B T}\right) \right) - k_B T \log Z,$$

and
$$F_1(p_S) = F(p_S) = \langle E \rangle - k_B T S(p_S)$$

is the "standard" free energy.

M. Horodecki and J. Oppenheim, *Fundamental limitations for quantum and nanoscale thermodynamics*, Nature Communications **4**, 2059 (2013).

2. Resource theory of athermality

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Landauer's Principle: if $p_S = (1, 0)$ and two identical energies, $W_{\text{extr}} = W_{\text{cost}} = k_B T \ln 2.$

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Fundamental thermodynamical irreversibility!

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But: in the thermodynamic limit,
$$n \text{ independent copies of } p_S$$

$$\lim_{n \to \infty} \frac{1}{n} F_{\alpha}^{\varepsilon}(p_S^{\otimes n}) = F(p_S).$$

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2. Resource theory of athermality

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Work extraction and work cost *n* independent copies of p_S But: in the thermodynamic limit, $\lim_{n \to \infty} \frac{1}{n} F_{\alpha}^{\varepsilon}(p_S^{\otimes n}) = F(p_S).$ n independent single particles eat bath, temperature eat bath, temperature eat bath, temperature eat bath, temperature heat bath, temperature T 2. Resource theory of athermality

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ideal gas;

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Remember: nowhere have we actually used that we have a gas.



2. Resource theory of athermality

A resource-theoretic approach to thermodynamics

The second law(s)

Theorem: A transition $p_S \to p'_S$ is possible if and only if $F_{\alpha}(p_S) \ge F_{\alpha}(p'_S)$ for all $\alpha \ge 0$.

All a-free energies must go down!

F. Brandao, M. Horodecki, N. Ng, J. Oppenheim, and S. Wehner, *The second laws of quantum thermodynamics*, Proc. Natl. Acad. Sci. USA **112** (2015)

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Consequence: some states are incomparable, i.e. neither $p_S \rightarrow p'_S$ nor $p'_S \rightarrow p_S$.

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Again, in the thermodynamic limit, it all collapses to $F(p_S) \ge F(p'_S)$.

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Again, in the thermodynamic limit, it all collapses to $F(p_S) \ge F(p'_S)$. Constant $E \Rightarrow$ entropy cannot decrease.

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3. Work from absence of correlations

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with Michele Pastena (Heidelberg)



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In standard thermodynamics, correlations are costly:





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In standard thermodynamics, correlations are costly:



Comes from subadditivity of entropy:

 $S(p_{AB}) \le S(p_A) + S(p_B).$

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One can "extract work from correlations".

3. Work from absence of correlations

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Resource theory ("single-shot thermodynamics"): recall as below.



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3. Work from absence of correlations

A resource-theoretic approach to thermodynamics
Resource theory ("single-shot thermodynamics"): recall as below. Constrained by infinitely many "2nd laws": $F_{\alpha}(p_S) \ge F_{\alpha}(p'_S)$ for all $\alpha \ge 0$.

Intuition: it should be even more difficult to do this instead:



A resource-theoretic approach to thermodynamics

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But: **Theorem**: This process is possible if and only if $F(p_S) \ge F(p'_S)$.

MM and M. Pastena, arXiv:1409.3258

3. Work from absence of correlations

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Intuition: it should be even more difficult to do this instead:



But: **Theorem**: This process is possible if and only if $F(p_S) \ge F(p'_S)$.

It's actually easier - stochastic independence can be "burnt like a fuel"!

MM and M. Pastena, arXiv:1409.3258

3. Work from absence of correlations

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Didn't our intuition say the opposite??



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However, in the resource theory of athermality, there is more then just "the" free energy.

For all $a \neq 1$ there are correlations with $F_{\alpha}(p_{AB}) < F_{\alpha}(p_A) + F_{\alpha}(p_B).$

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However, in the resource theory of athermality, there is more then just "the" free energy.

For all $a \neq 1$ there are correlations with $F_{\alpha}(p_{AB}) < F_{\alpha}(p_A) + F_{\alpha}(p_B).$

Knowling less makes you "less *a*-confused" and allows you to do more.



3. Work from absence of correlations

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Conclusions

- There are good reasons for formulating thermodynamics as a resource theory.
- Very general approach; reproduces standard thermo results in the thermodynamic limit, but refines them.
- Produces sometimes very surprising results for "small" systems → does nature do that?



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