

A derivation of quantum theory from physical requirements

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Perimeter Institute for Theoretical Physics, Waterloo (Canada)



Joint work with Lluís Masanes
arXiv: 1004.1483



Outline

1. Motivation

Why??

2. General Probabilistic Theories

3. The Axioms

What do they mean?

4. Derivation of the Hilbert space formalism

Why are qubits 3D-balls??

5. What's beyond QT?

I. Motivation

John A. Wheeler, New York Times, Dec. 12 2000:

„Quantum physics [...] has explained the structure of atoms and molecules, [...] the behavior of semiconductors [...] and the comings and goings of particles from neutrinos to quarks.

*Successful, yes, but mysterious, too.
Why does the quantum exist?“*



The New York Times

I. Motivation

ANNALS OF PHYSICS **194**, 336–386 (1989)

Testing Quantum Mechanics

STEVEN WEINBERG*

*Theory Group, Department of Physics,
University of Texas, Austin, Texas 78712*

Received March 6, 1989

This paper presents a general framework for introducing nonlinear corrections into ordinary quantum mechanics, that can serve as a guide to experiments that would be sensitive to such corrections. In the class of generalized theories described here, the equations that determine the time-dependence of the wave function are no longer linear, but are of Hamiltonian type. Also, wave functions that differ by a constant factor represent the same physical state and satisfy the same time-dependence equations. As a result, there is no difficulty in combining separated subsystems. Prescriptions are given for determining the states in which observables have definite values and for calculating the expectation values of observables for general states, but the calculation of probabilities requires detailed analysis of the method of measurement. A study is presented of various possibilities



I. Motivation

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Volume 143, number 1,2

PHYSICS LETTERS A

1 January 1990

WEINBERG'S NON-LINEAR QUANTUM MECHANICS AND SUPRALUMINAL COMMUNICATIONS

N. GISIN

Group on Applied Physics, University of Geneva, 1211 Geneva 4, Switzerland

Received 16 October 1989; accepted for publication 3 November 1989

Communicated by J.P. Vigièr

We show with an example that Weinberg's general framework for introducing non-linear corrections into quantum mechanics allows for arbitrarily fast communications.

Recently Weinberg has proposed a general framework for introducing non-linear corrections into ordinary quantum mechanics [1,2]. Although we fully support his emphasis on the importance of testing quantum mechanics, we would like in this Letter to draw attention to the difficulty of modifying quantum mechanics without introducing arbitrarily fast actions at a distance. Below we show how to construct, within Weinberg's framework, an arbitrarily fast telephone line. In ordinary quantum mechanics

to know what such an apparatus is... do you know what is inside your phone?) In order to simplify we consider only a single-bit message. The two directions z and u are in the xz -plane orthogonal to the incoming flow of particles, and are 45° from each other. The way the inhomogeneous magnetic field acts on the particles is well-known from experimental evidence. After the apparatus there are two counters. For each particle one of the counters will click. This click will be amplified until all readers of



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Our results:

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- Methods to construct natural consistent **modifications of quantum theory.**

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Builds on:

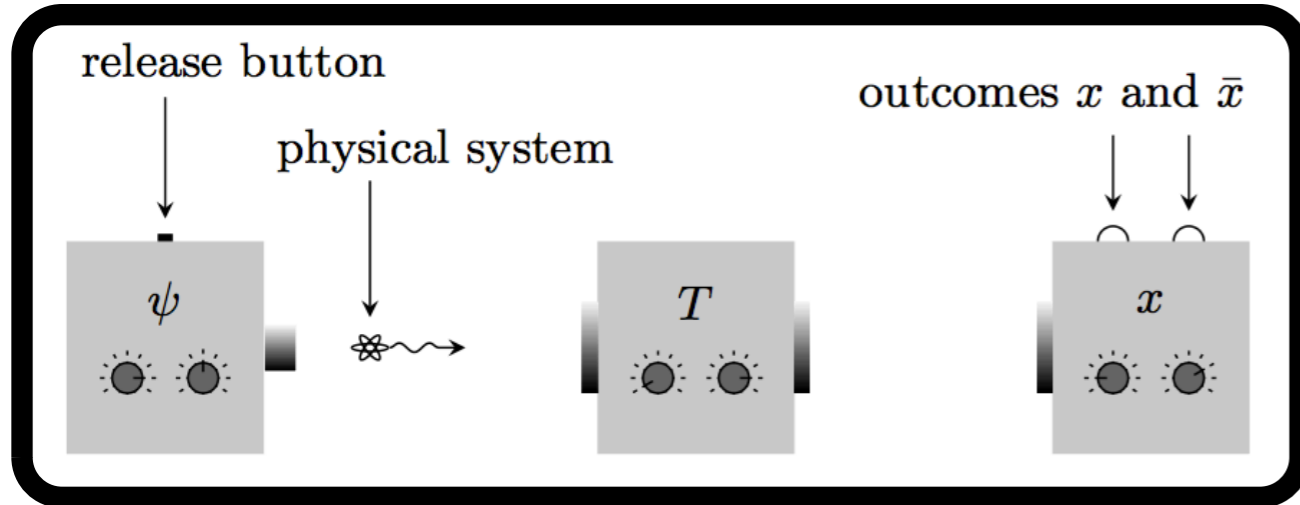
- L. Hardy, *Quantum Theory From Five Reasonable Axioms*, 2001
- B. Dakić and Č. Brukner, *Quantum Theory and Beyond: Is Entanglement Special?*, 2009

See also:

- G. Chiribella et al., *Informational derivation of Q.T.*, 2010
- L. Hardy, *Reformulating and Reconstructing Q.T.*, 2011



Basic physical / operational assumptions

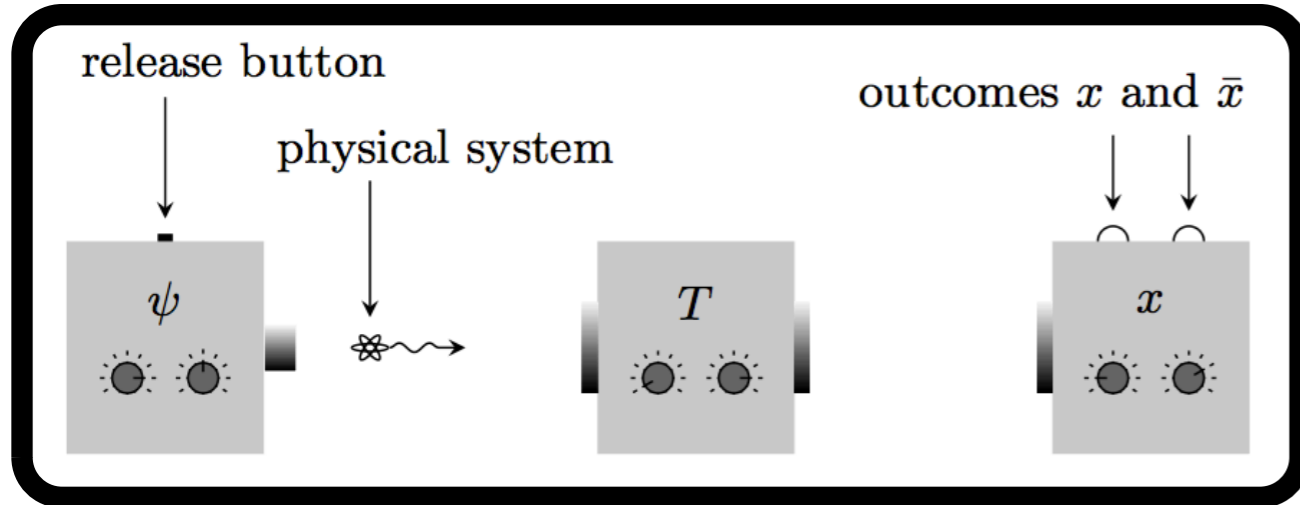
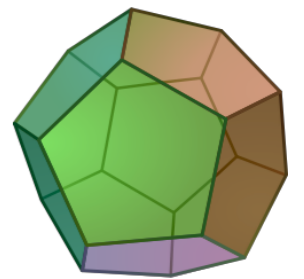


- States, transformations, and measurements with **outcome probabilities**.
- Combined systems: **no-signalling**.

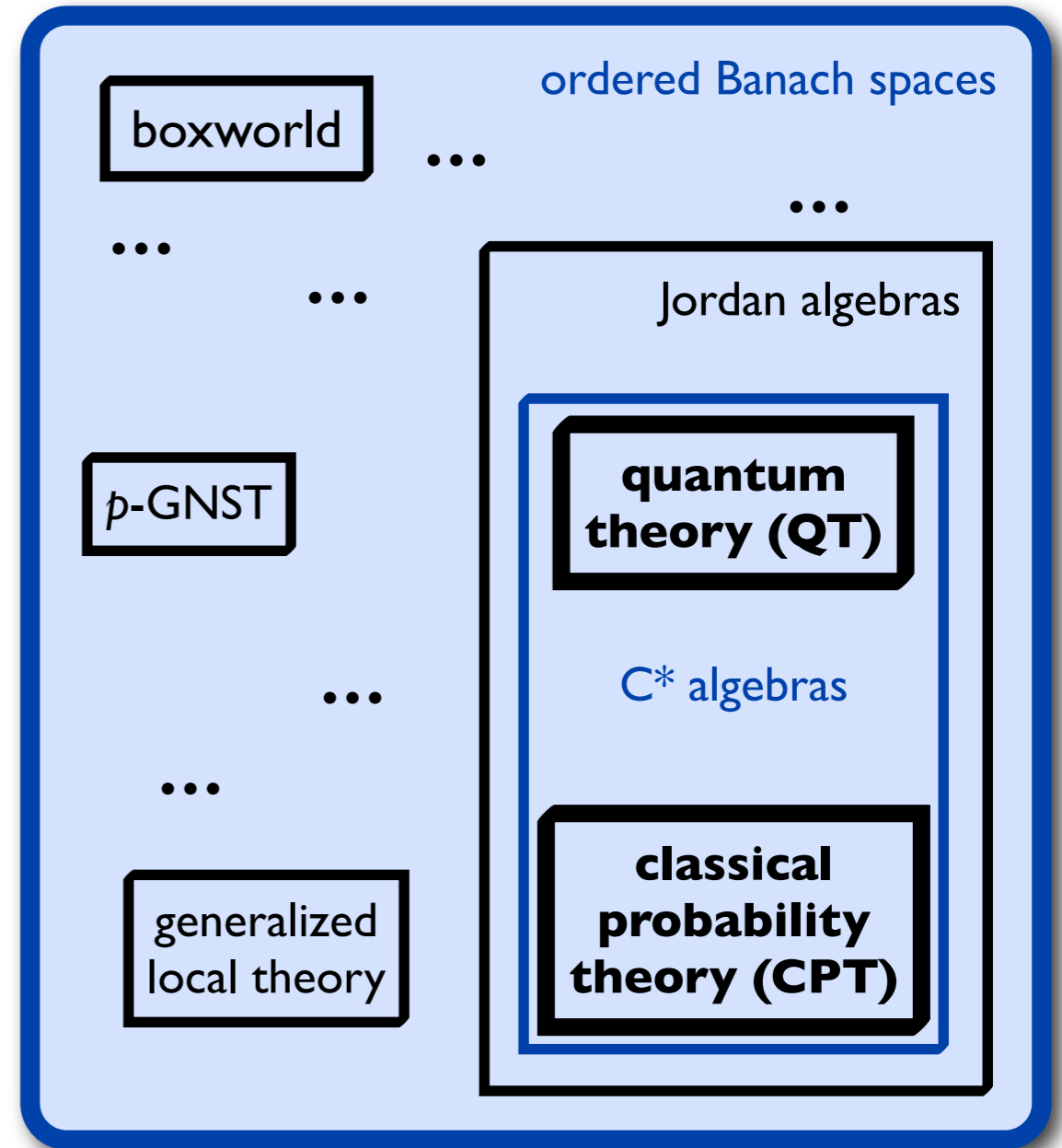
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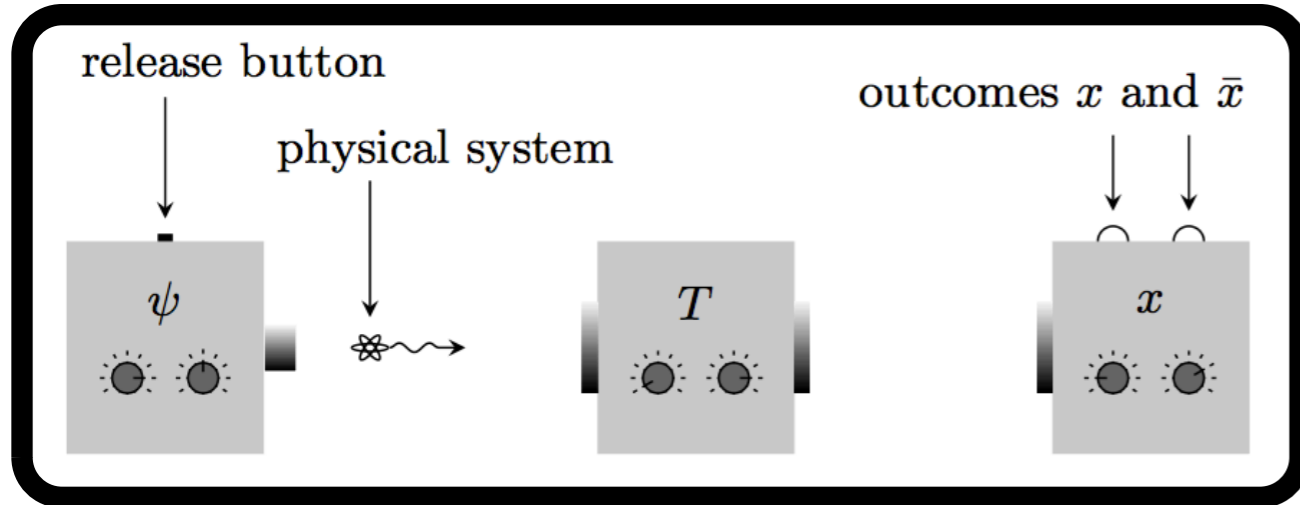
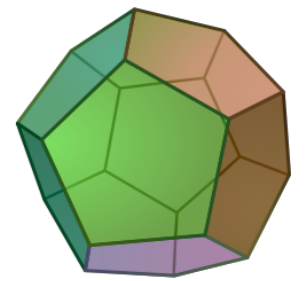


- **No** Hilbert spaces, complex numbers,...
- State spaces: **arbitrary convex sets**.
- Many ways to **combine systems**.

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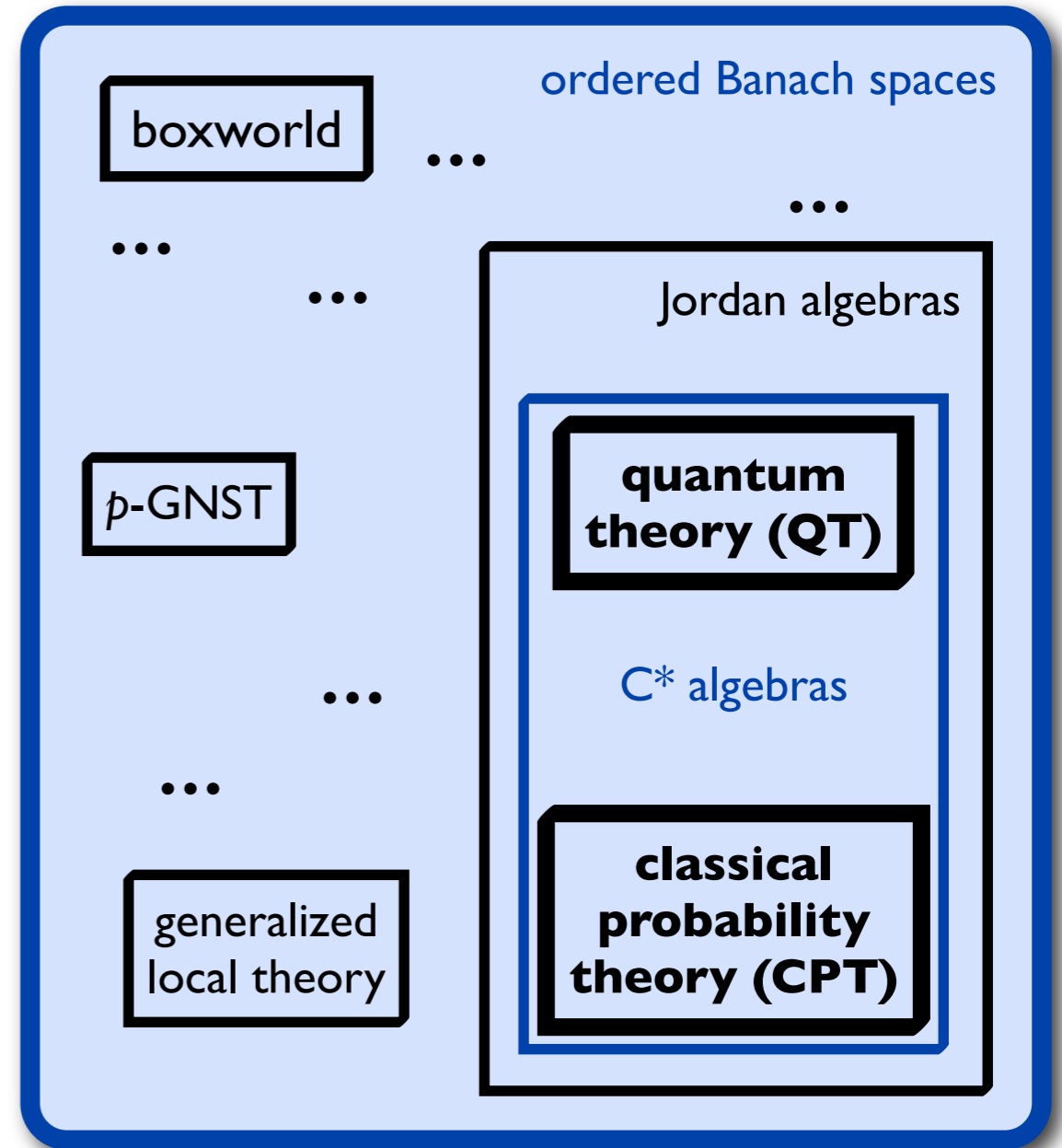
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The Axioms:

- I. Local tomography
- II. Reversibility
- III. Subspace axiom
- IV. Finite-dimensionality
- V. All measurements allowed

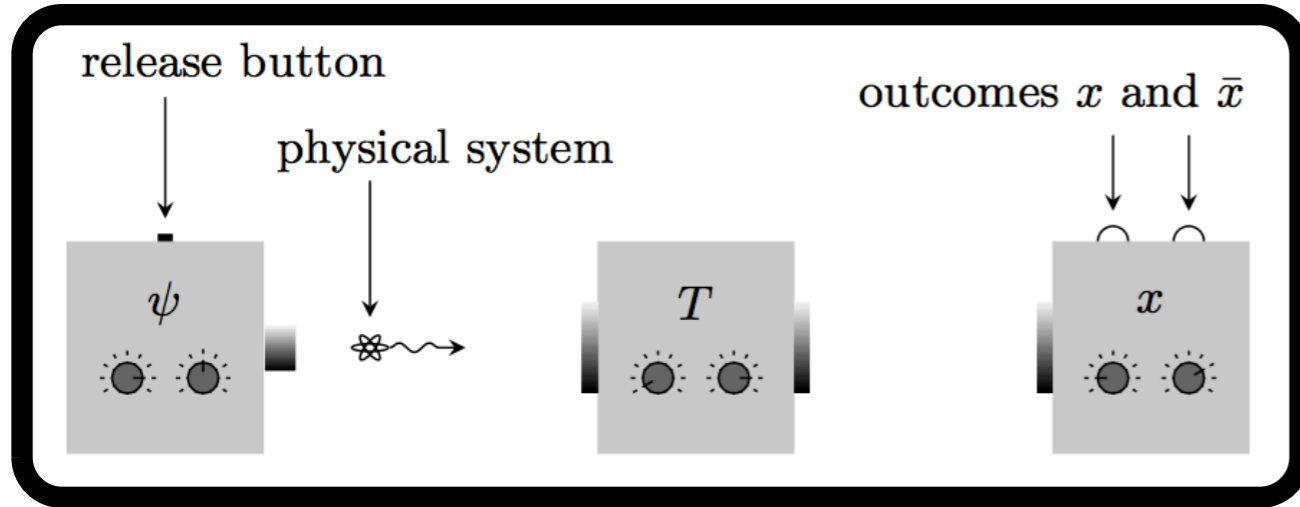
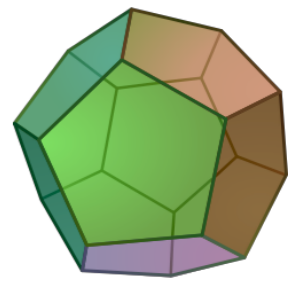


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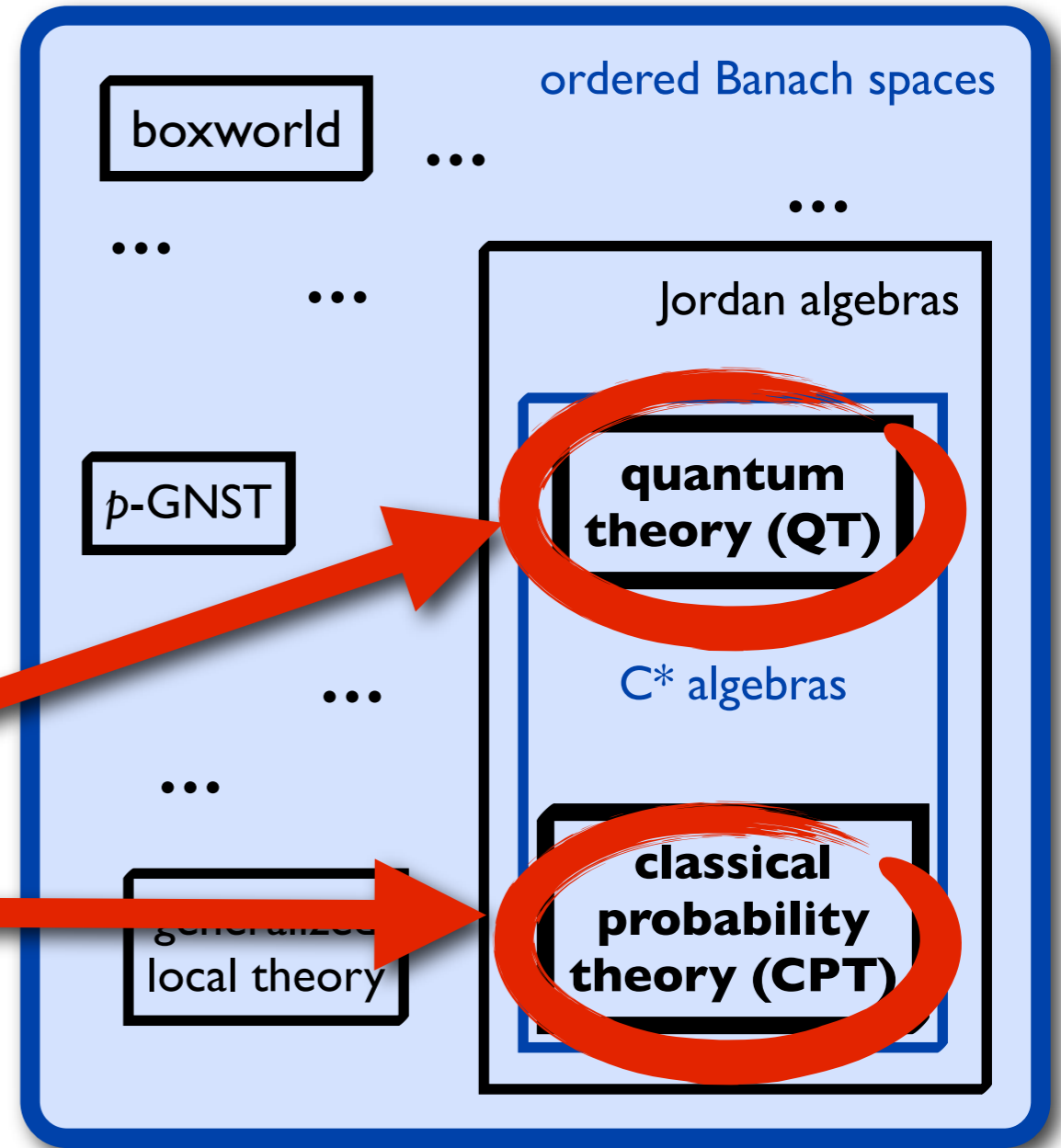
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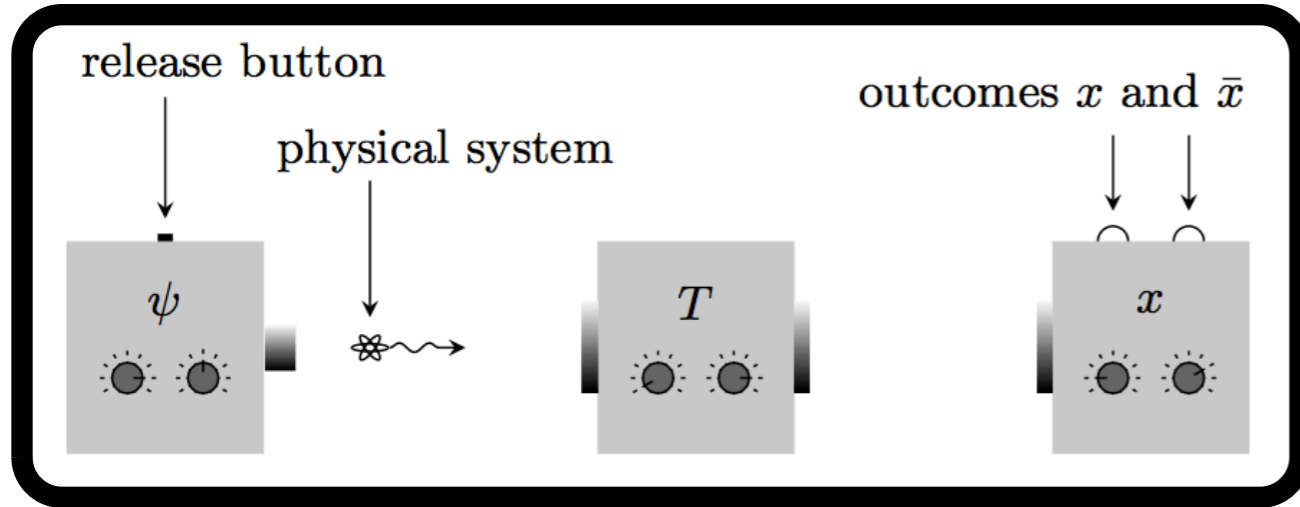
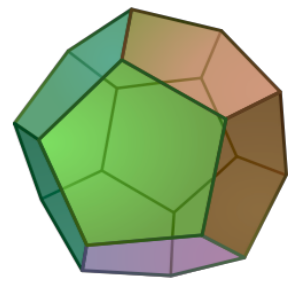
Determine CPT+QT uniquely!

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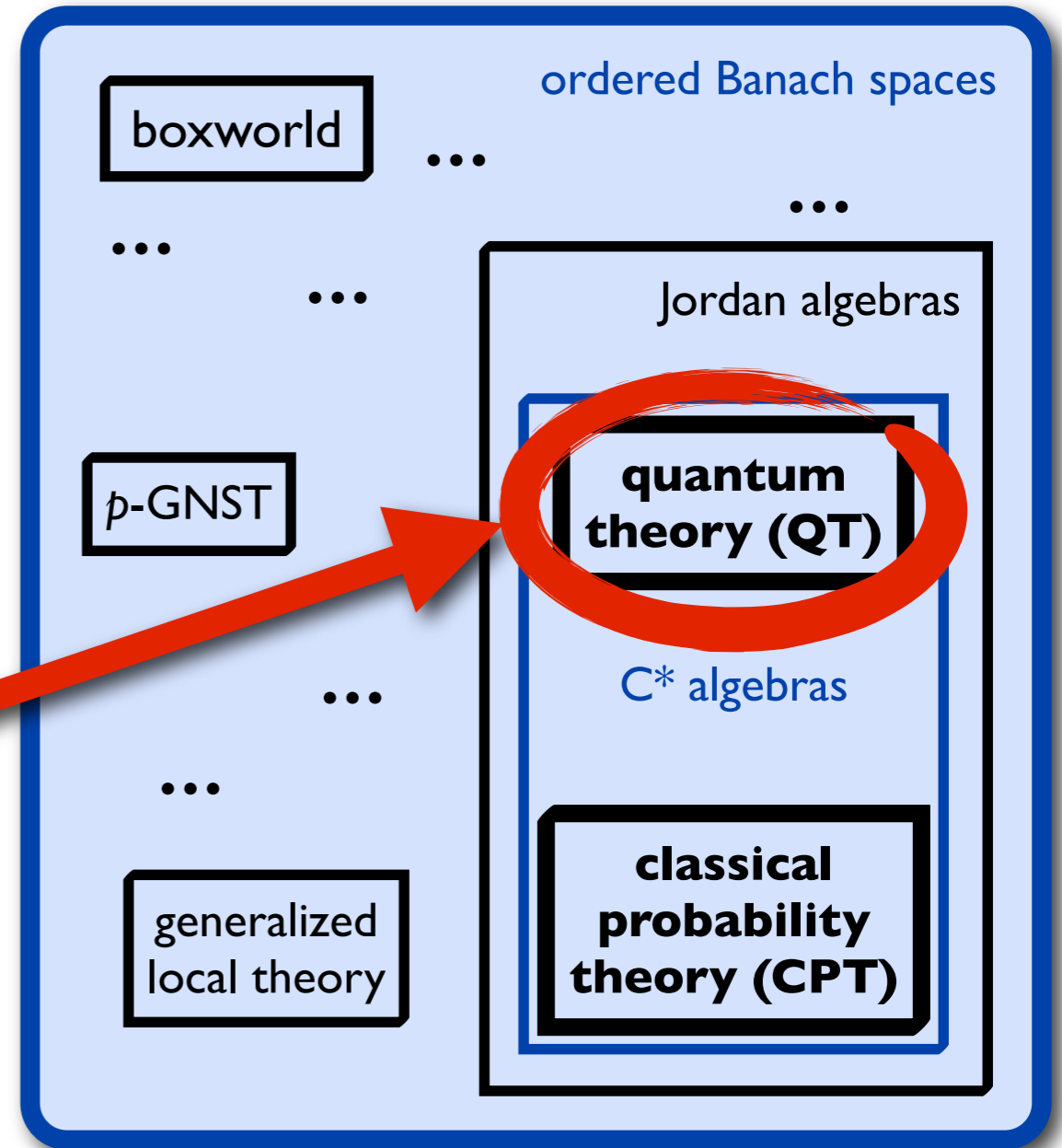
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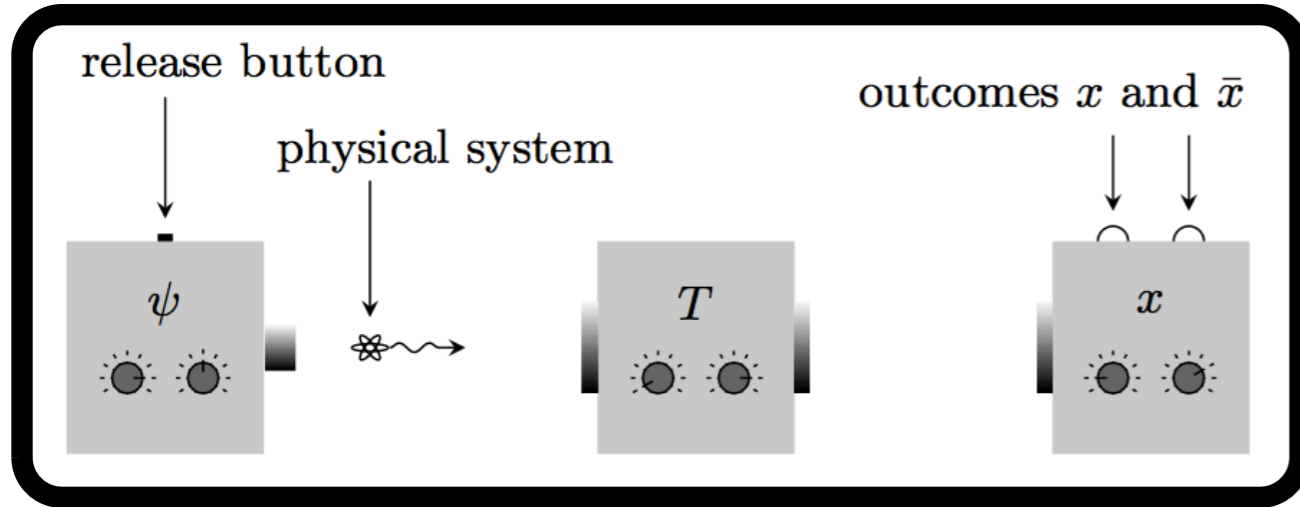
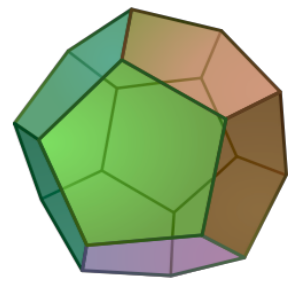
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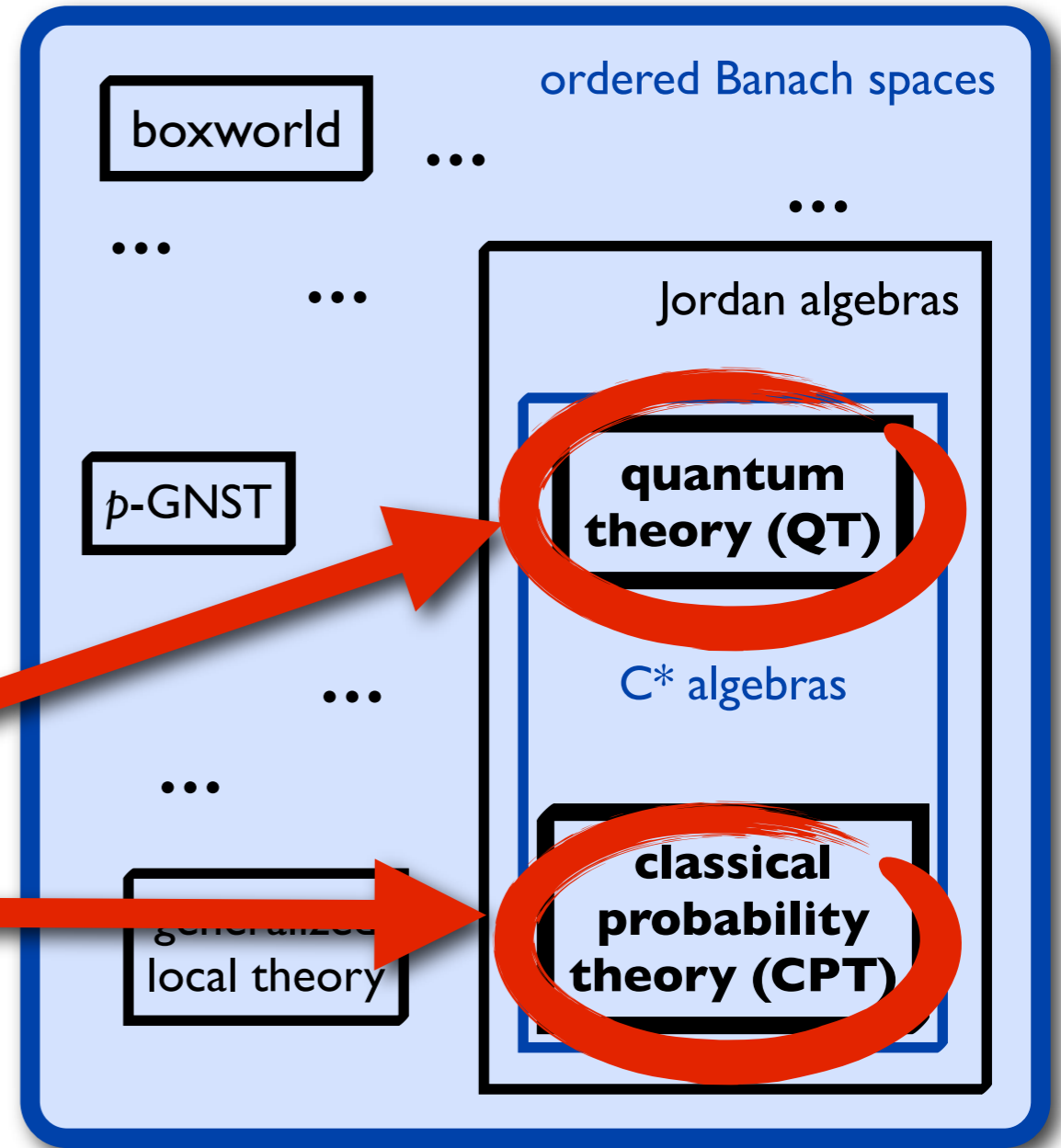
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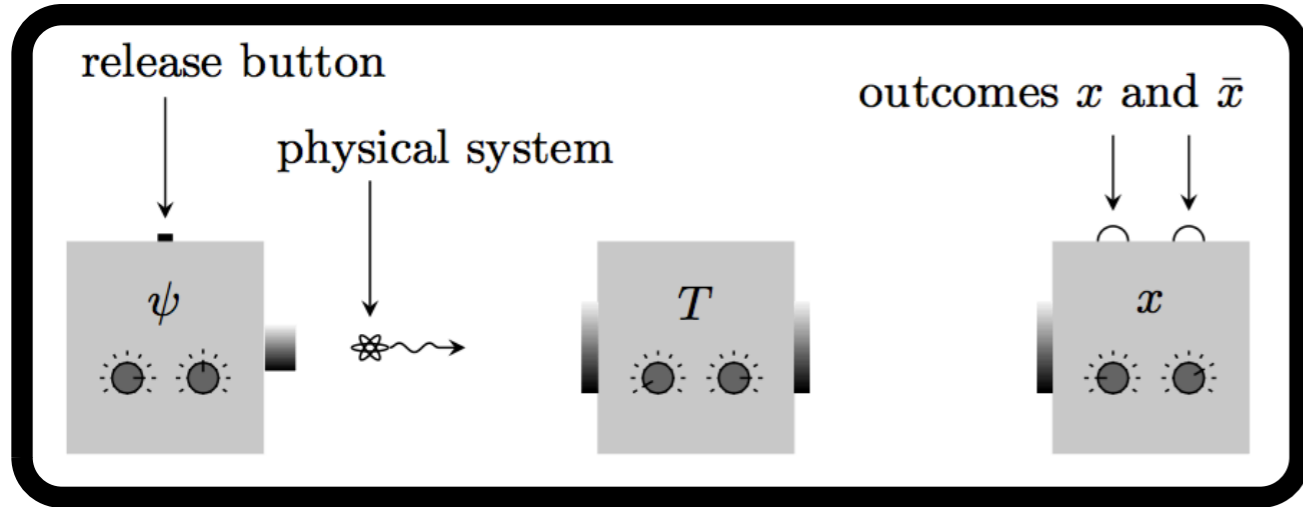
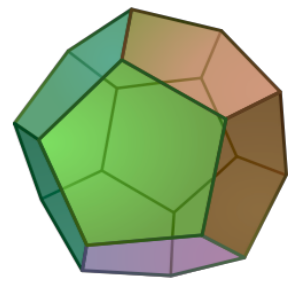
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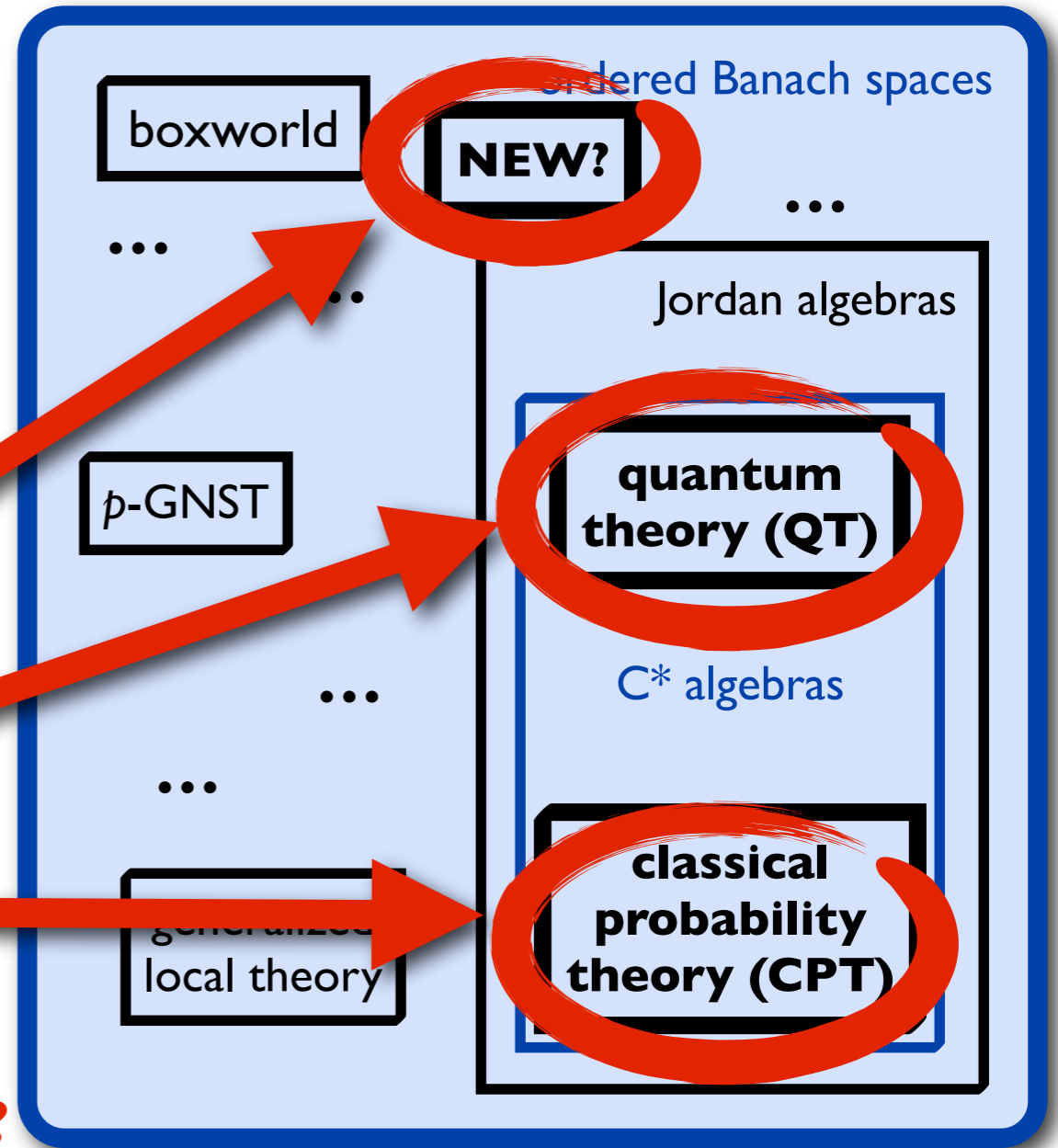
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Natural theories beyond QT?

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What our results are **not**:



- They offer **no resolution of the measurement problem.**
- **No new interpretation** of quantum theory.
- We **assume that probabilities** exist.
- Only **finite-dimensional QT** so far.
- **Only abstract QT**, no mechanics / field theory.

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Mechanics
(Hamiltonian,
phase space,...)

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Probability
Theory

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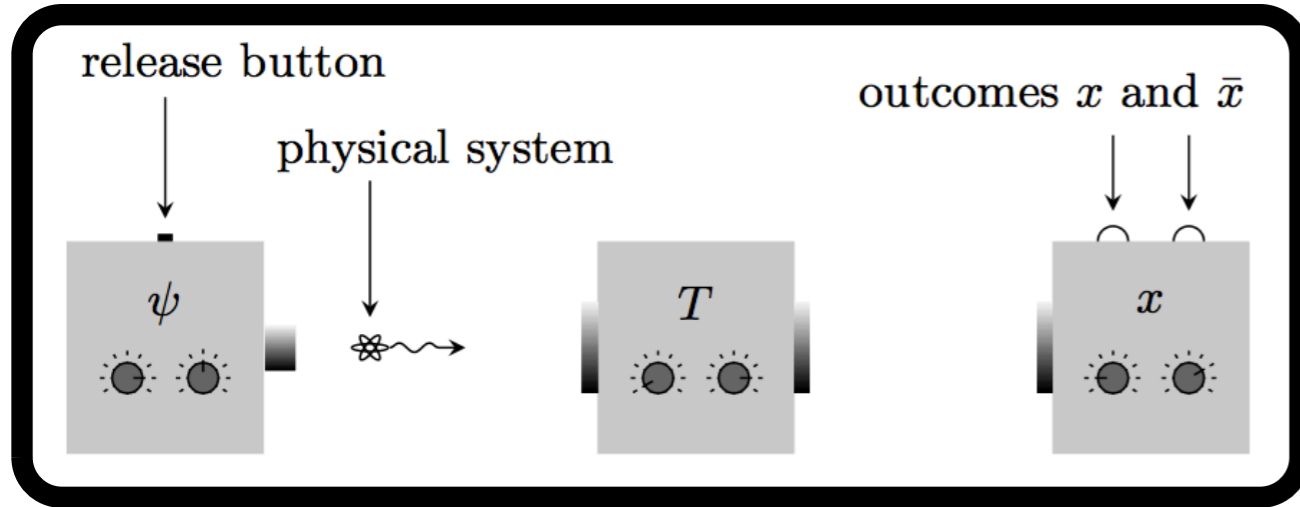
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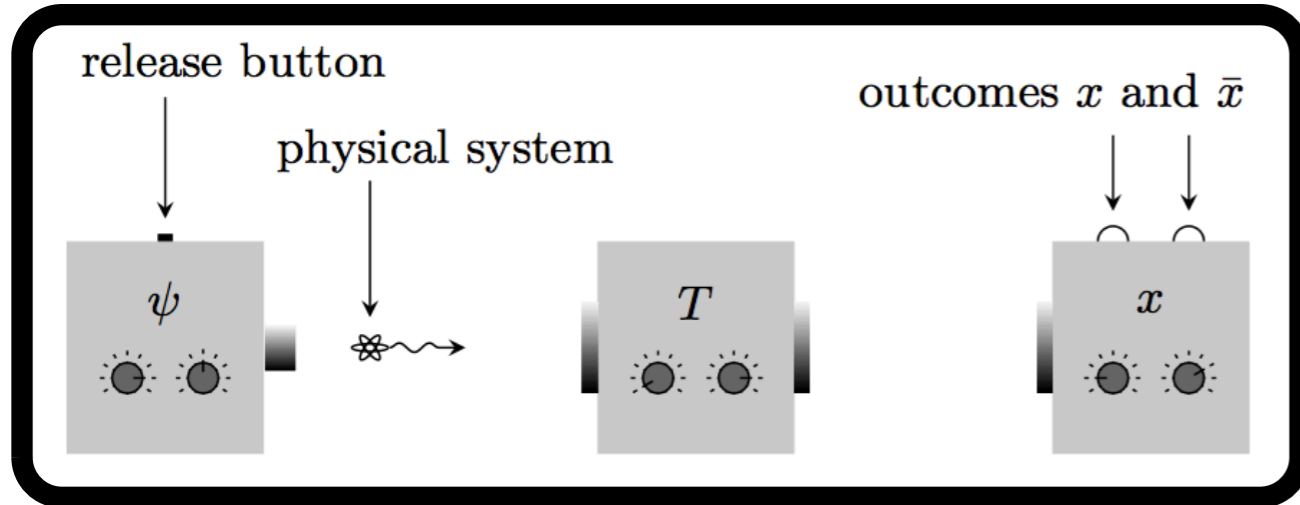
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2. General Probabilistic Theories



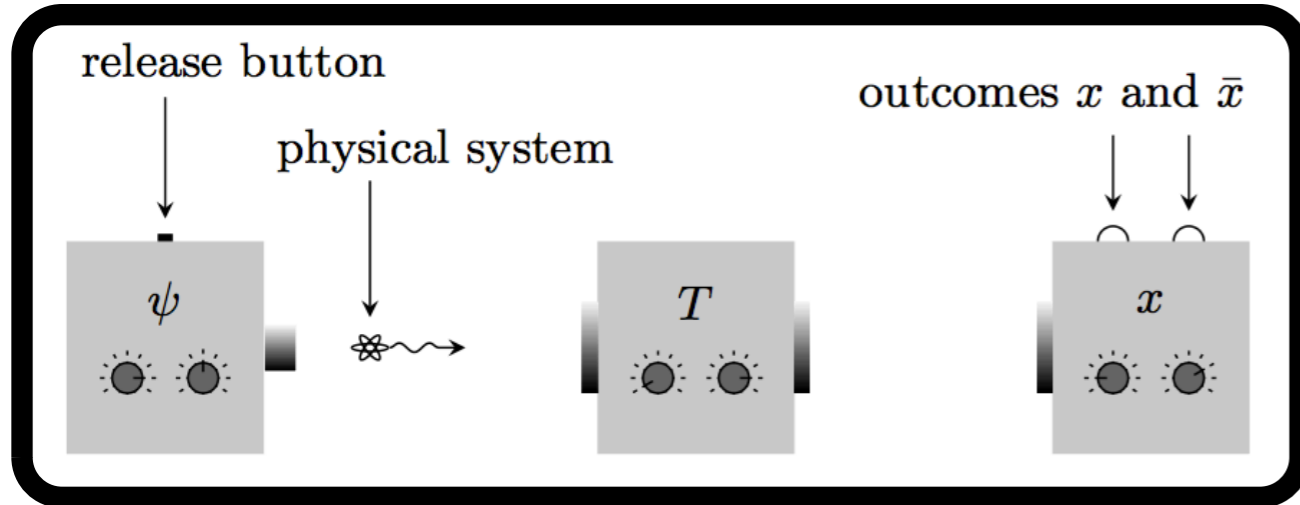
2. General Probabilistic Theories



(Unnormalized) state ω =
list of all probabilities of „yes“-
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$$\omega = (p_1, p_2, p_3, p_4, p_5, p_6, \dots)$$

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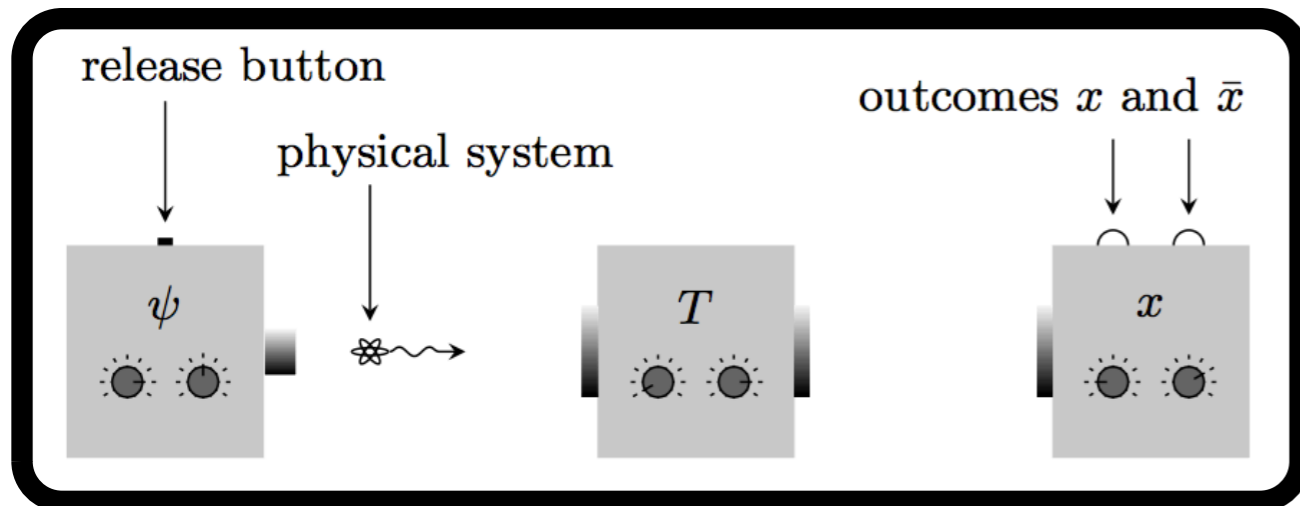
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Sometimes, all ω span a finite-dimensional subspace. Ex.: Qubit

- What's the prob. of „spin up“ in **X**-direction?
- What's the prob. of „spin up“ in **Y**-direction?
- What's the prob. of „spin up“ in **Z**-direction?
- Is the particle there **at all**?

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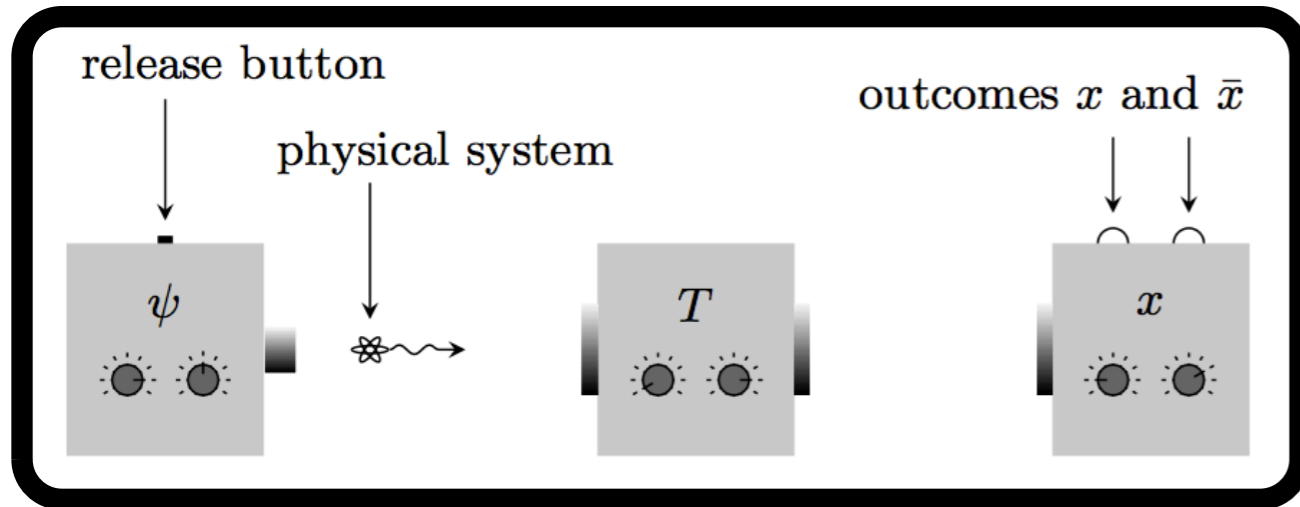
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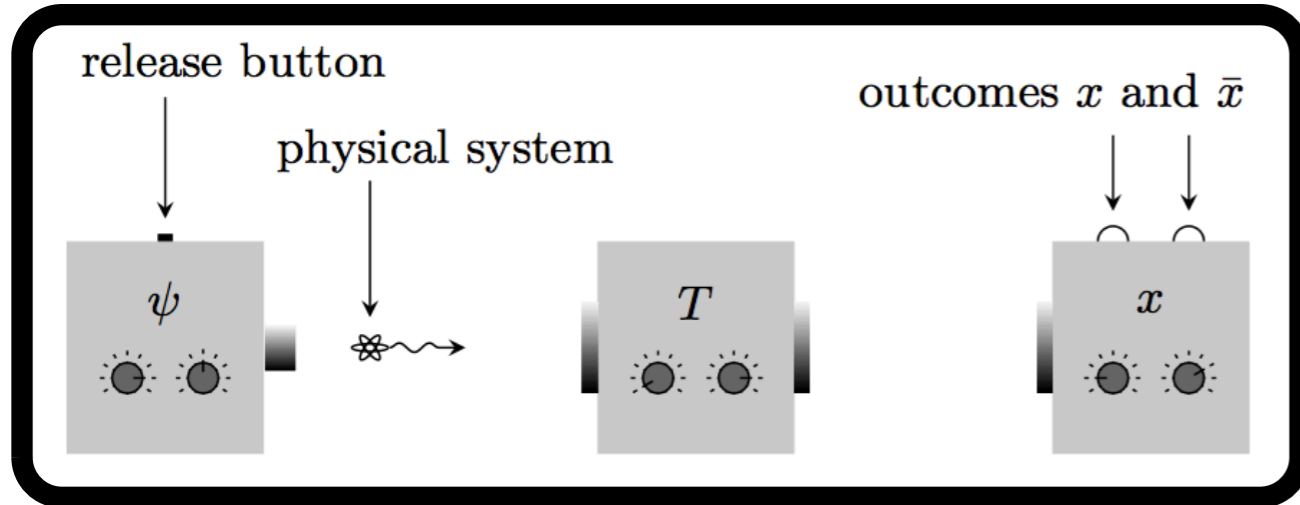
Axiom IV: All state spaces are finite-dimensional.

2. General Probabilistic Theories



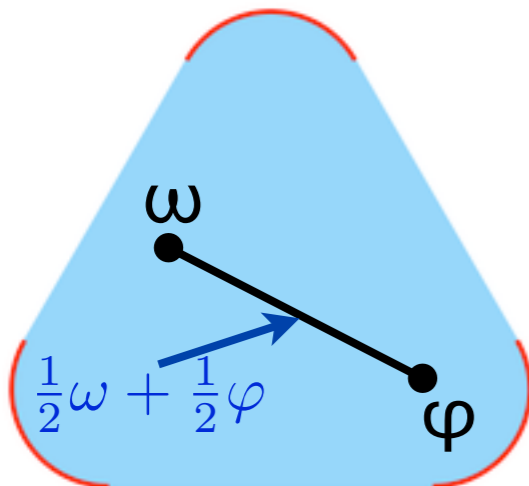
Prepare state ω or φ with prob. $\frac{1}{2}$. Result: $\frac{1}{2}\omega + \frac{1}{2}\varphi$

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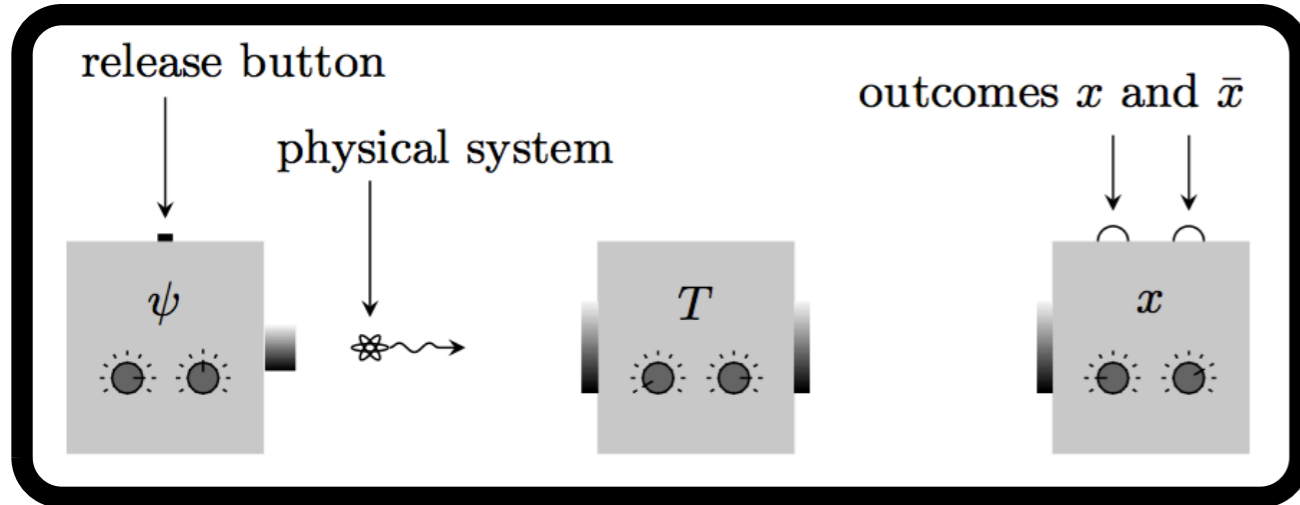


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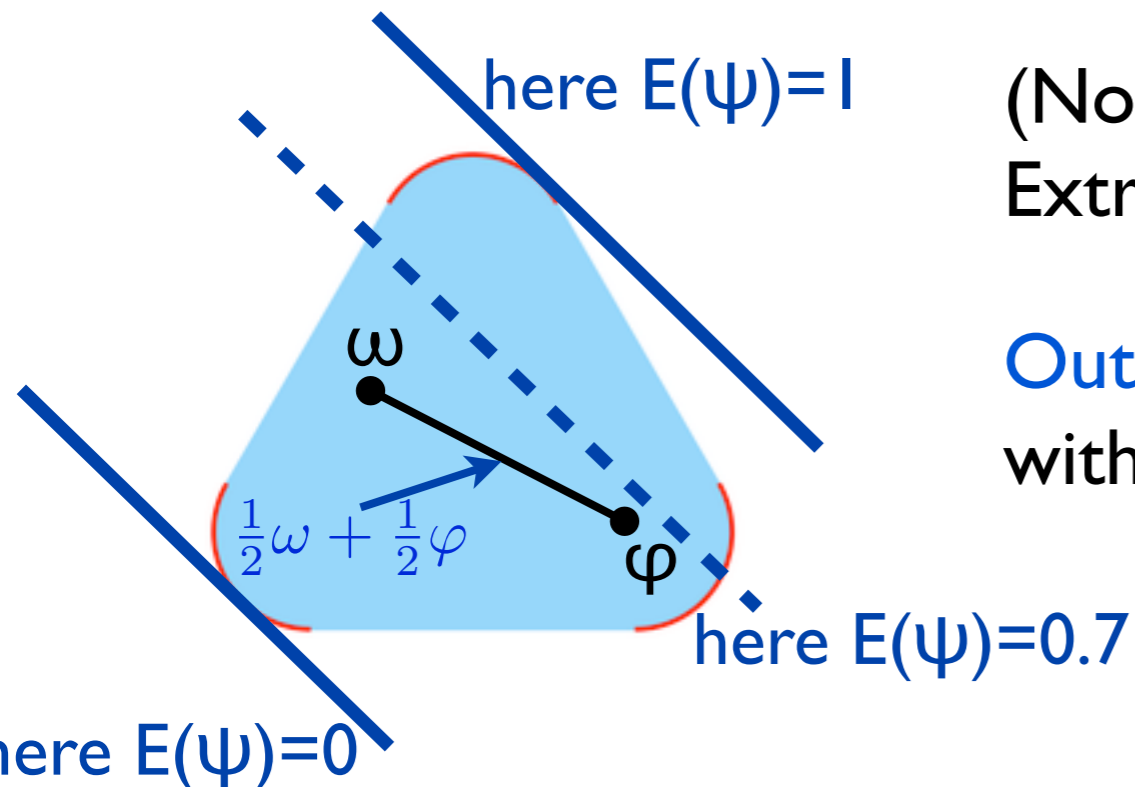
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Extremal points are **pure states**, others **mixed**.



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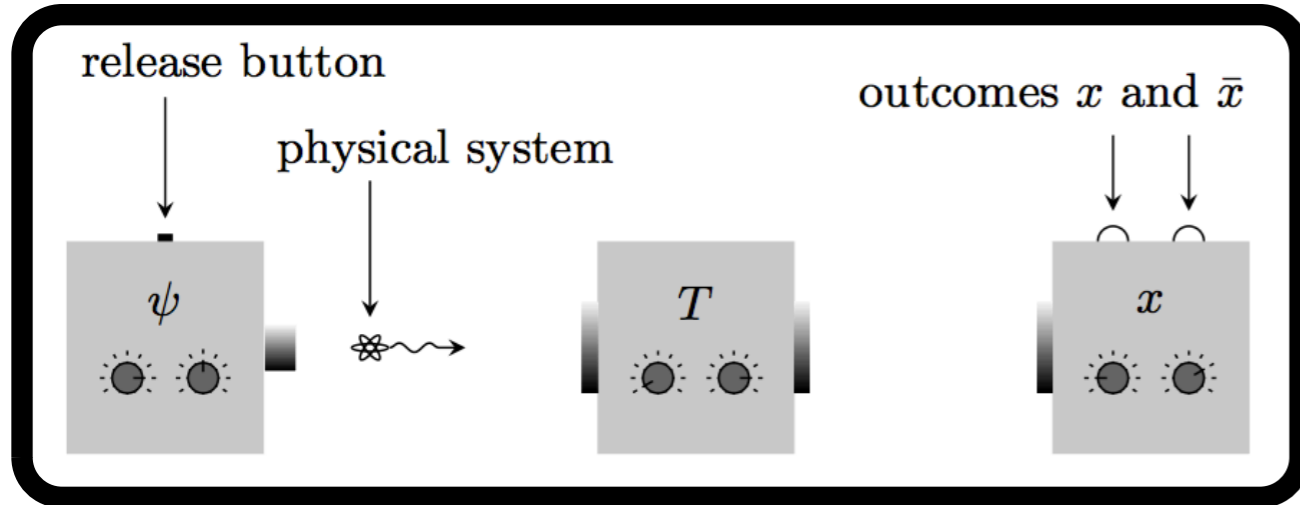
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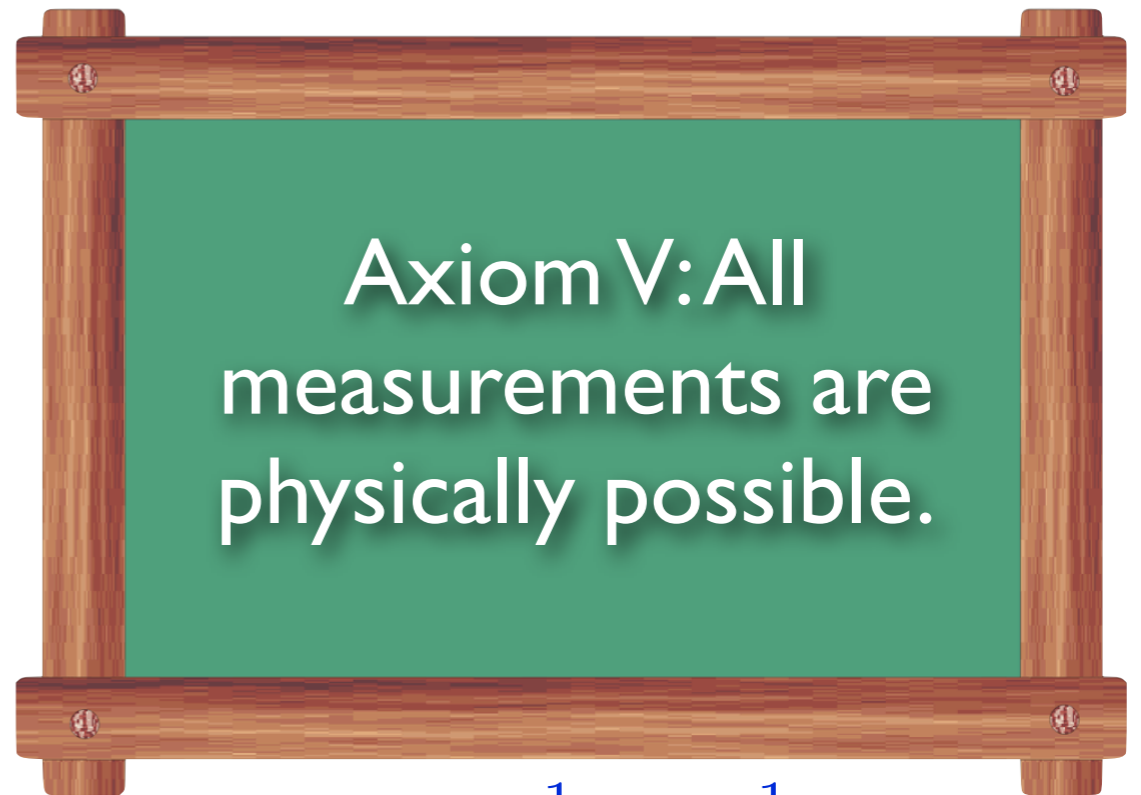
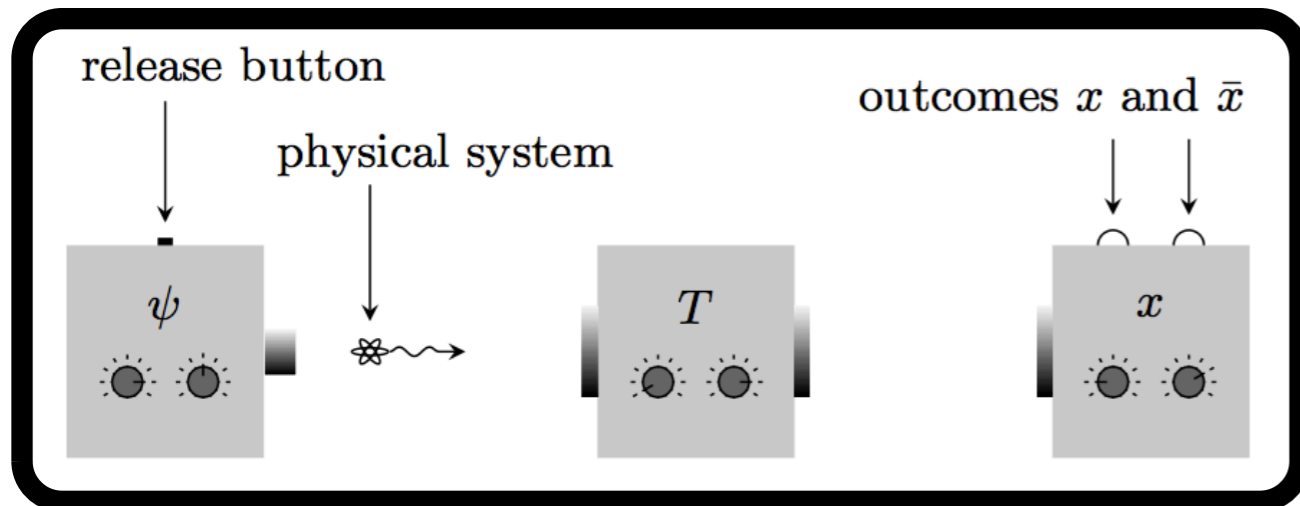
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Measurements are (E_1, E_2, \dots, E_k)
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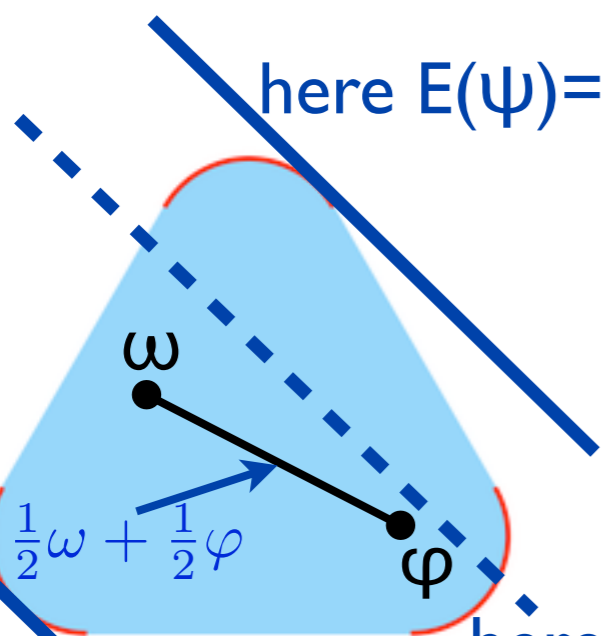
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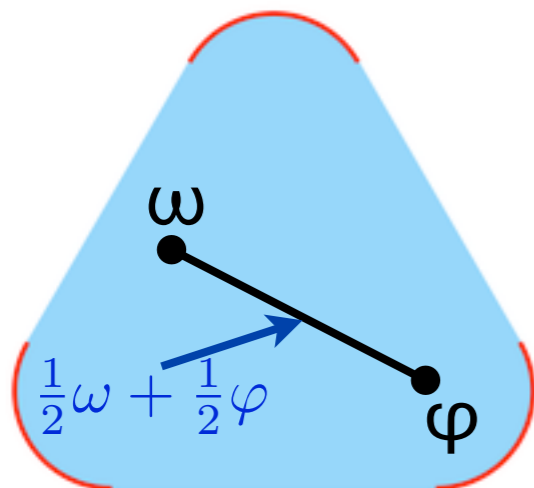
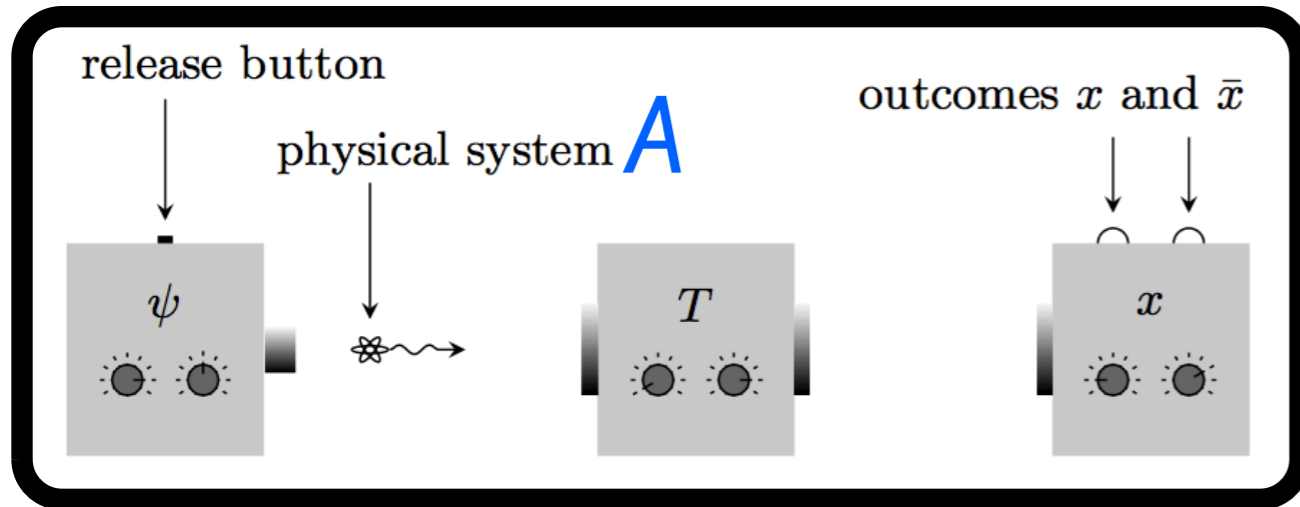
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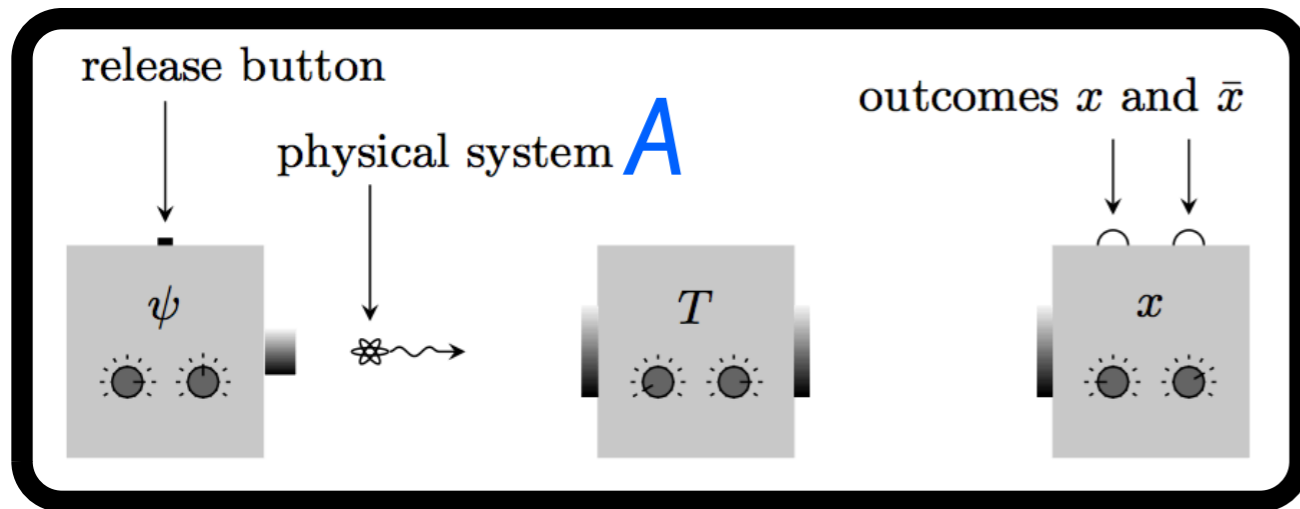


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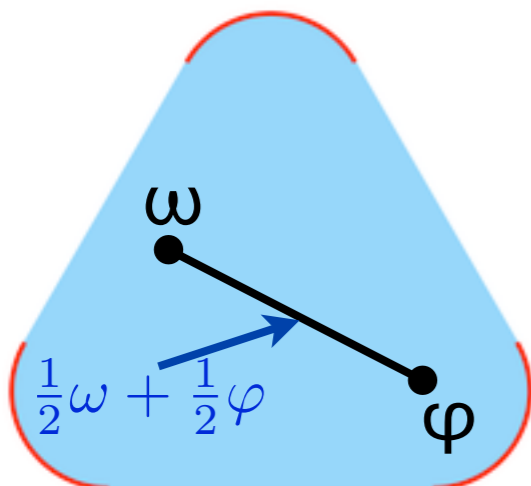


Normalized state space Ω_A

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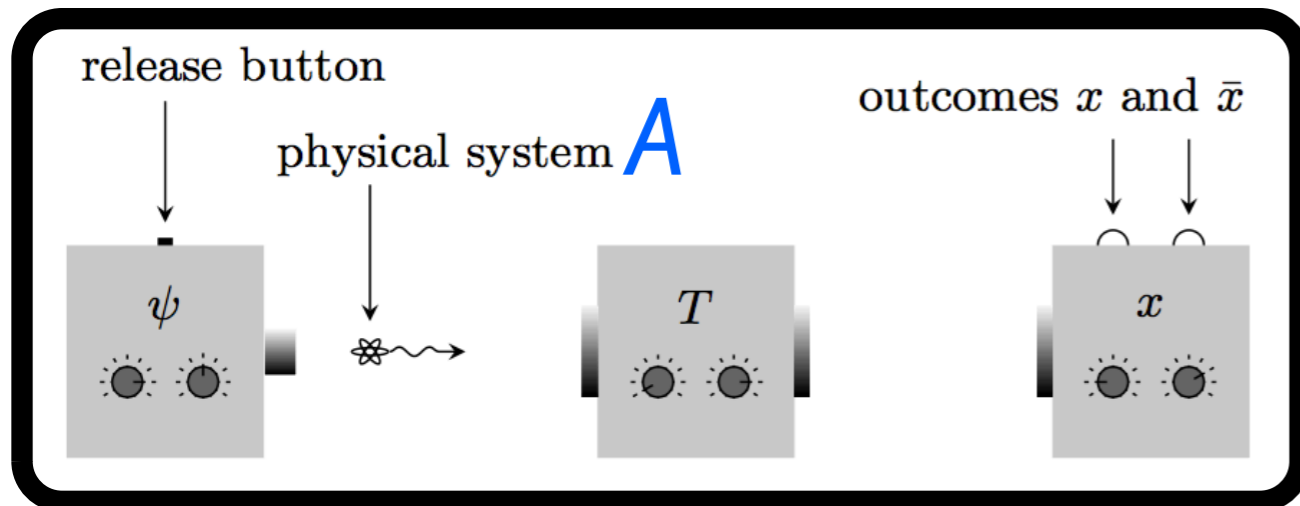


Transformations T map (unnormalized) states to states, and are **linear**.



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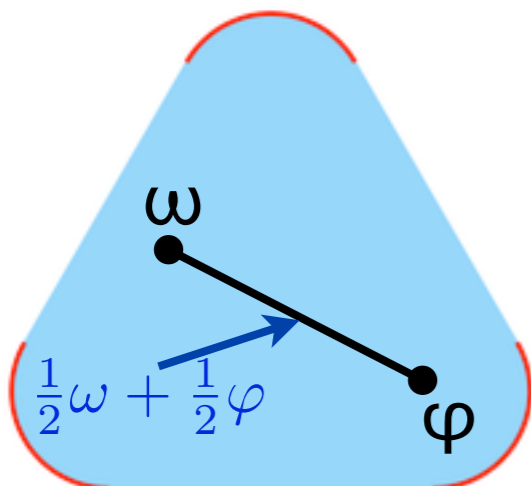
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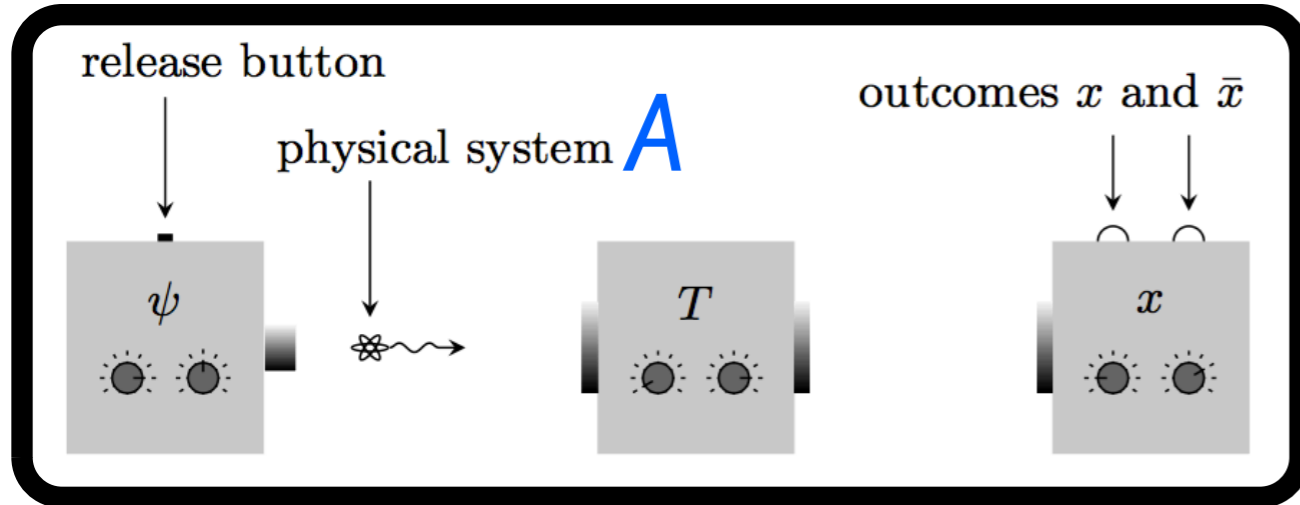
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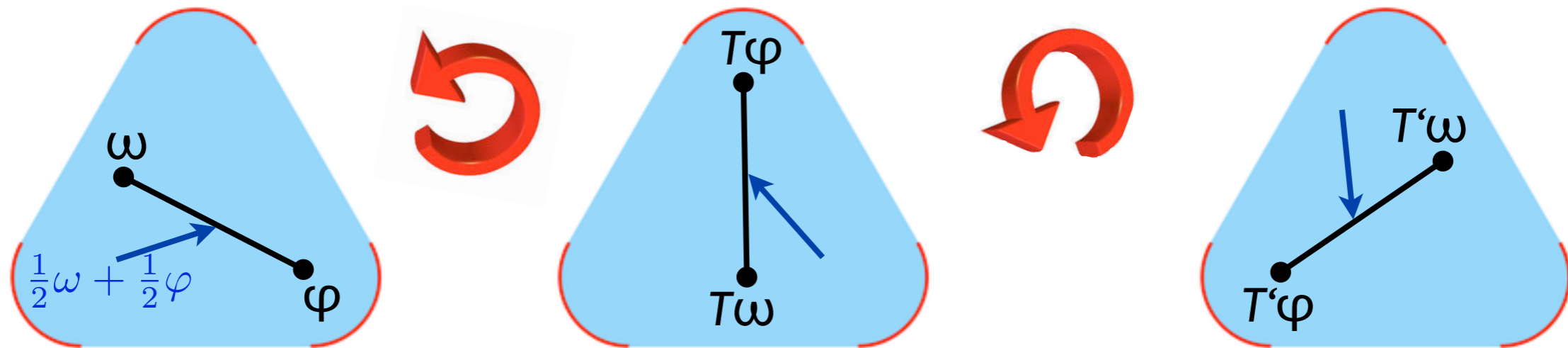
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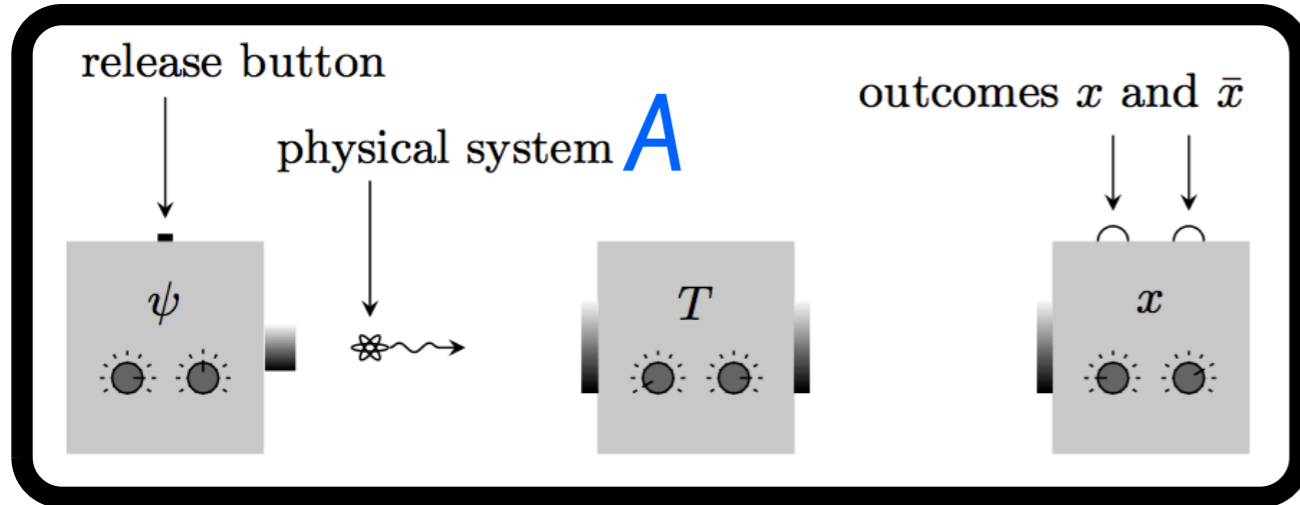
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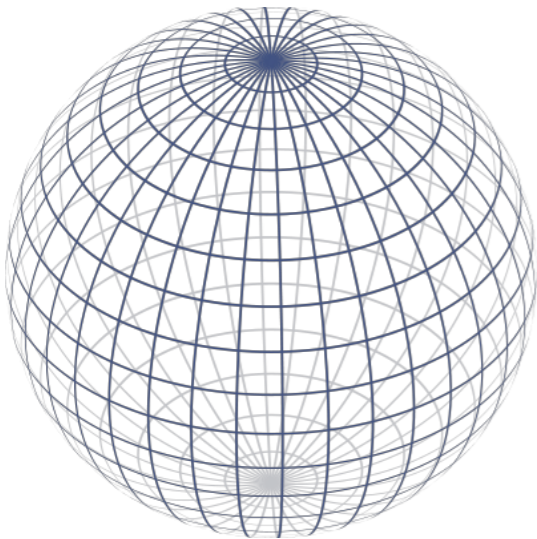


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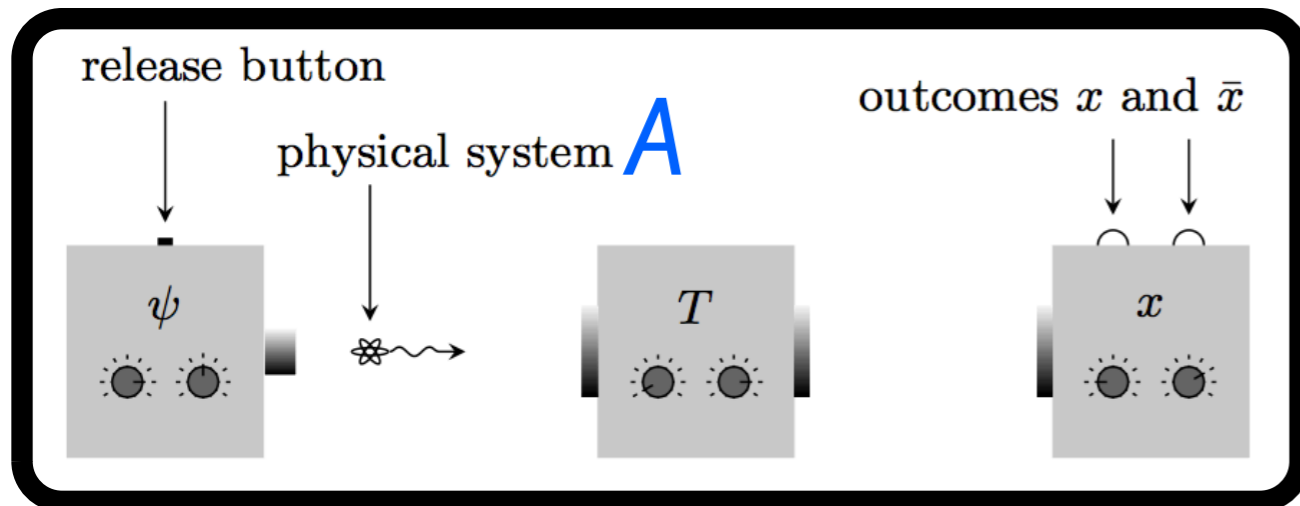


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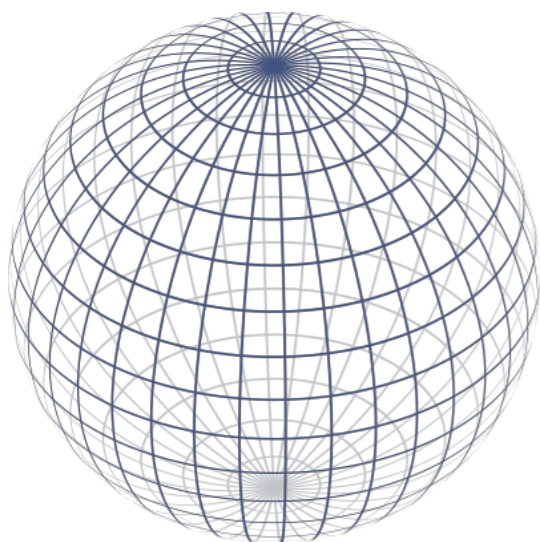


Qubit: Ω_A is the 3D unit ball,
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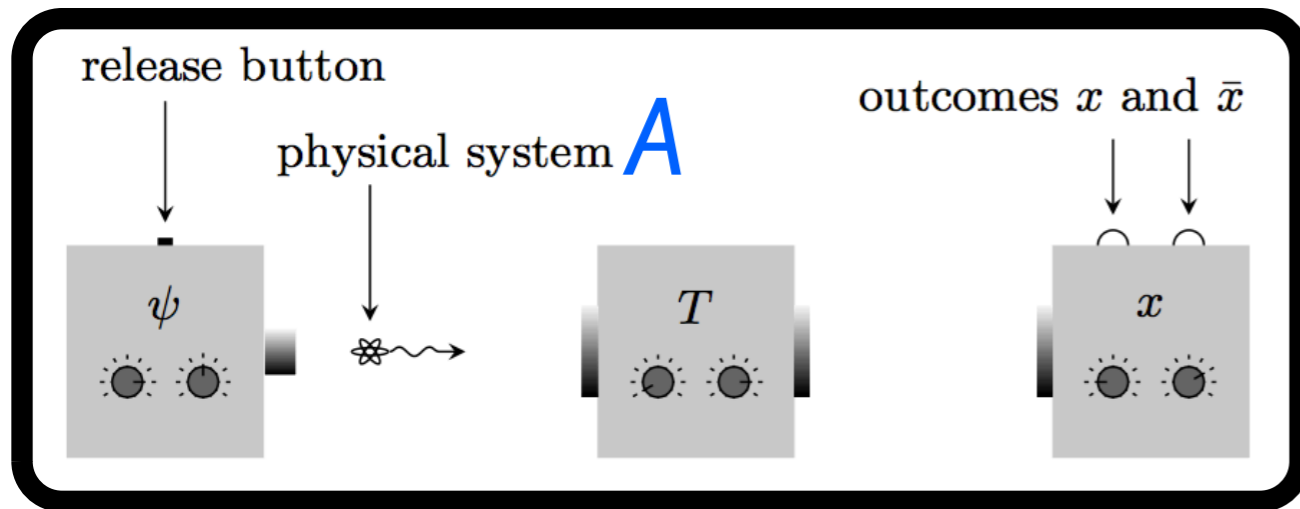


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\Rightarrow A **system** is a pair $(\Omega_A, \mathcal{G}_A)$.

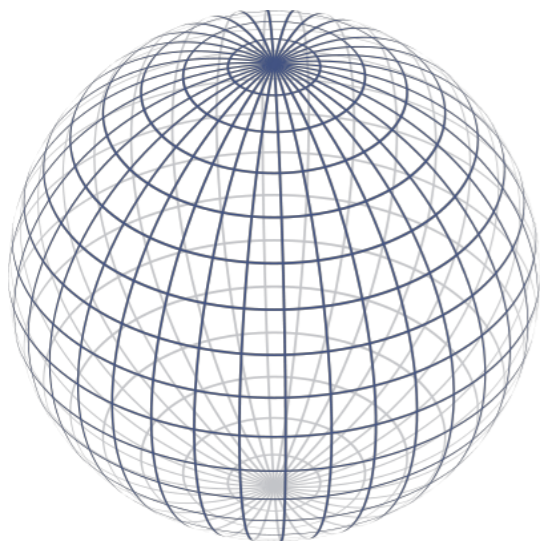
state space \nearrow Ω_A \nwarrow reversible transformations \mathcal{G}_A

2. General Probabilistic Theories



Axiom II (Reversibility):
 If φ and ω are **pure**, then
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 with $T\varphi = \omega$.

Not all symmetries have to be in \mathcal{G}_A .

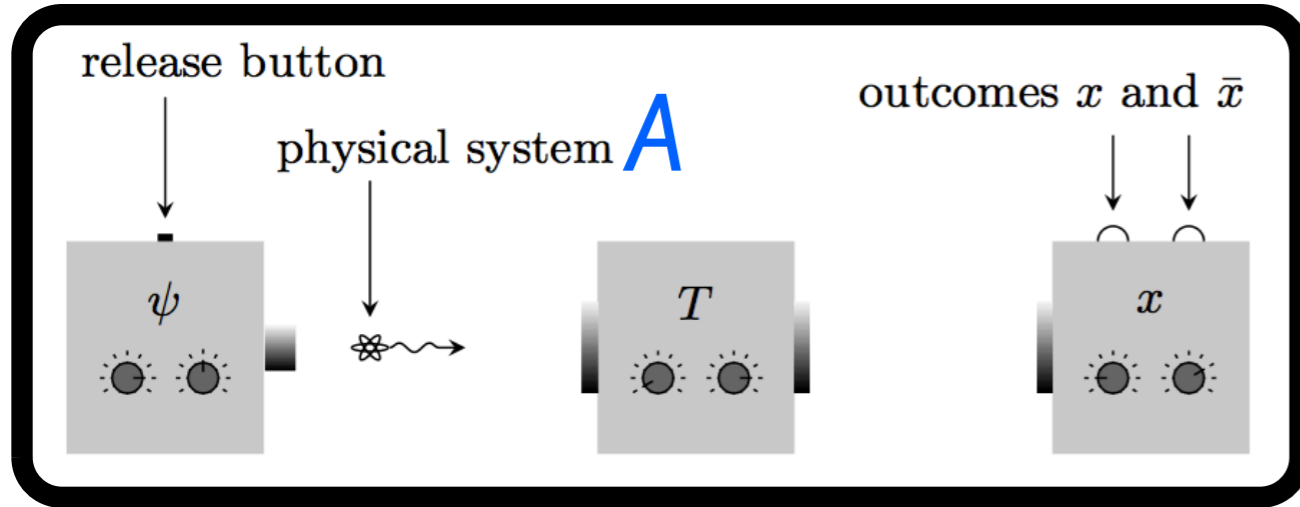


Qubit: Ω_A is the 3D unit ball,
 $\mathcal{G}_A = SO(3)$ (no reflections!)

\Rightarrow A **system** is a pair $(\Omega_A, \mathcal{G}_A)$.

state space \nearrow Ω_A \nwarrow reversible transformations \mathcal{G}_A

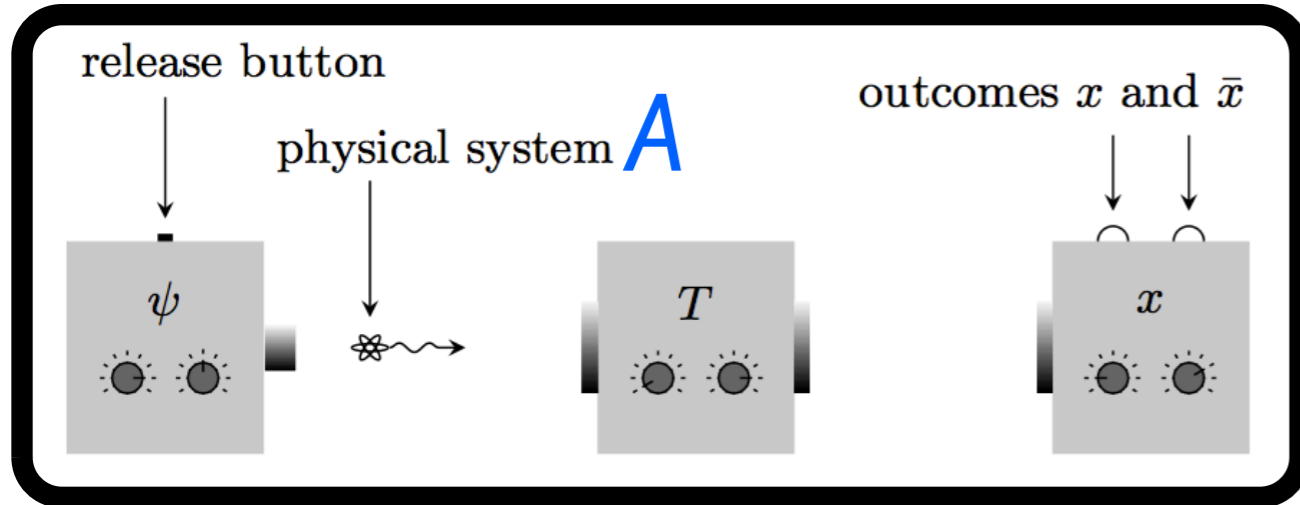
2. General Probabilistic Theories



Axiom II (Reversibility):
If φ and ω are **pure**, then
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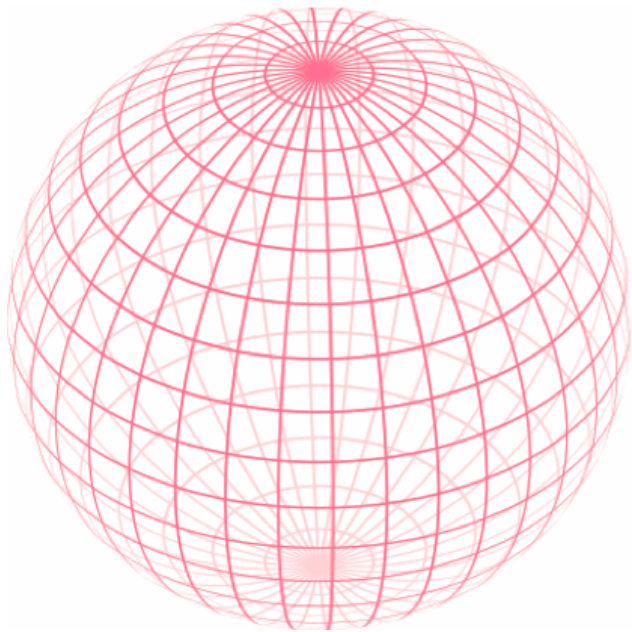
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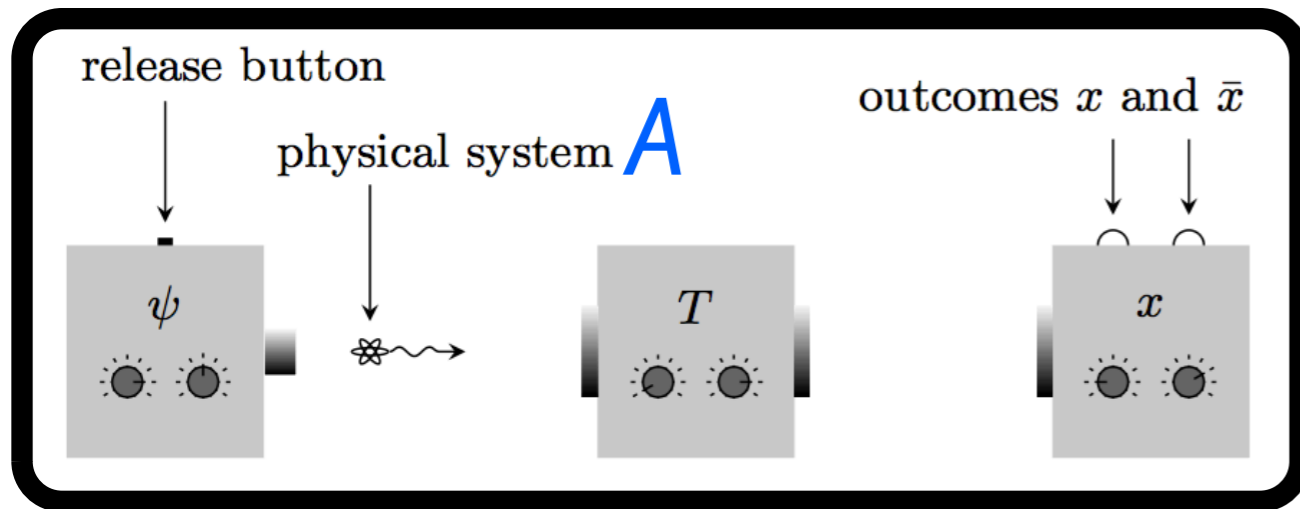


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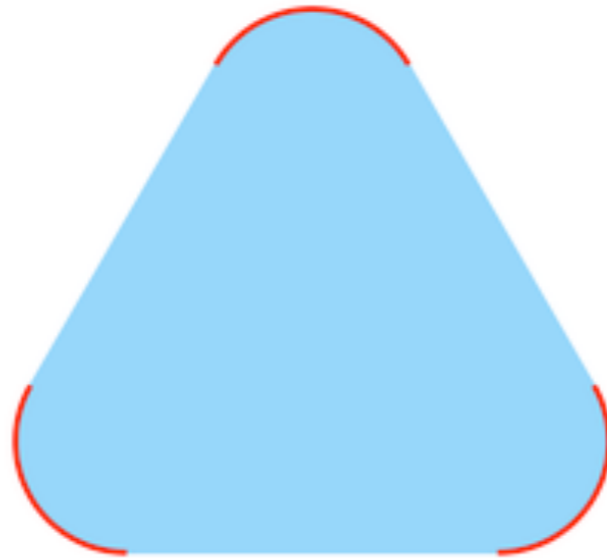
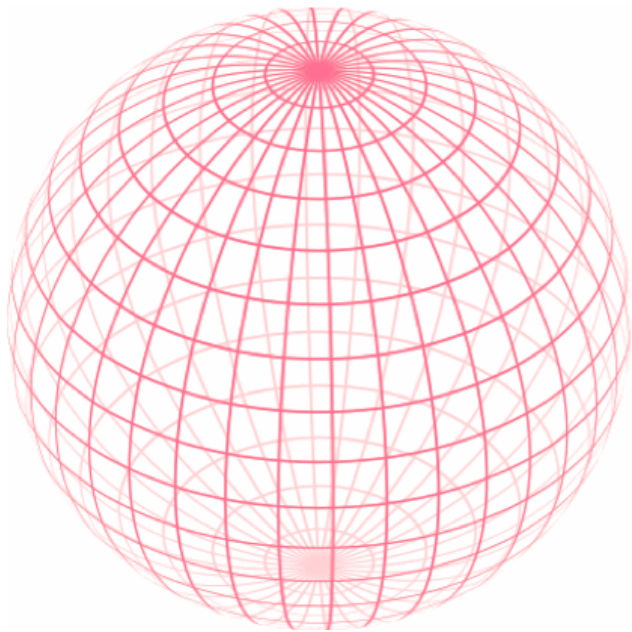


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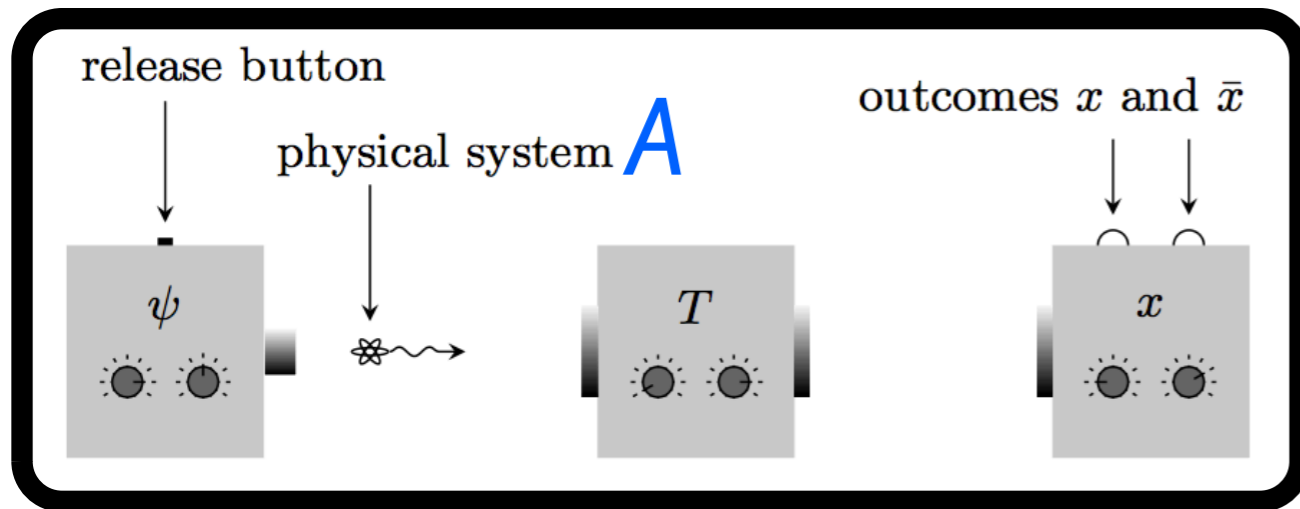


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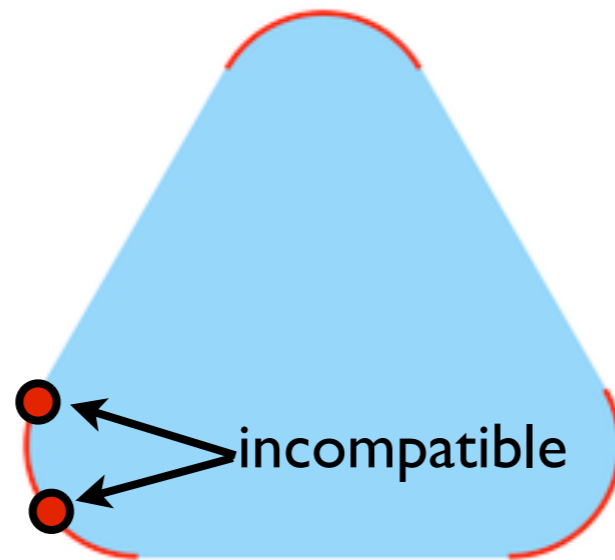
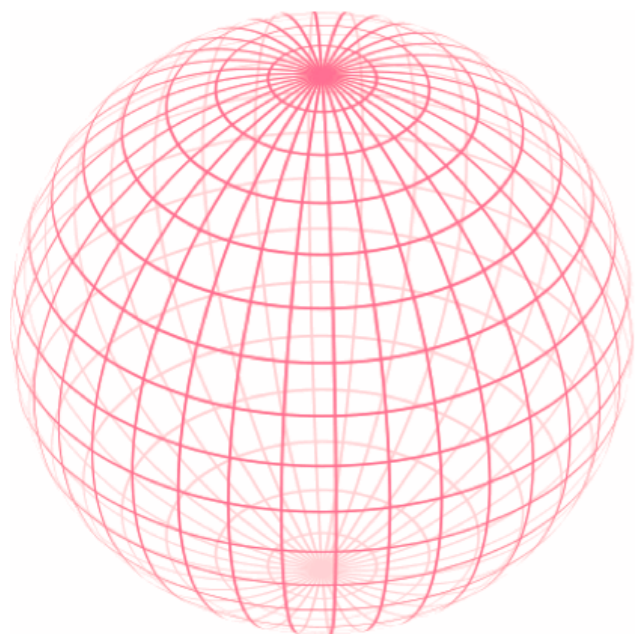


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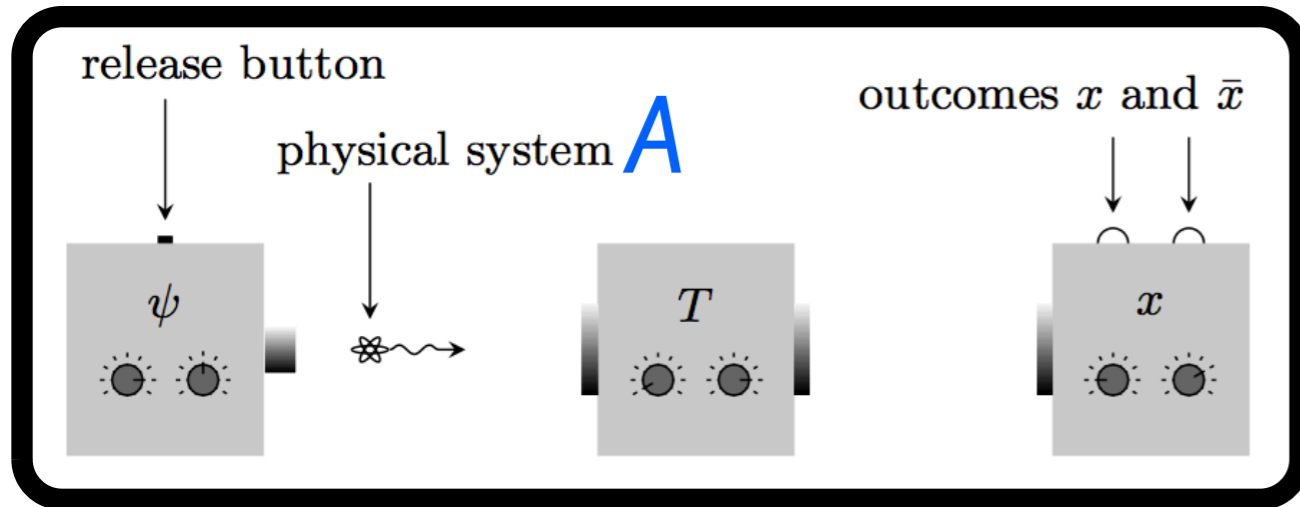


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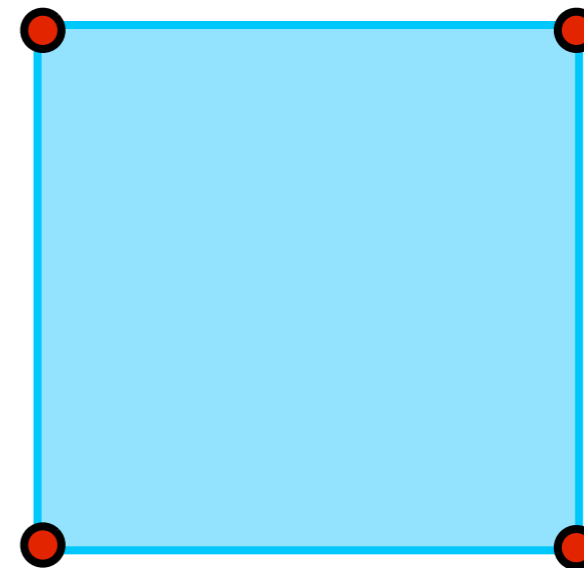
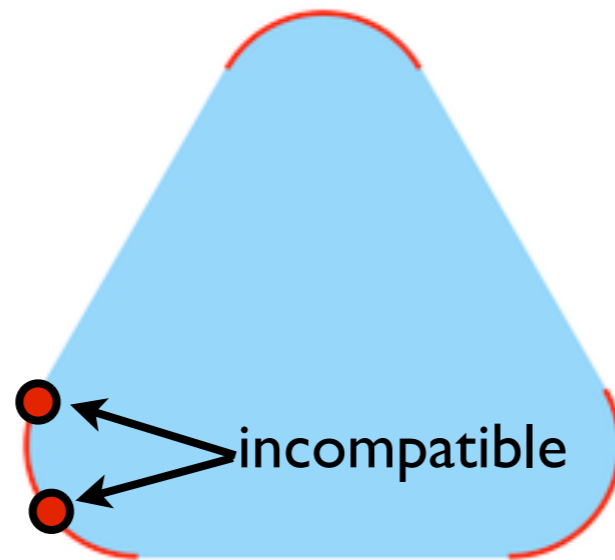
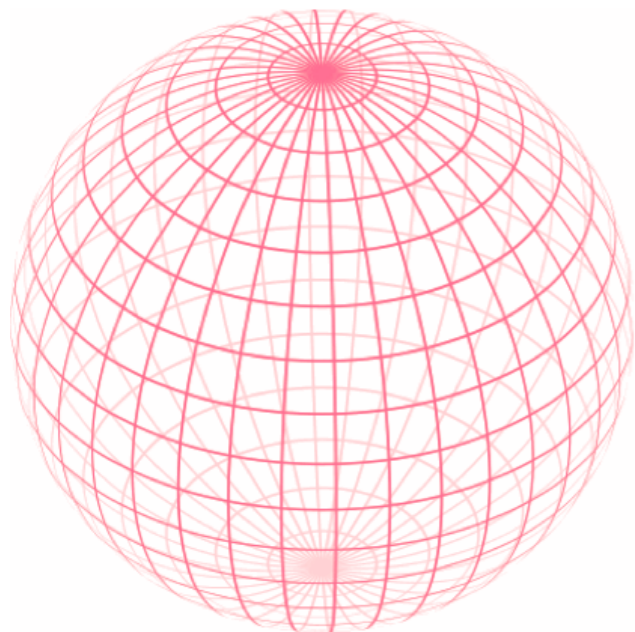


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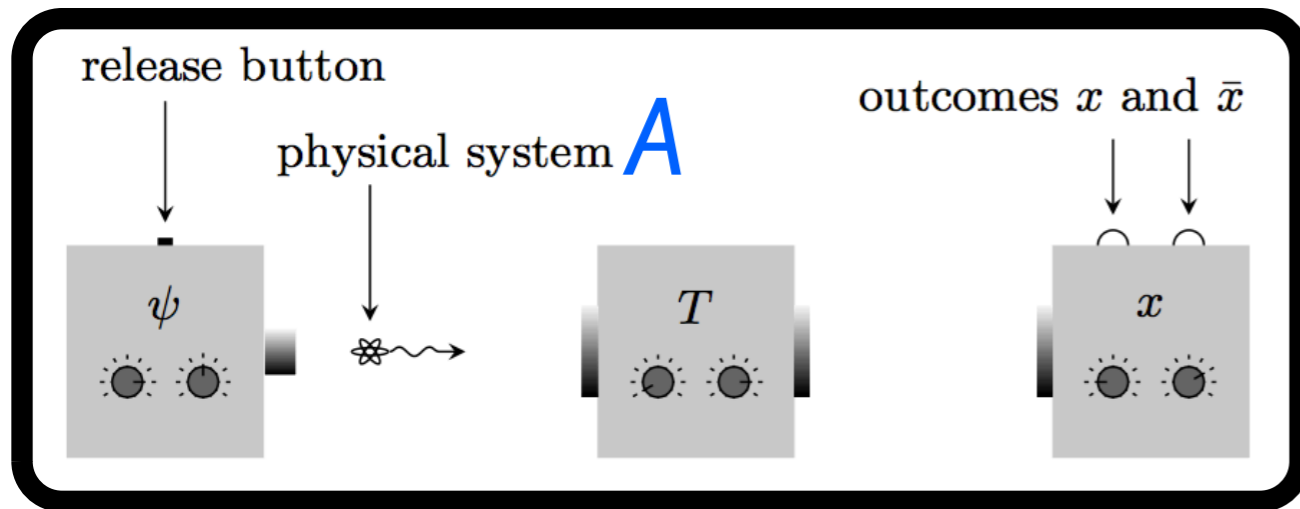


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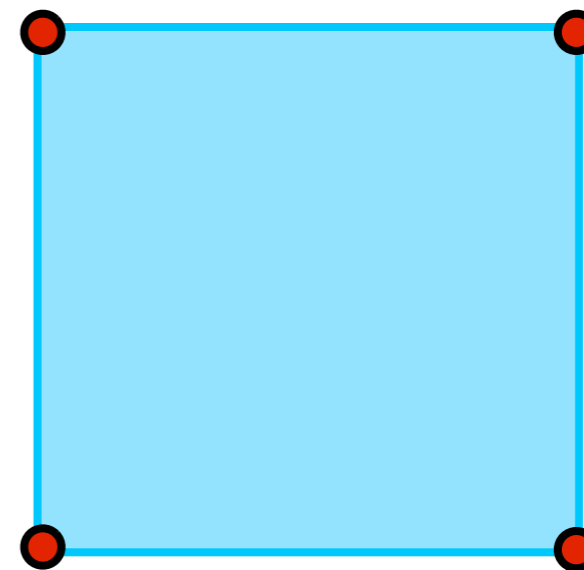
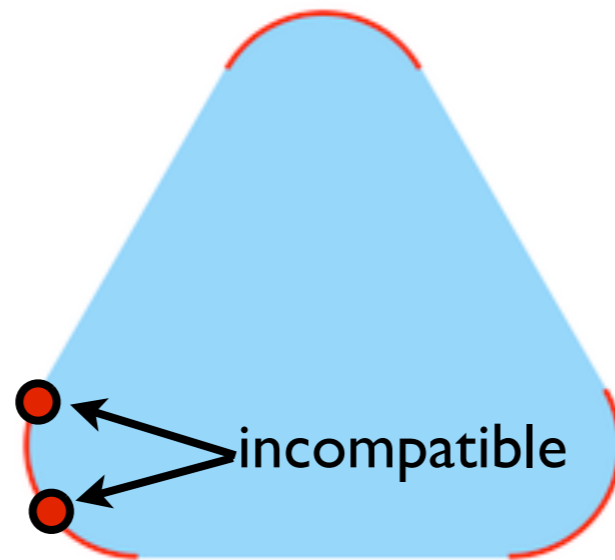
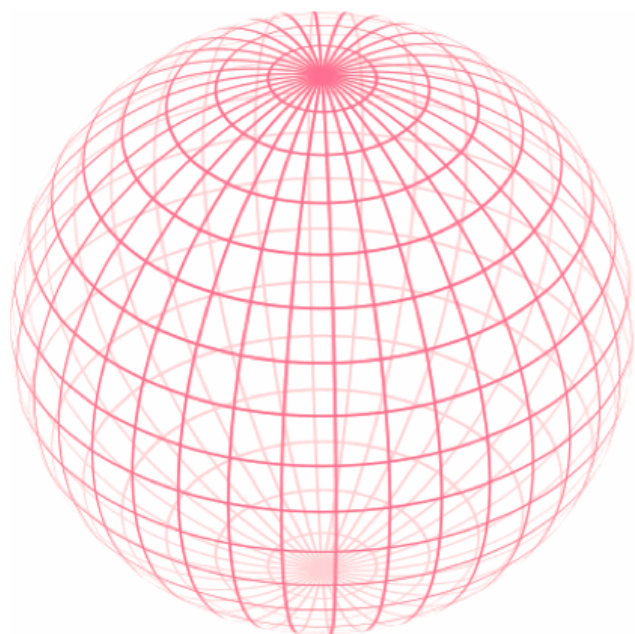


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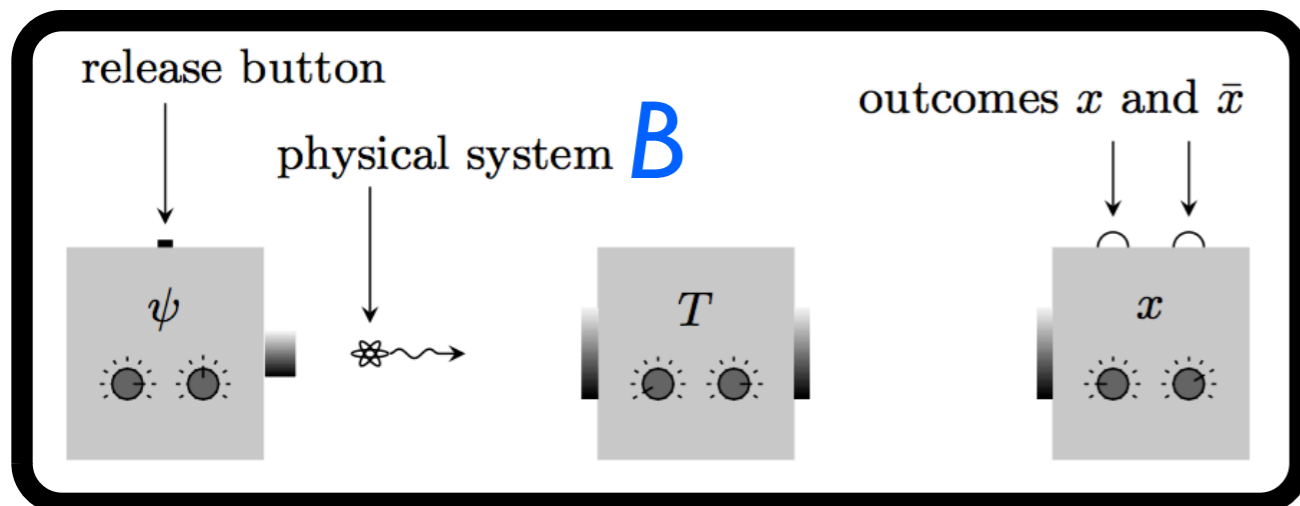
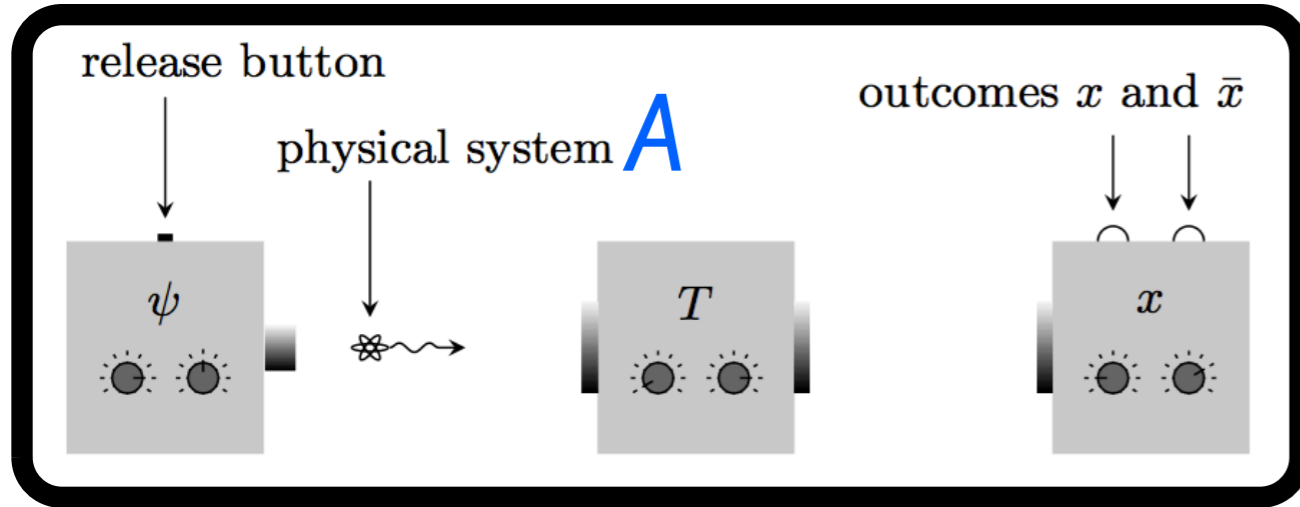


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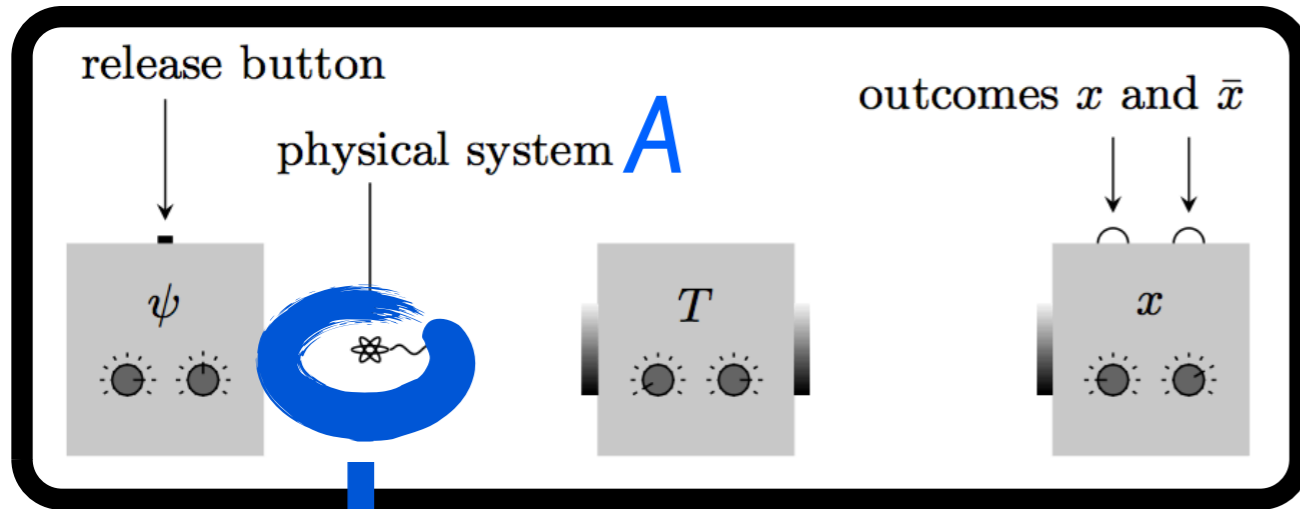
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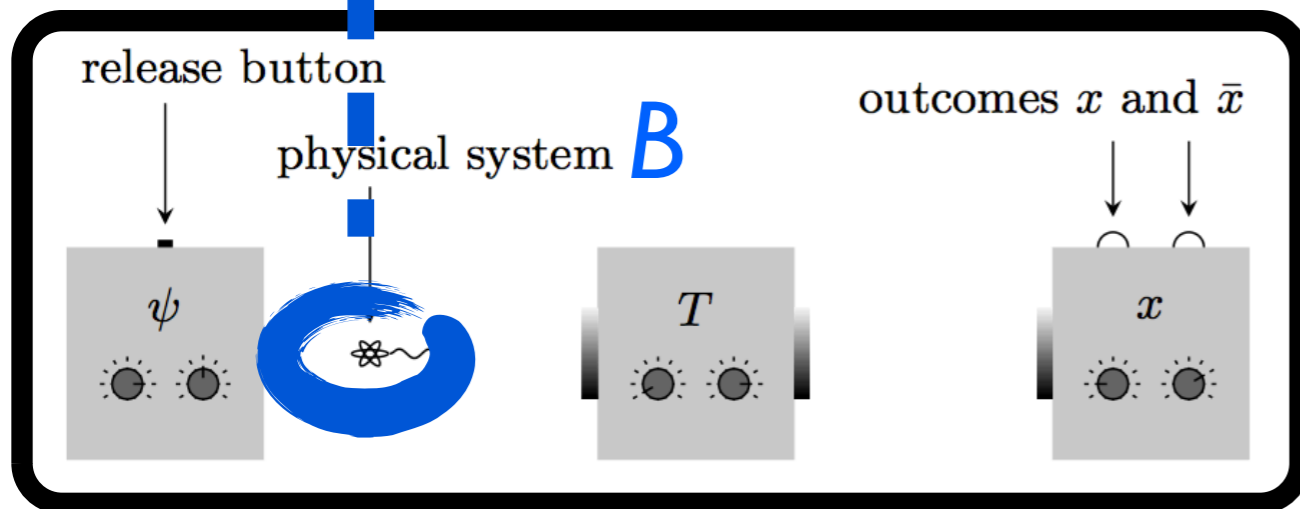
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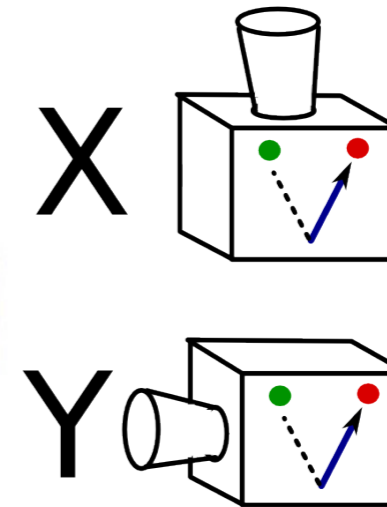
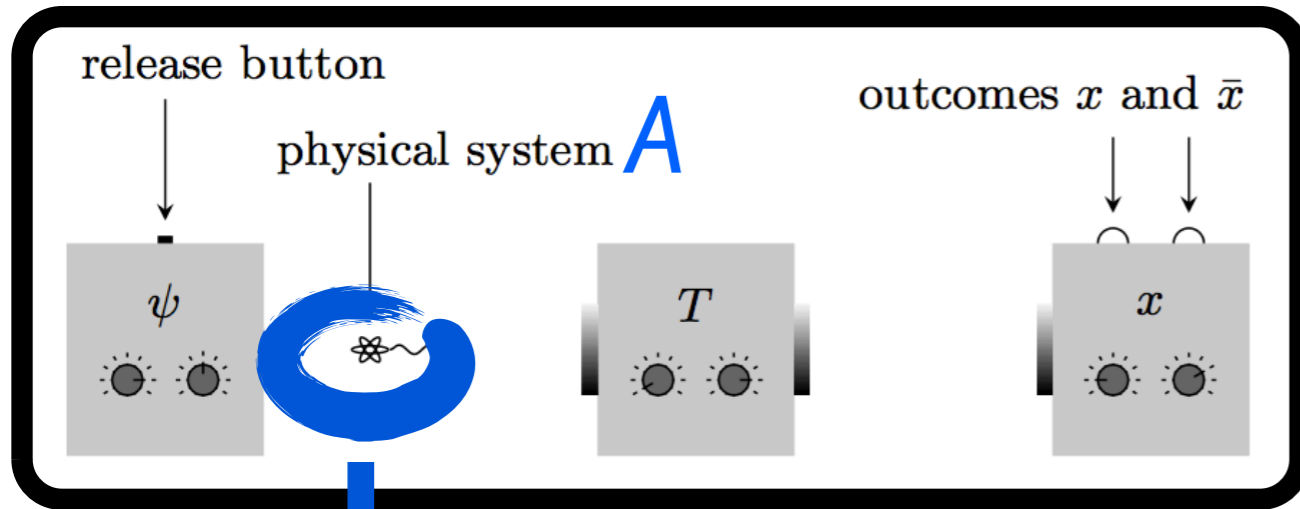
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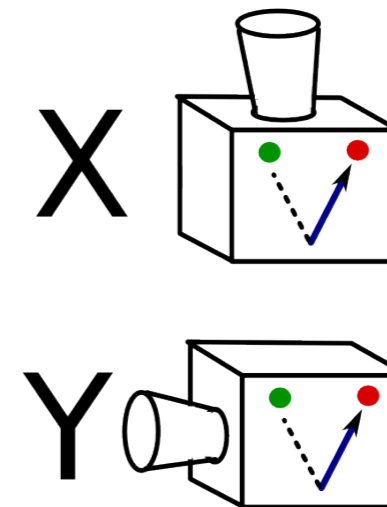
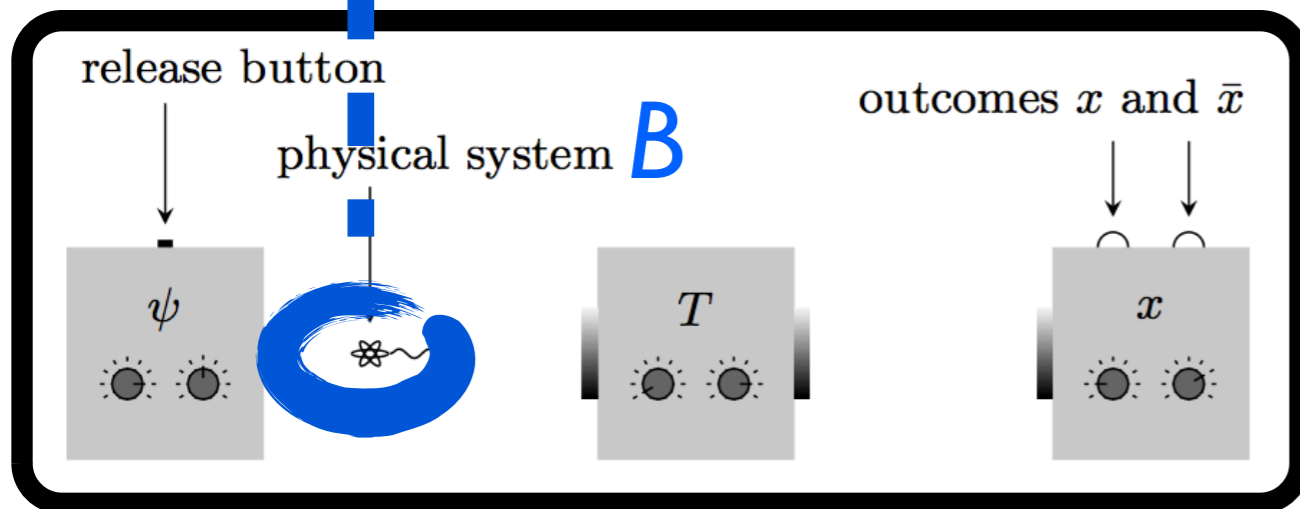
state on AB :
correlations



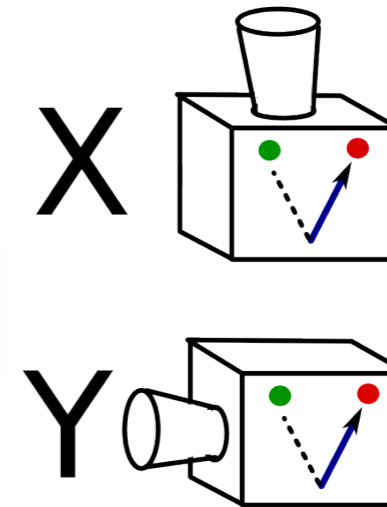
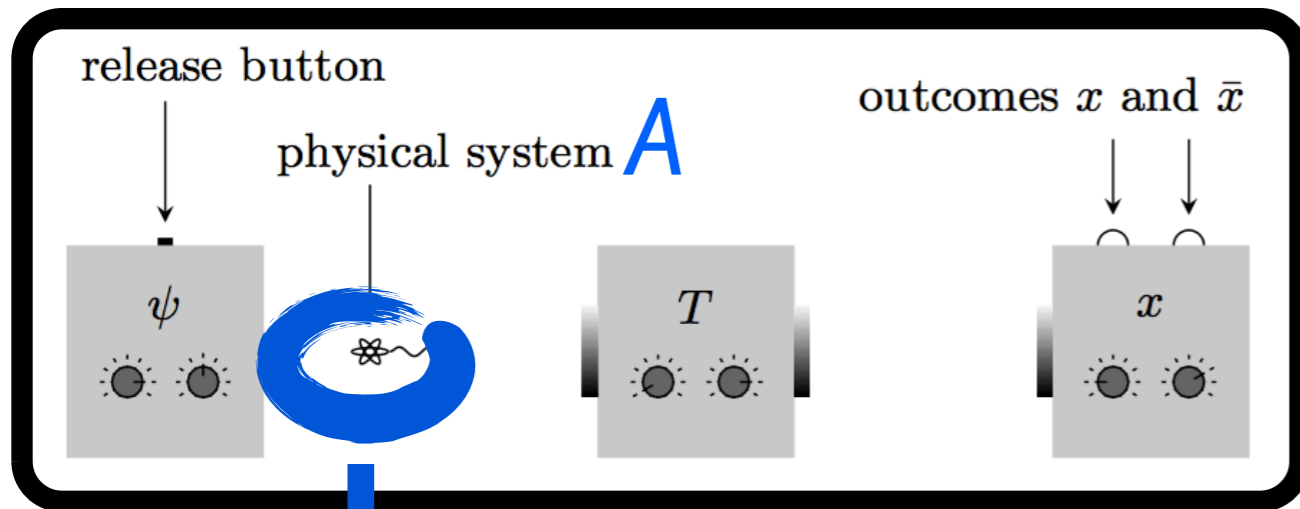
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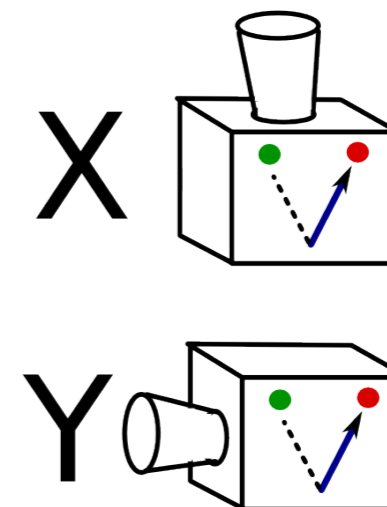
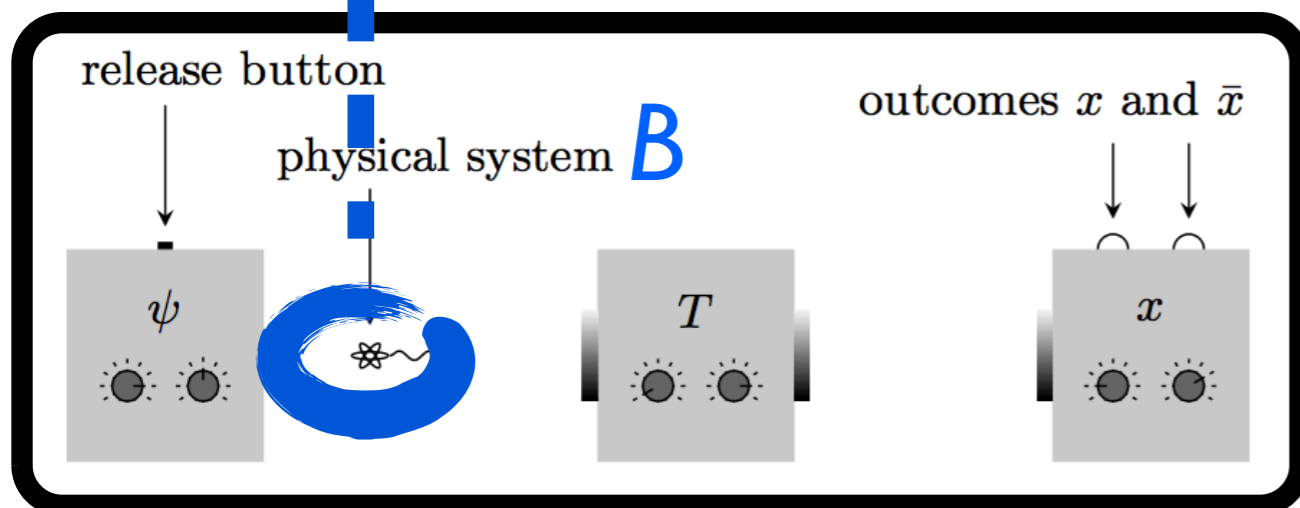


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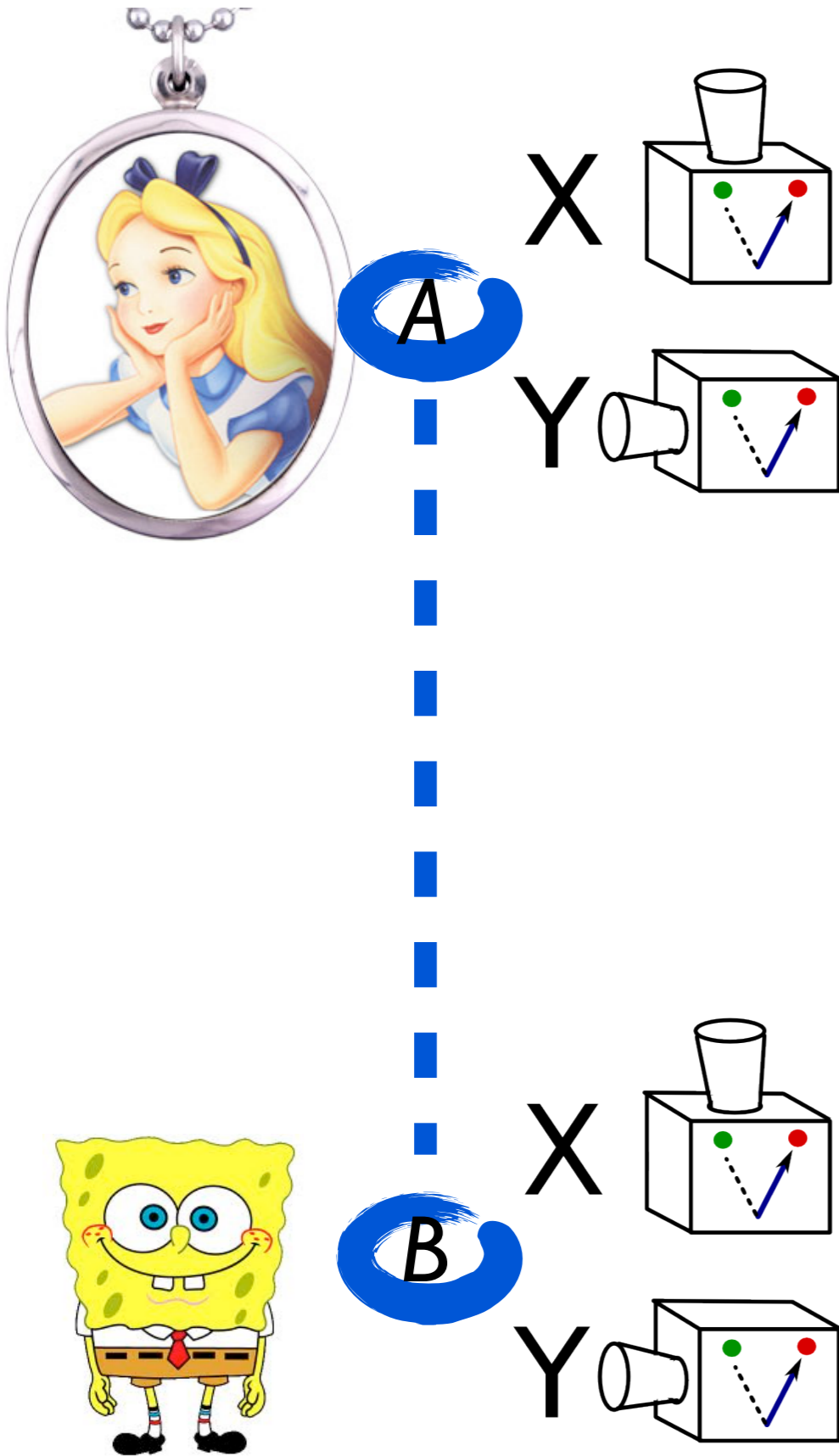


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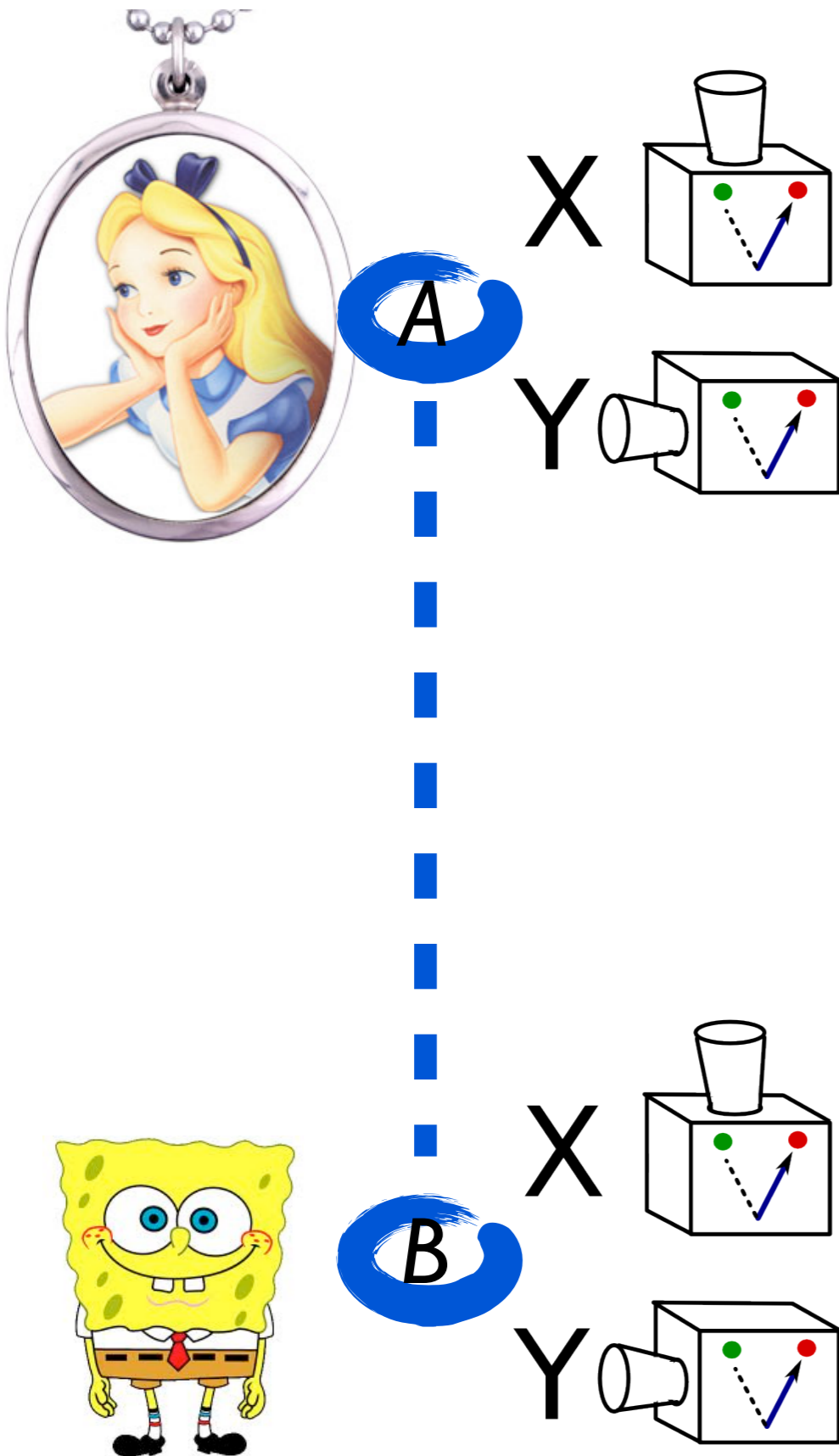
No-signalling condition:
Alice's probabilities **do not depend** on
Bob's choice of measurement.



2. General Probabilistic Theories

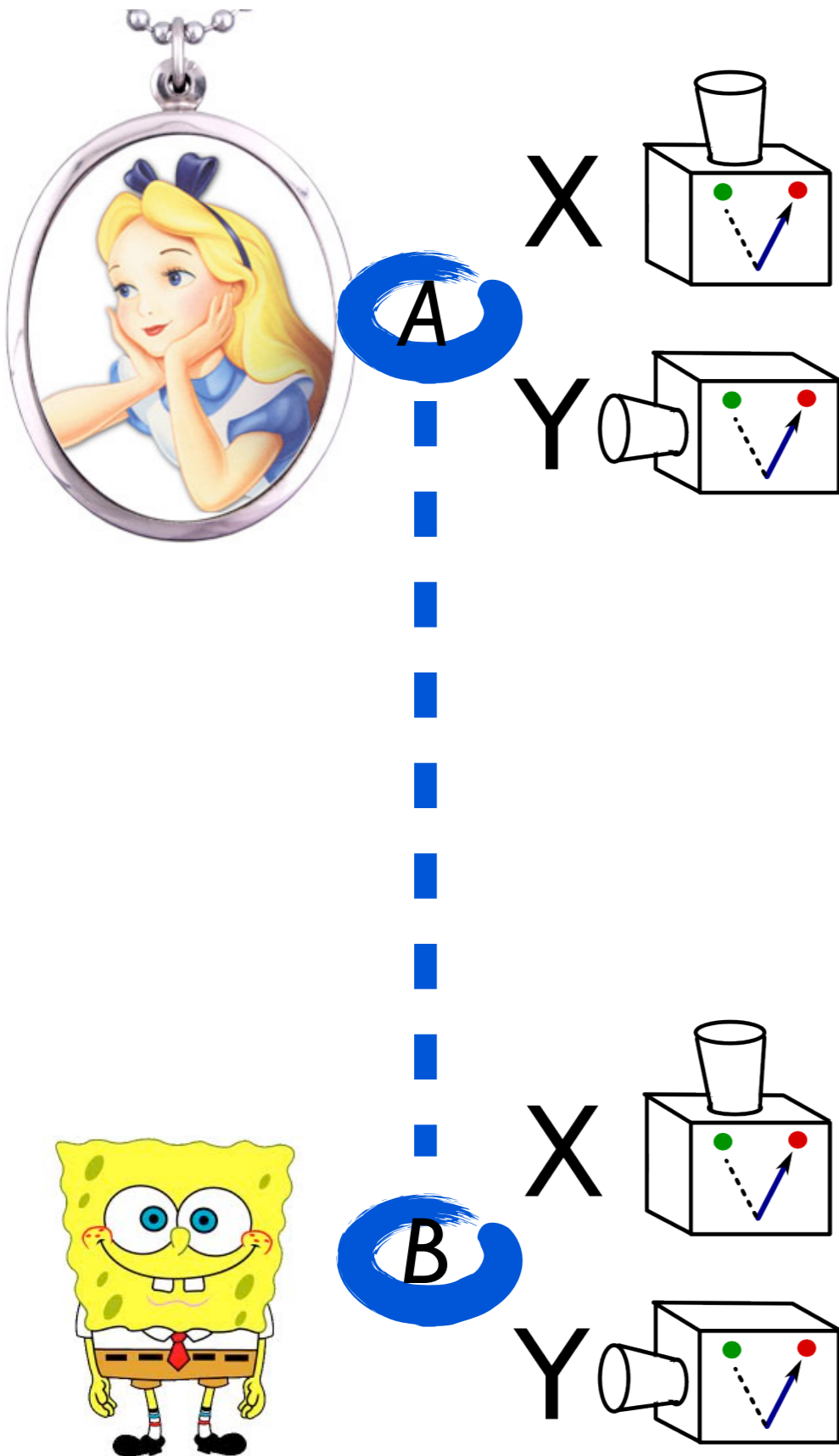


2. General Probabilistic Theories



Axiom I: States on AB are uniquely determined by correlations of local measurements on A, B .

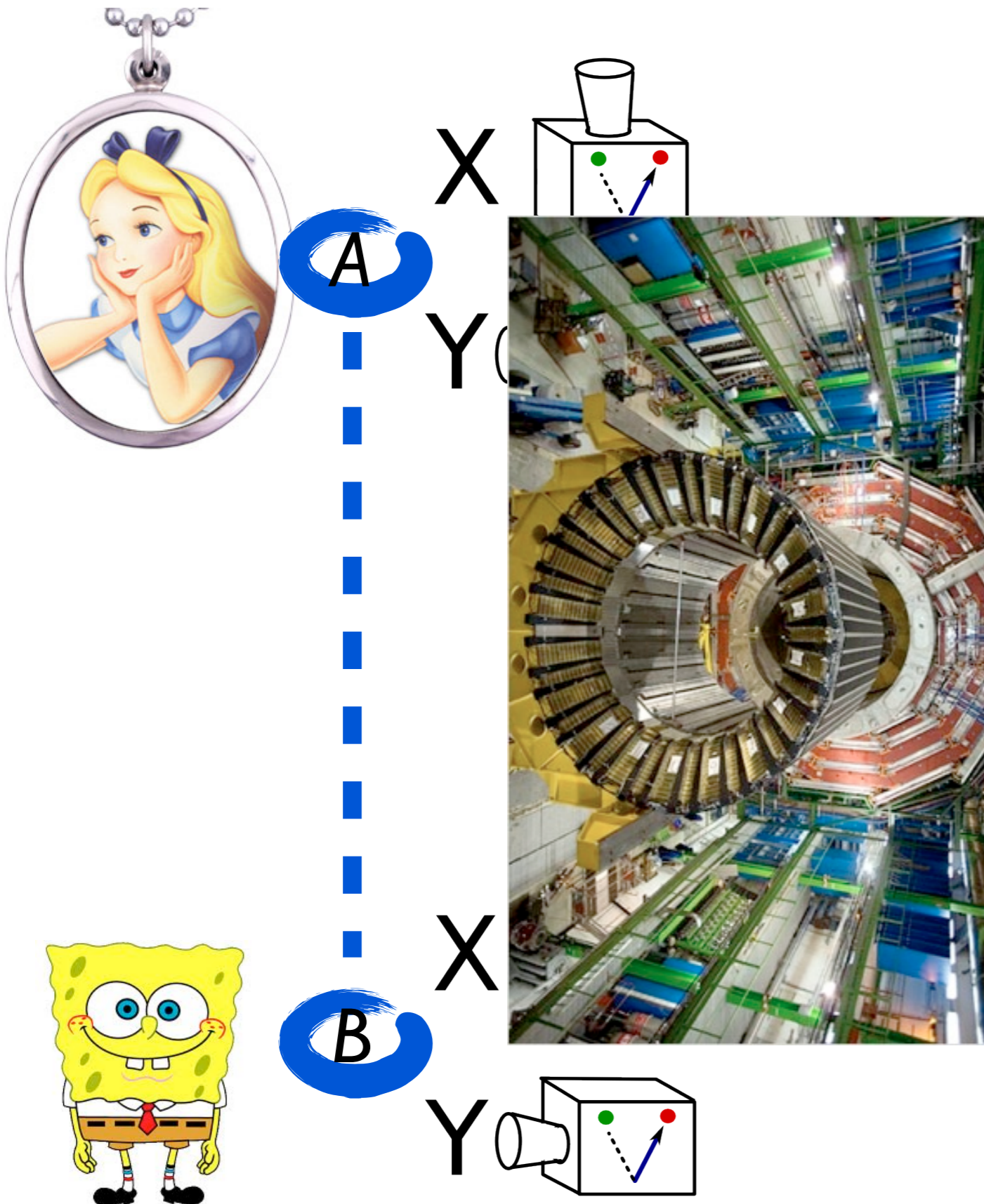
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No non-local measurements necessary.

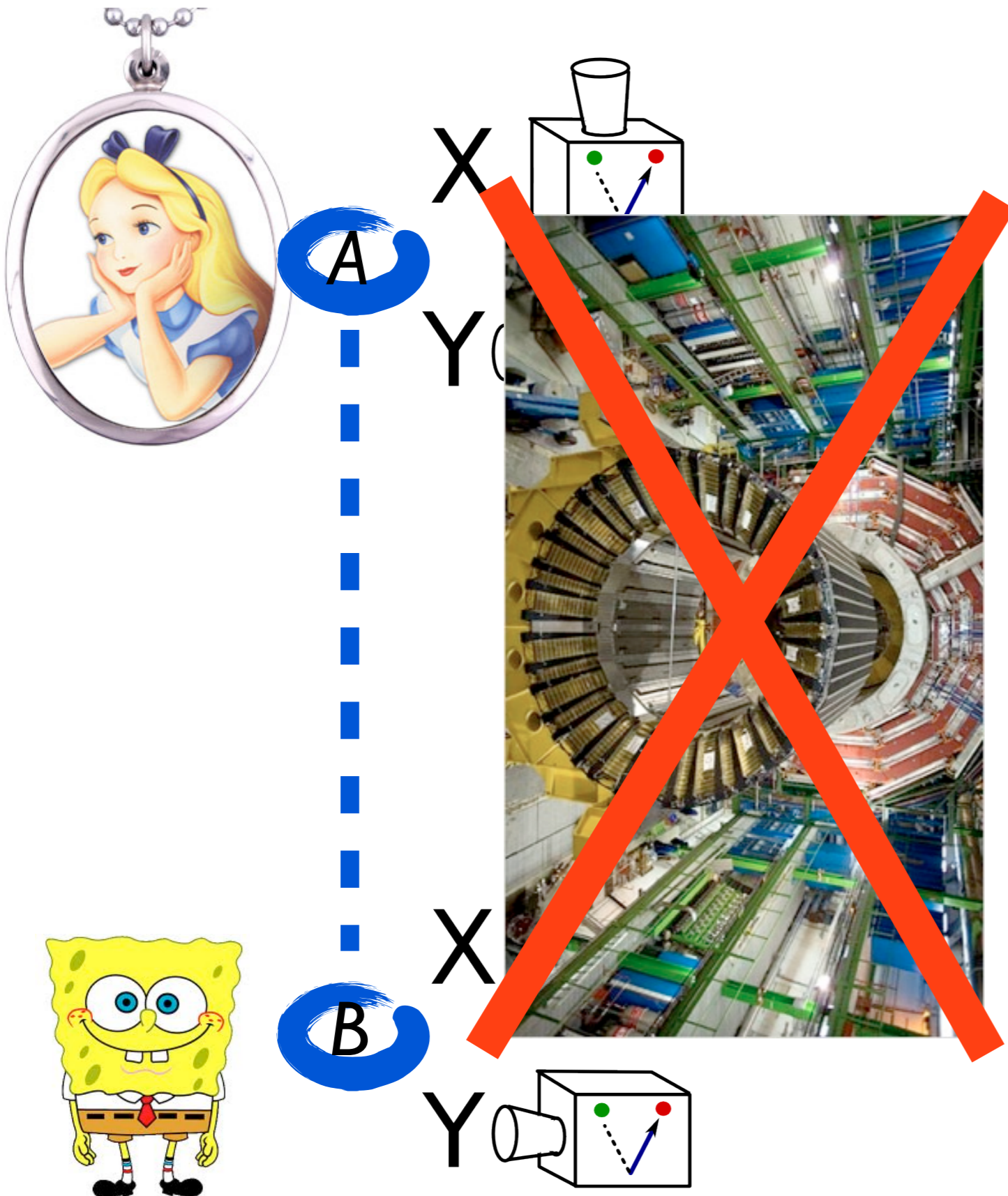
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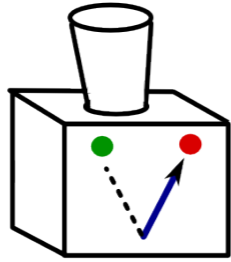
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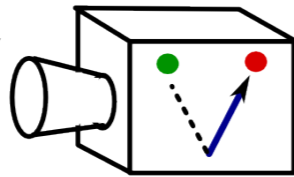


A

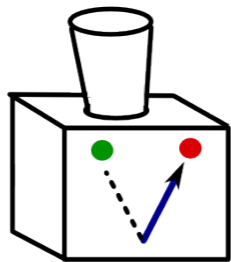
X



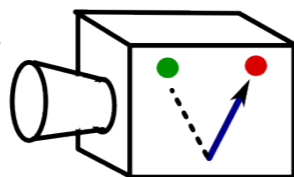
Y



X



Y



B

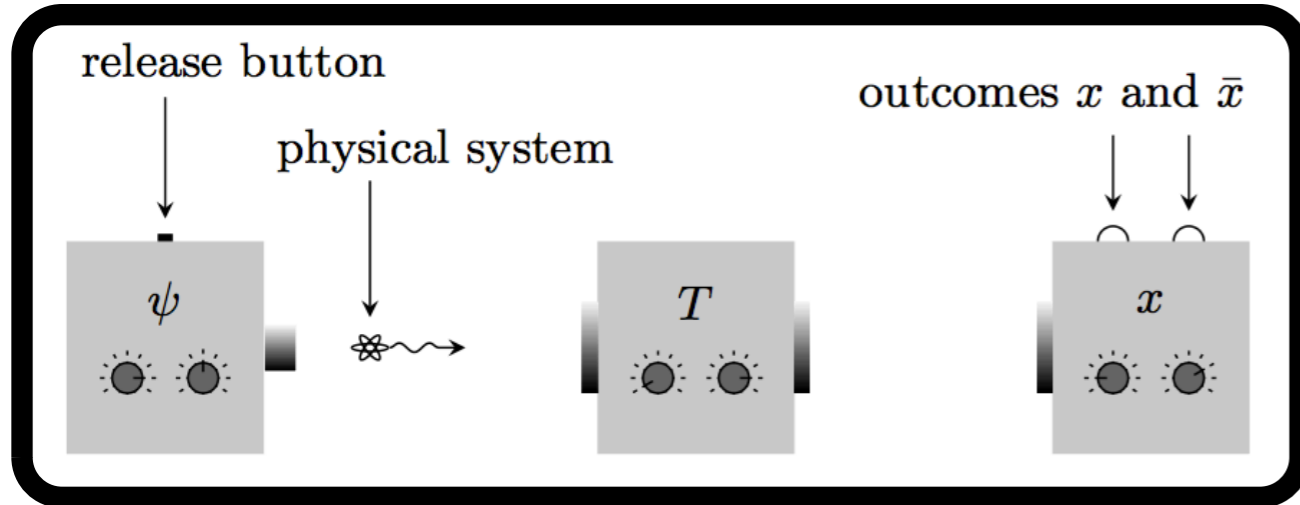


Axiom I: States on AB are uniquely determined by correlations of local measurements on A,B.

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No non-local measurements necessary.

Global state space $\Omega_{AB} \subset A \otimes B$ but not uniquely fixed!

Basic physical / operational assumptions

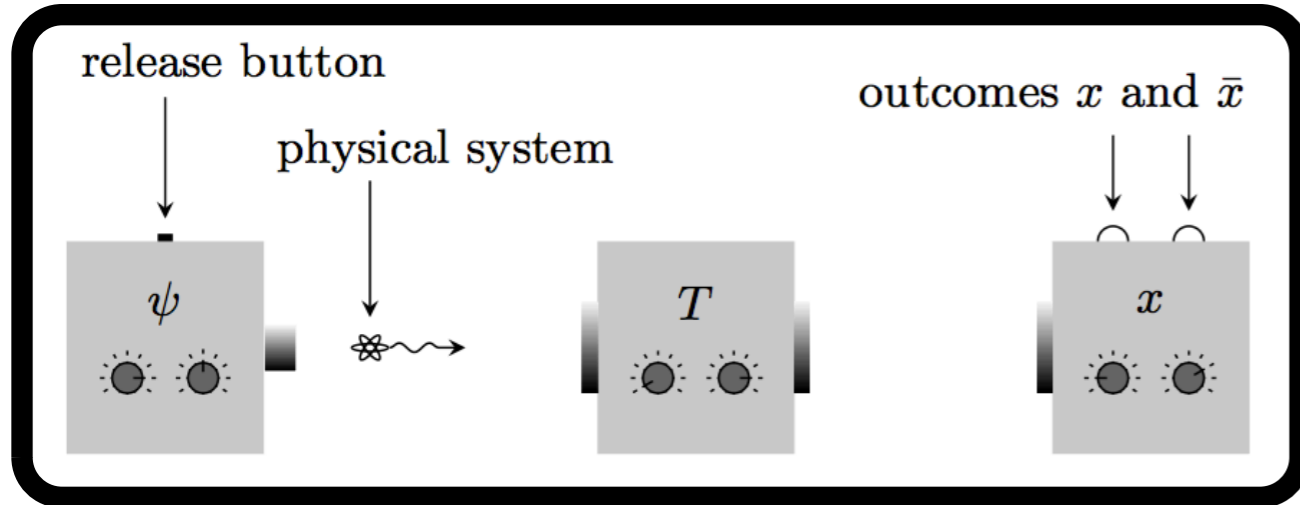
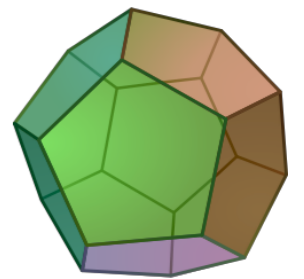


- States, transformations, and measurements with **outcome probabilities**.
- Combined systems: **no-signalling**.

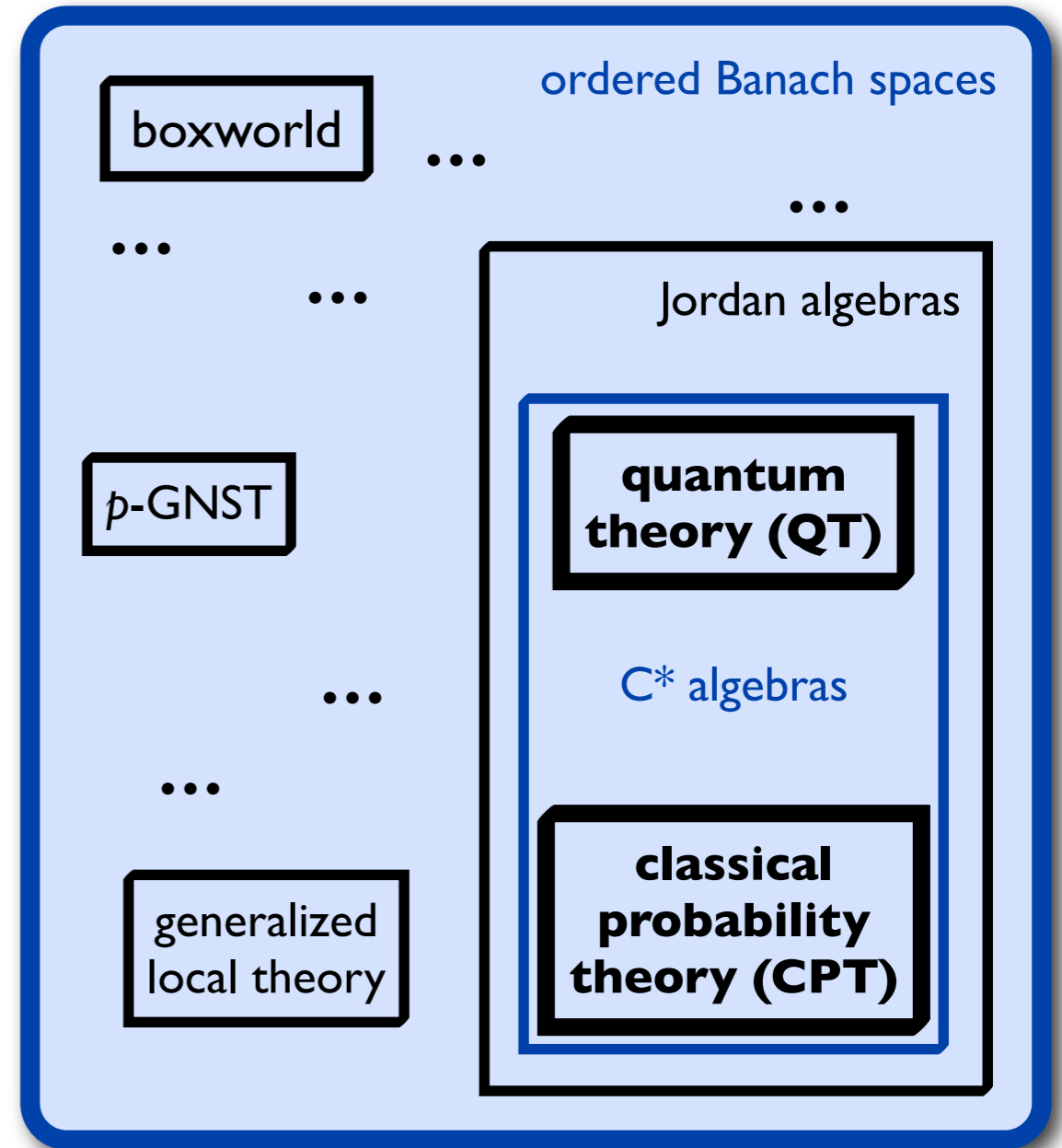
Basic physical / operational assumptions



General probabilistic theories



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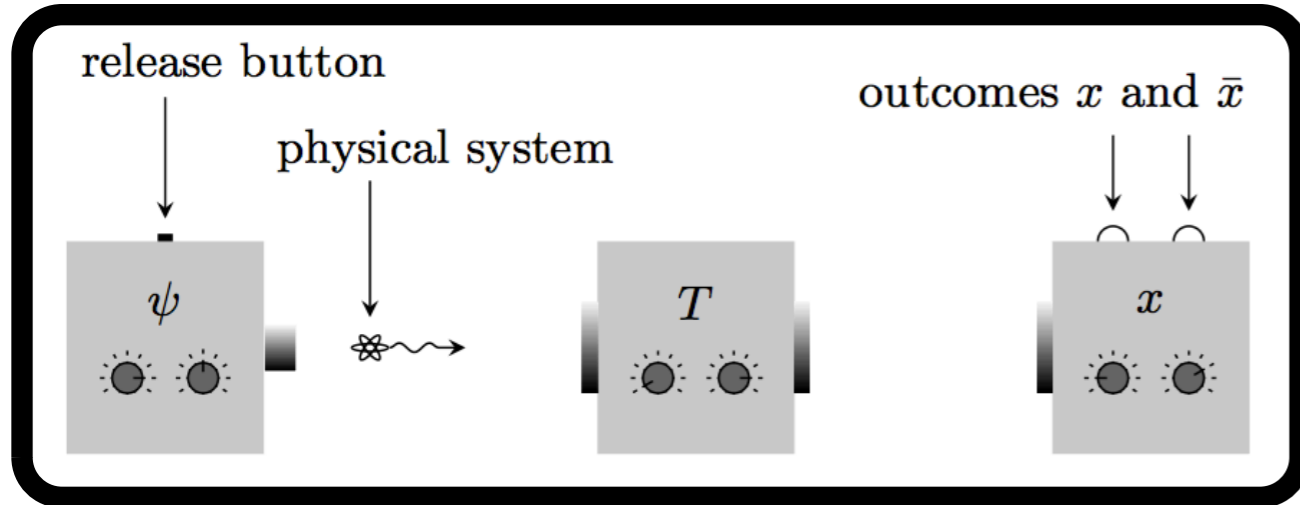
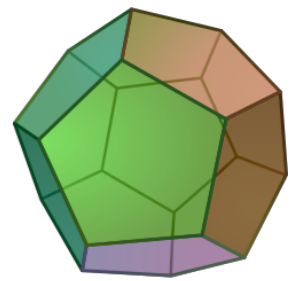


- **No** Hilbert spaces, complex numbers,...
- State spaces: **arbitrary convex sets**.
- Many ways to **combine systems**.

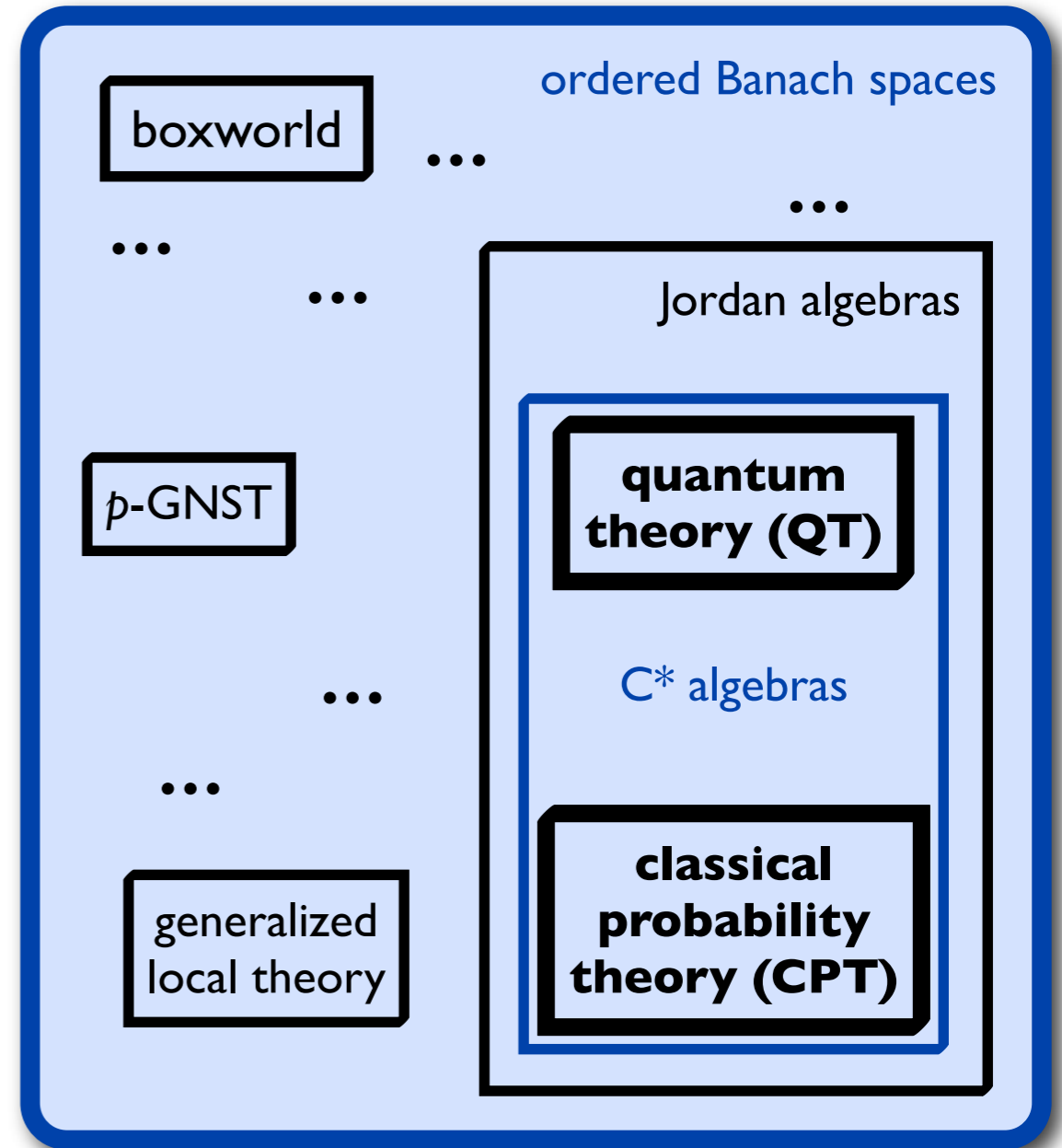
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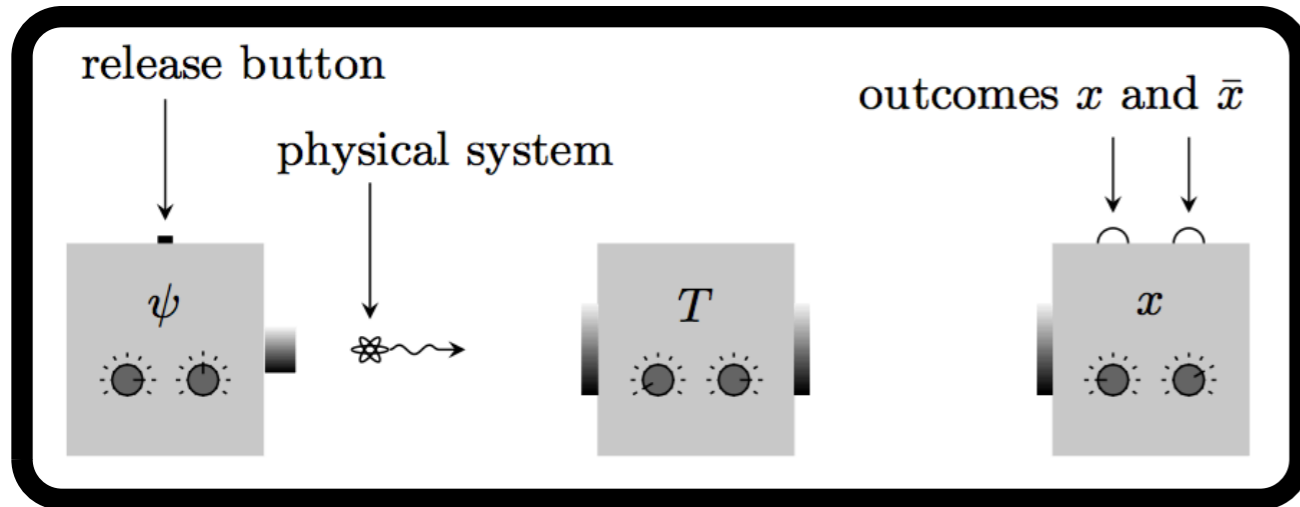
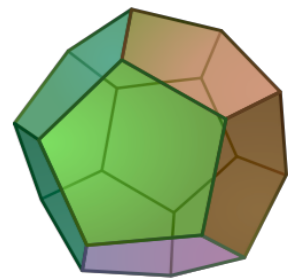
The Axioms:

- I. Local tomography
- II. Reversibility
- III. Subspace axiom
- IV. Finite-dimensionality
- V. All measurements allowed

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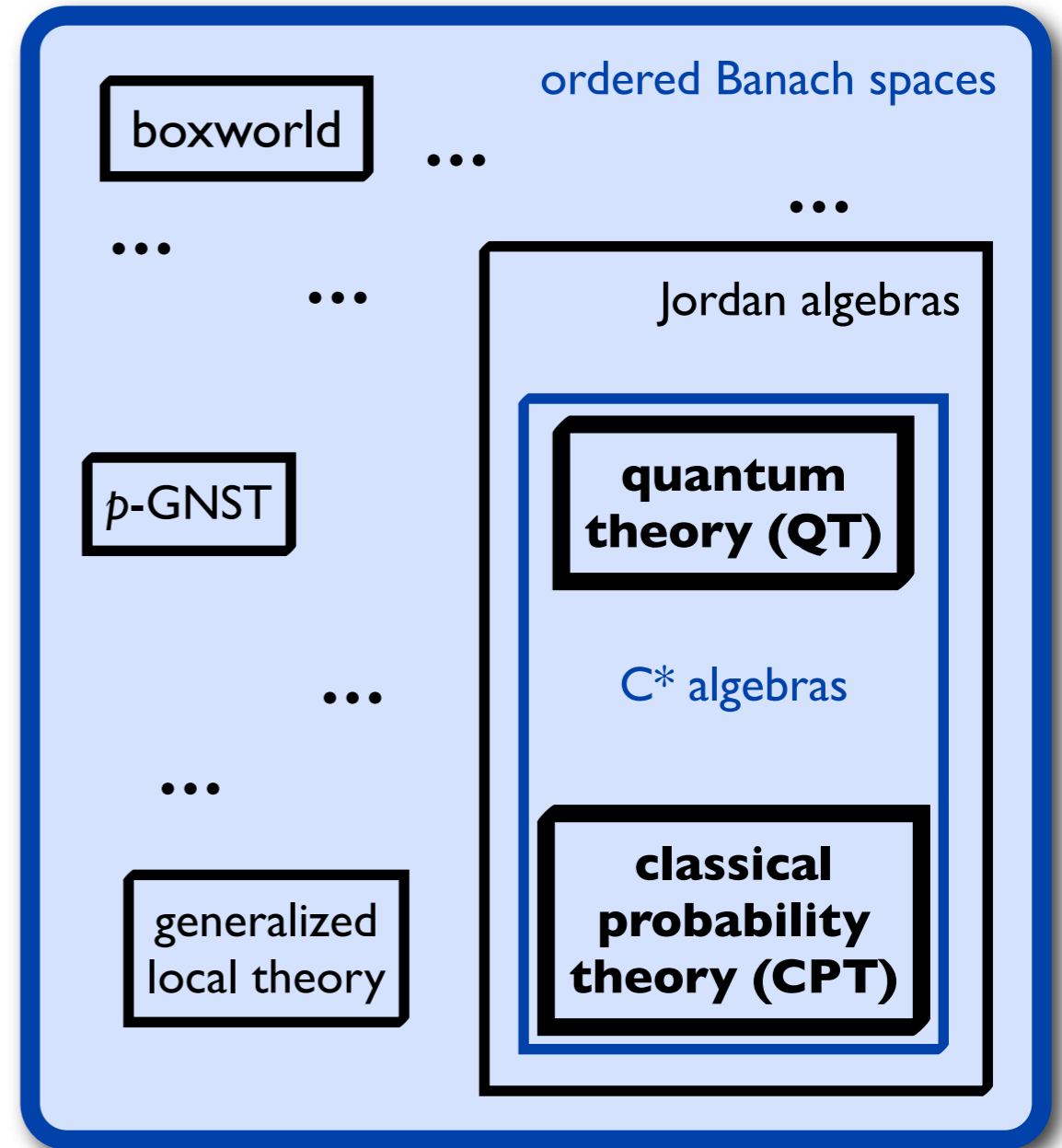
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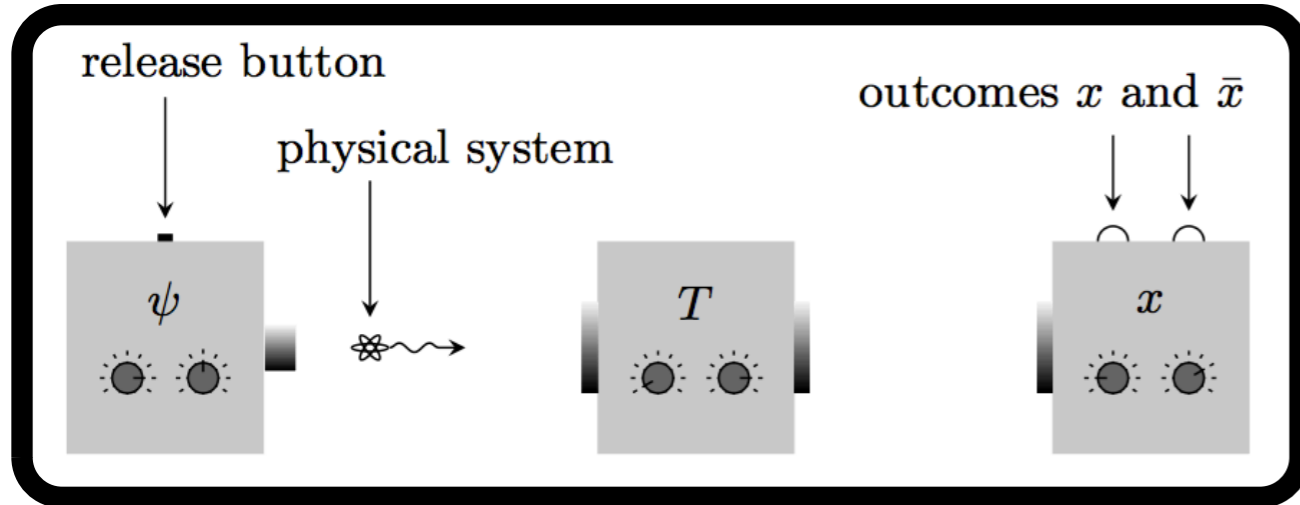
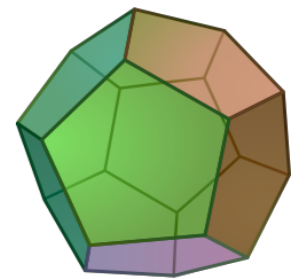


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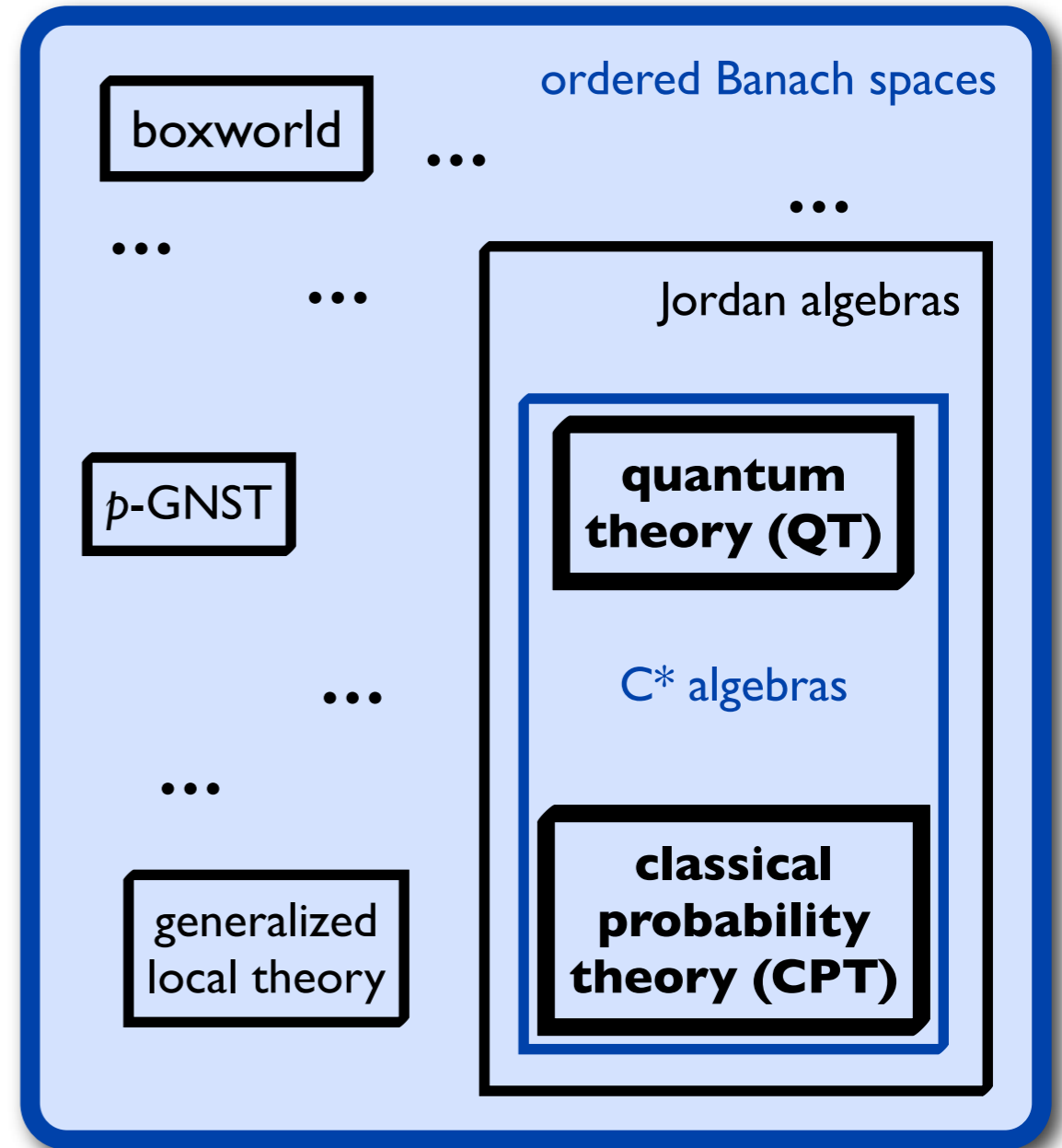
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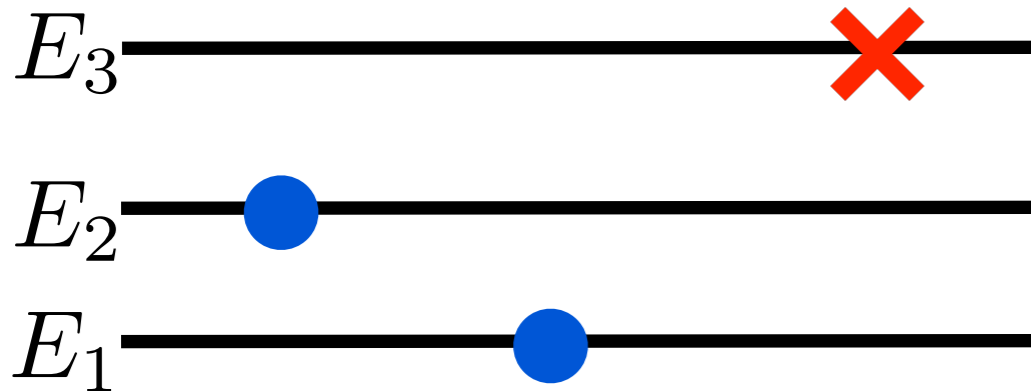
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3. The Subspace Axiom

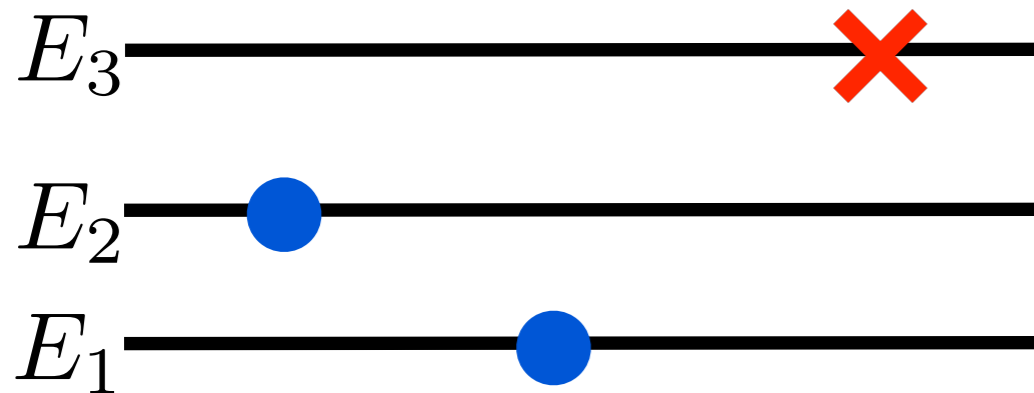
Some 3-level system:



Impossible to have system in 3rd level
 \Rightarrow find particle there with **probab. 0**

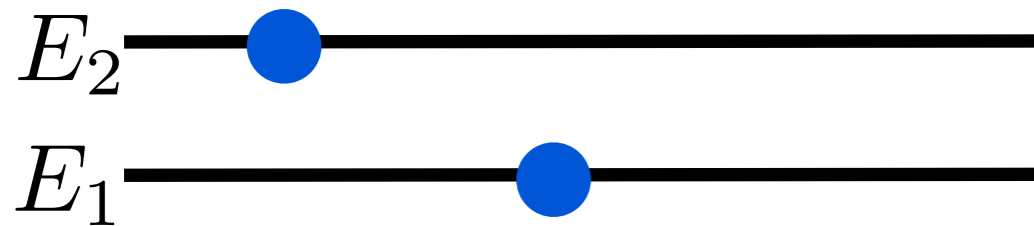
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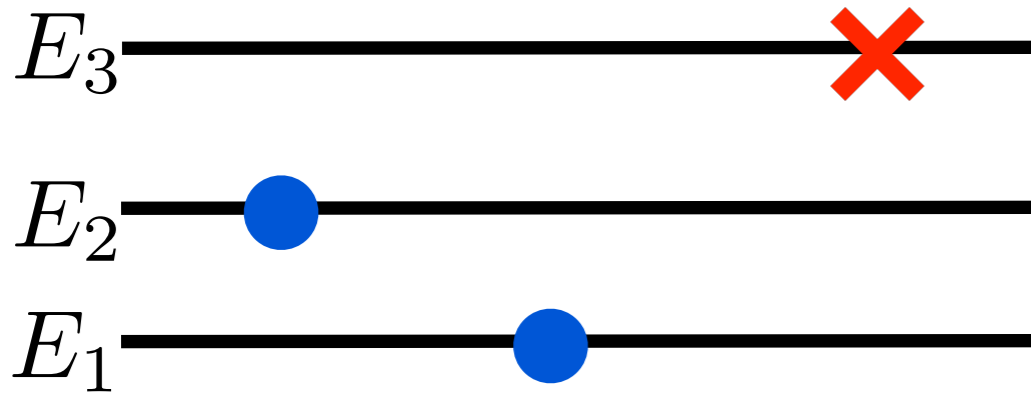
=



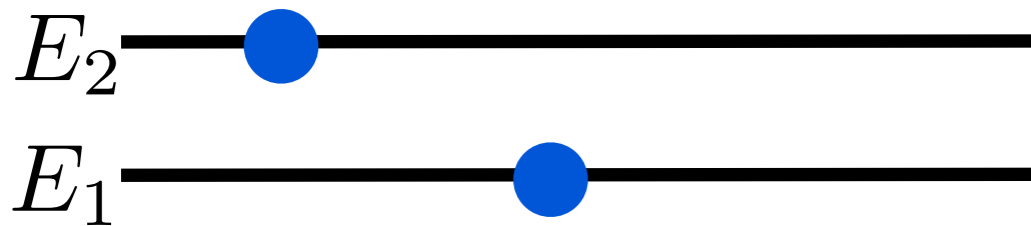
2-level system.

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=



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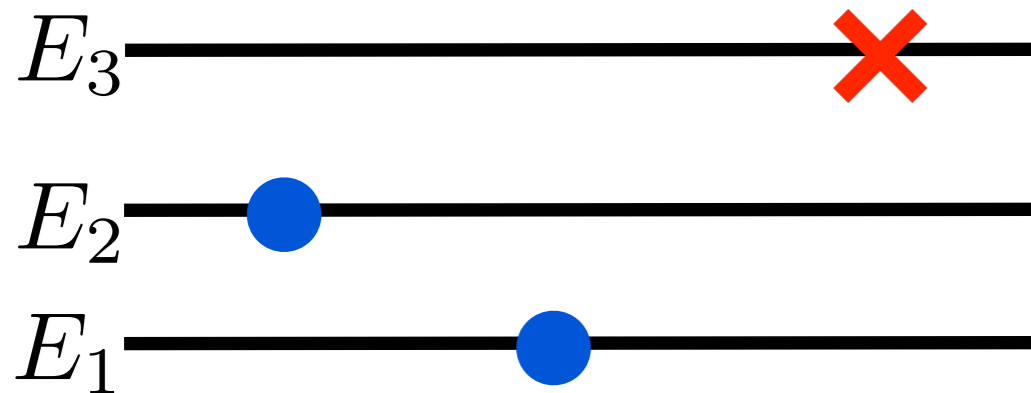
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QT: $\rho^{(3)} = \begin{pmatrix} \bullet & \bullet & 0 \\ \bullet & \bullet & 0 \\ 0 & 0 & 0 \end{pmatrix} \longrightarrow \rho^{(2)} = \begin{pmatrix} \bullet & \bullet \\ \bullet & \bullet \end{pmatrix}$

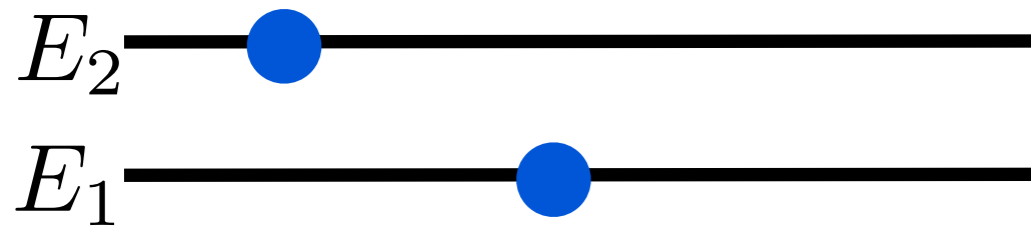
CPT: $P^{(3)} = (P_1, P_2, 0) \longrightarrow P^{(2)} = (P_1, P_2)$

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CPT: $P^{(3)} = (P_1, P_2, 0) \longrightarrow P^{(2)} = (P_1, P_2)$

Otherwise, physics would be affected by **impossible potentialities**.



3. The Subspace Axiom

Axiom III: Let Ω_N and Ω_{N-1} be systems with capacities N and $N-1$. If (E_1, \dots, E_N) is a complete measurement on Ω_N , then the set of states ω with $E_N(\omega) = 0$ is equivalent to Ω_{N-1} .

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Capacity N of Ω = maximal # of perfectly distinguishable states.

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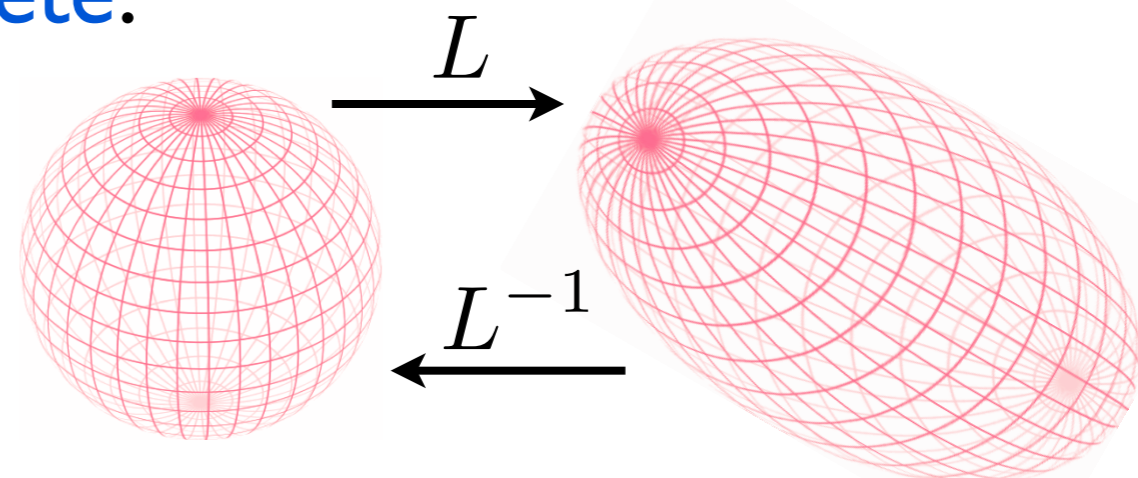
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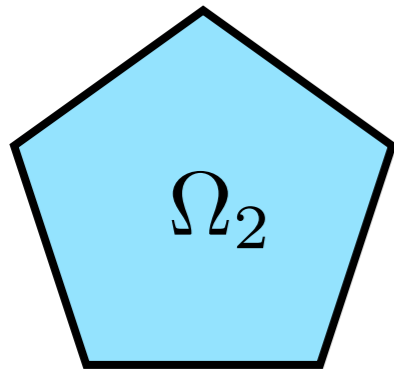
If $n = N$ then (E_1, \dots, E_n) is **complete**.

Equivalent = same state spaces up to a linear map (physically the same!)



4. Derivation of the Hilbert space formalism

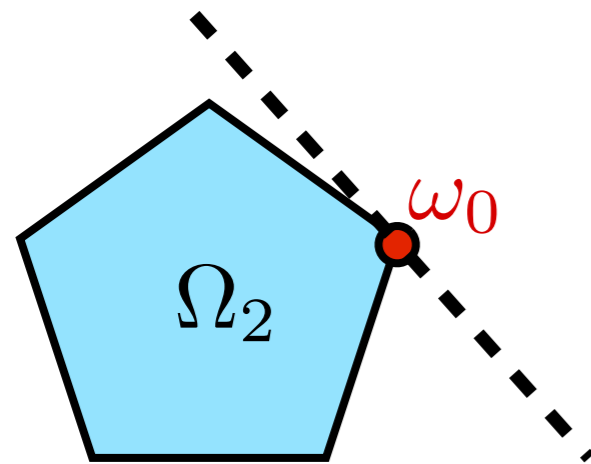
Why a **bit** is described by a **ball**:



capacity 2 (bit)

4. Derivation of the Hilbert space formalism

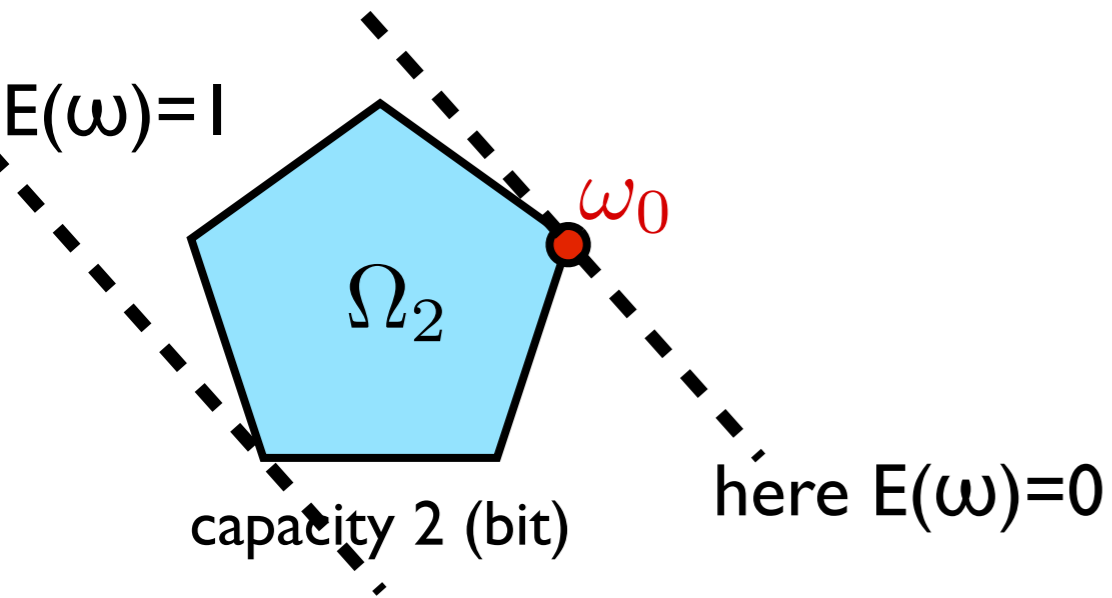
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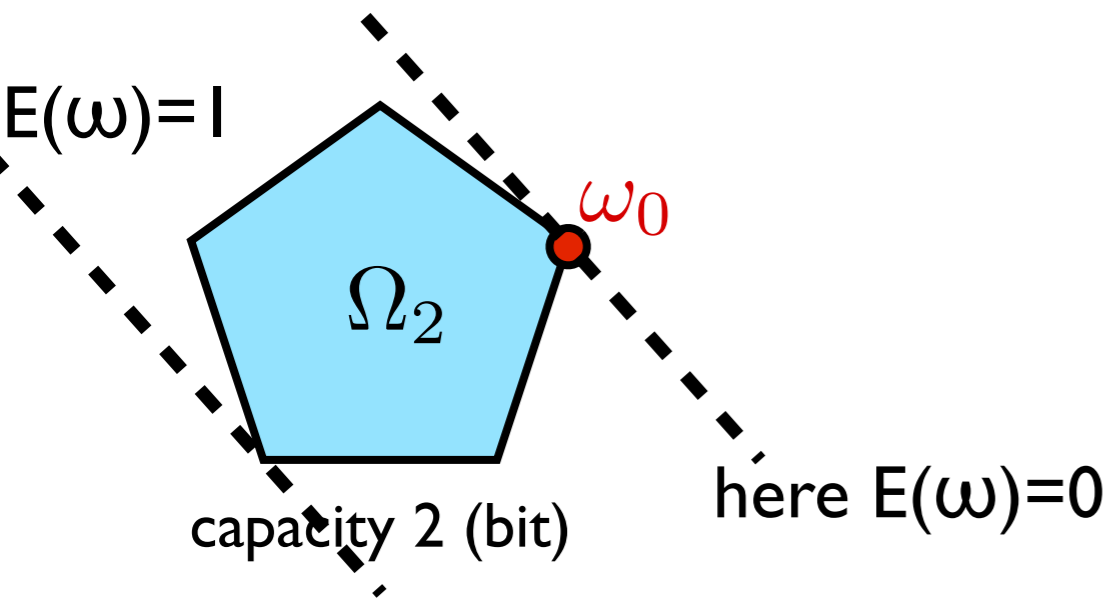


$(I-E, E)$ is complete measurement.

$$\Rightarrow \{\omega : E(\omega) = 0\} = \{\omega_0\} \sim \Omega_1.$$

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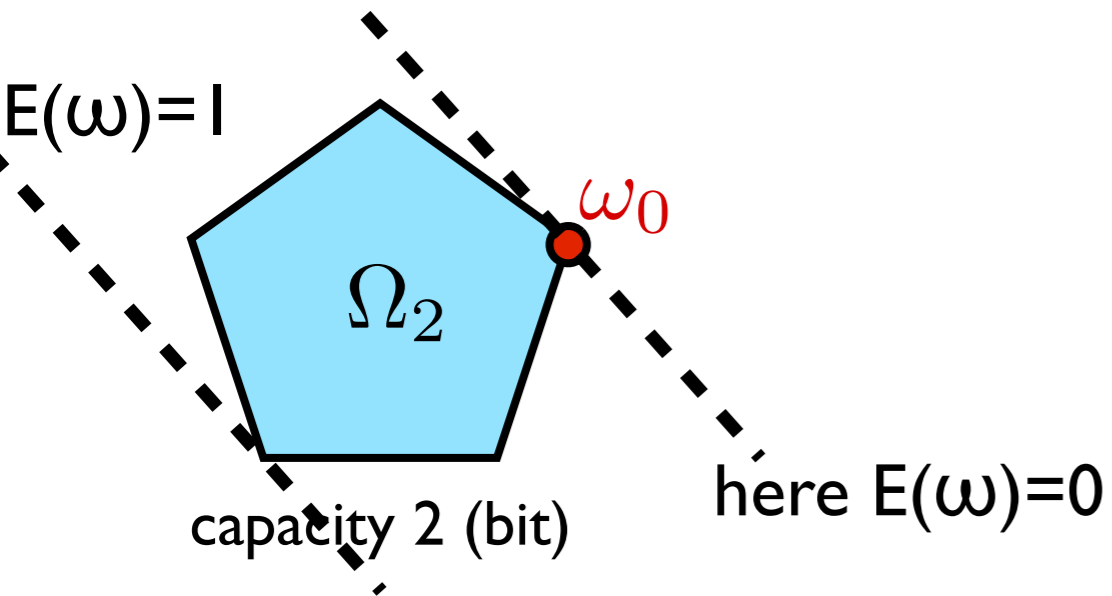
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$\Rightarrow \Omega_1$ contains a **single** state.

4. Derivation of the Hilbert space formalism

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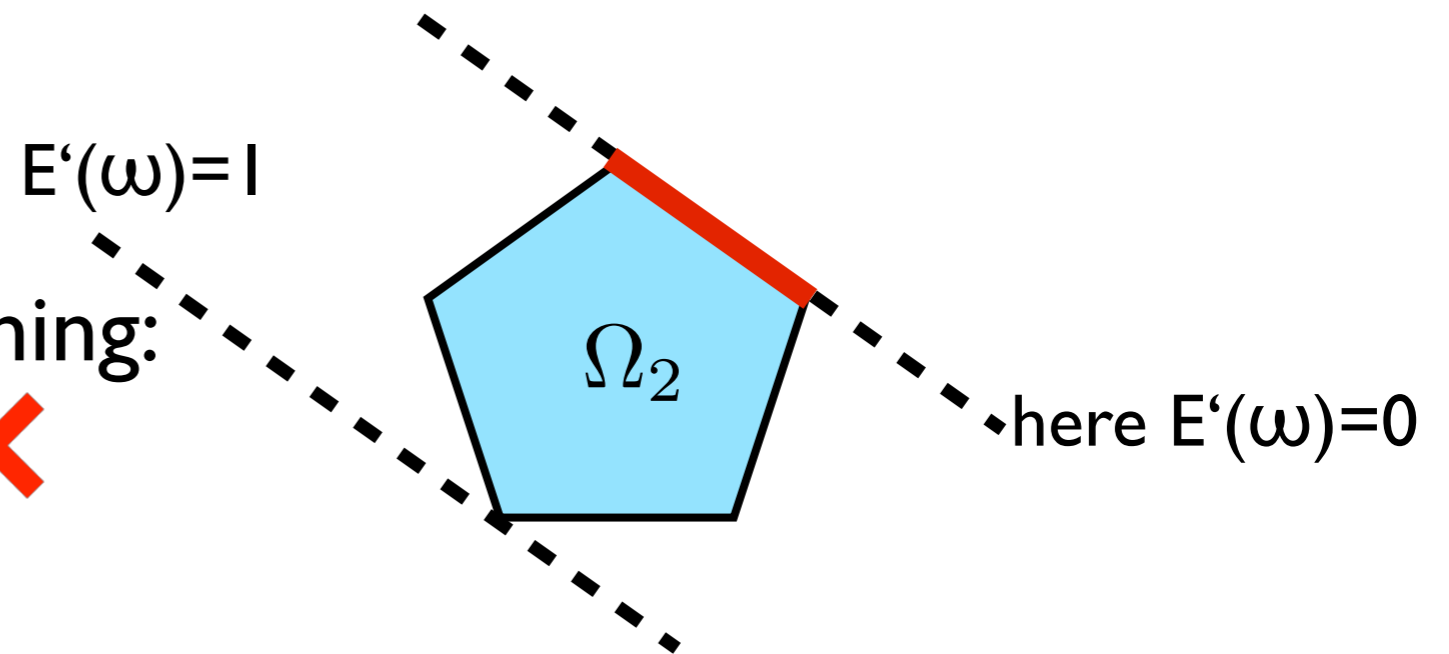


$(I-E, E)$ is complete measurement.

$$\Rightarrow \{\omega : E(\omega) = 0\} = \{\omega_0\} \sim \Omega_1.$$

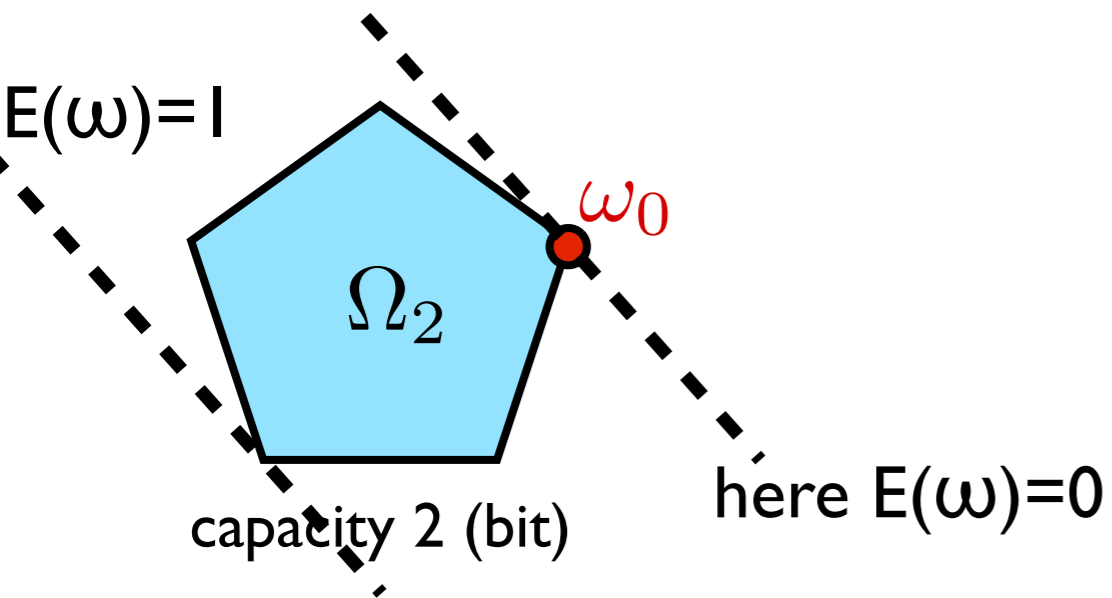
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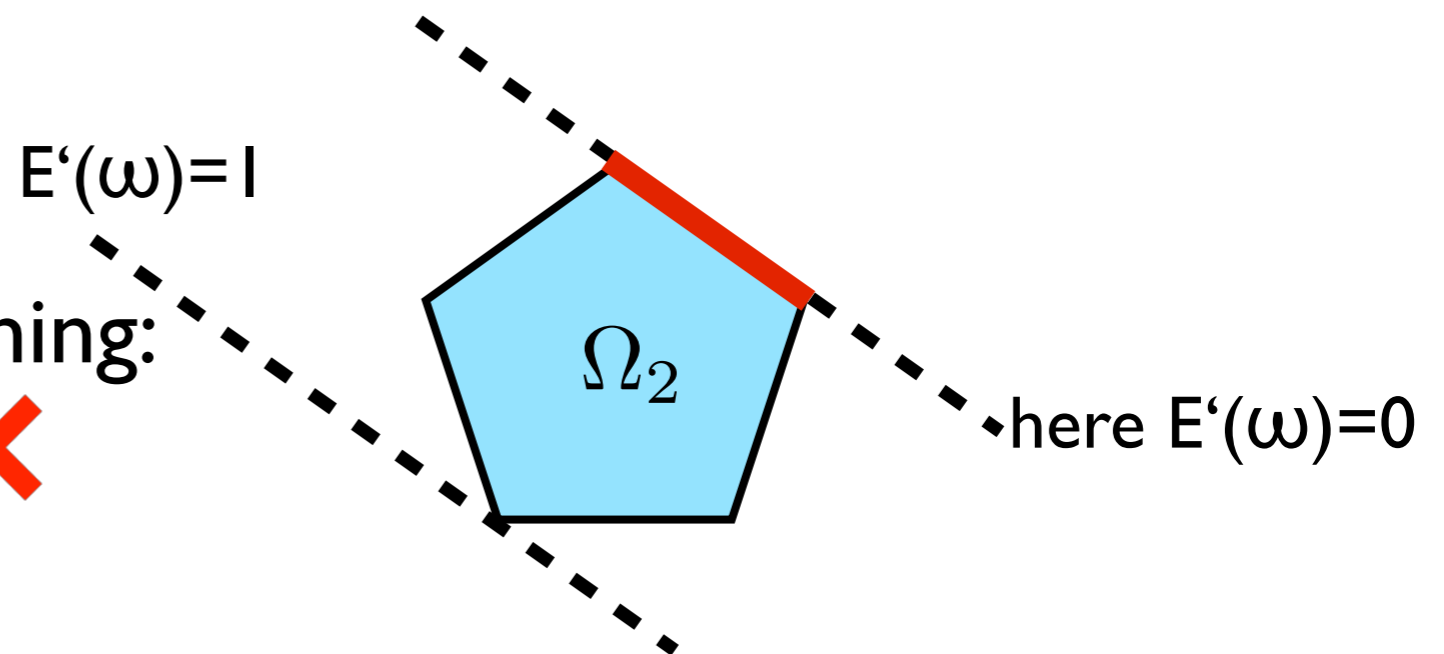
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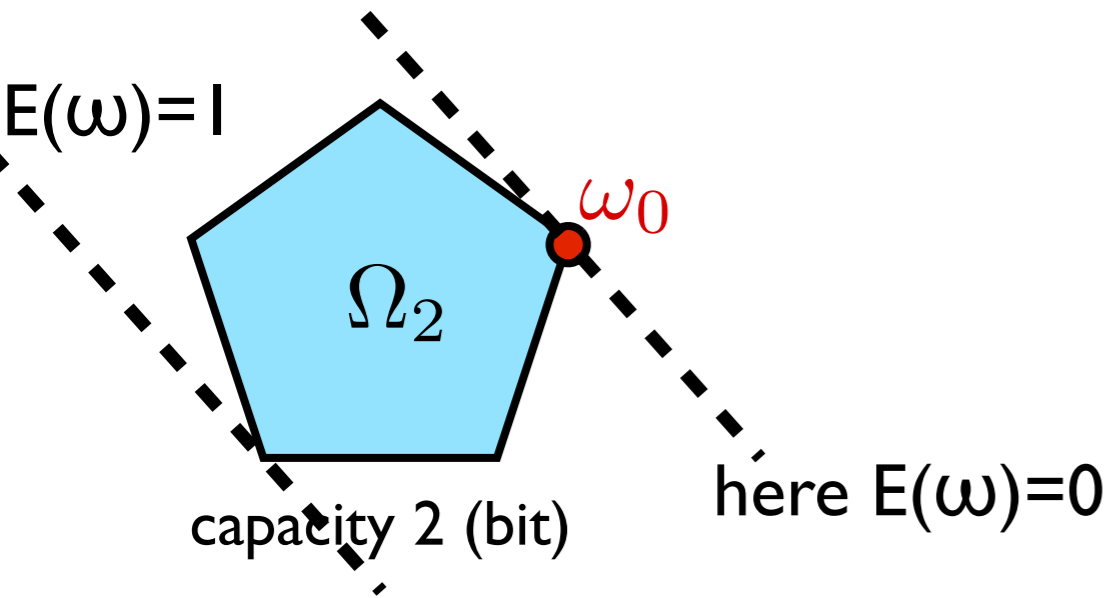


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4. Derivation of the Hilbert space formalism

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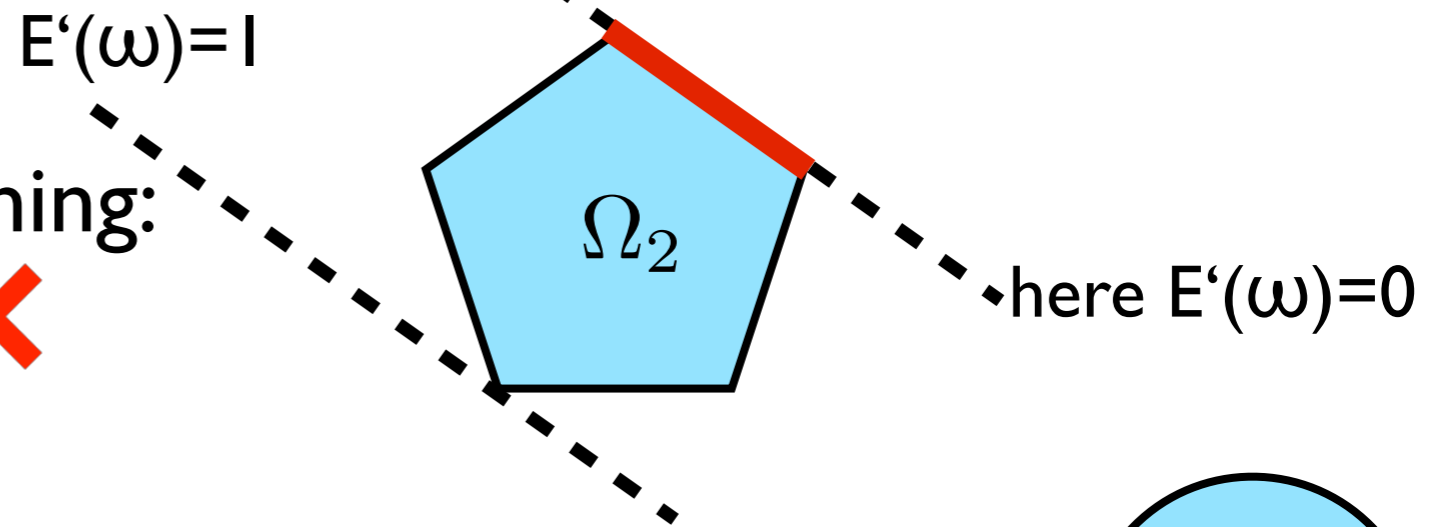
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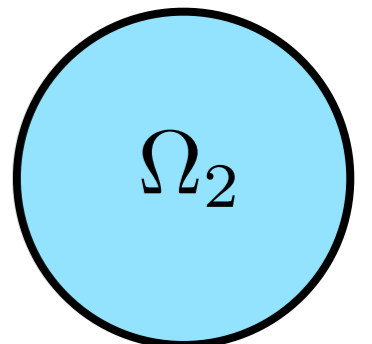
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Reversibility axiom $\Rightarrow \Omega_2$ is a ball.



4. Derivation of the Hilbert space formalism

Prove step by step (using the axioms):

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$$\mu = \frac{1}{N} \sum_{i=1}^N \omega_i,$$

- capacity $N_{AB} = N_A N_B$ and bit ball dimension

$$\dim(\Omega_2) = 2^r - 1 \in \{1, 3, 7, 15, 31, \dots\}.$$

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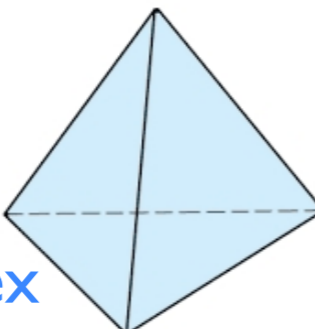
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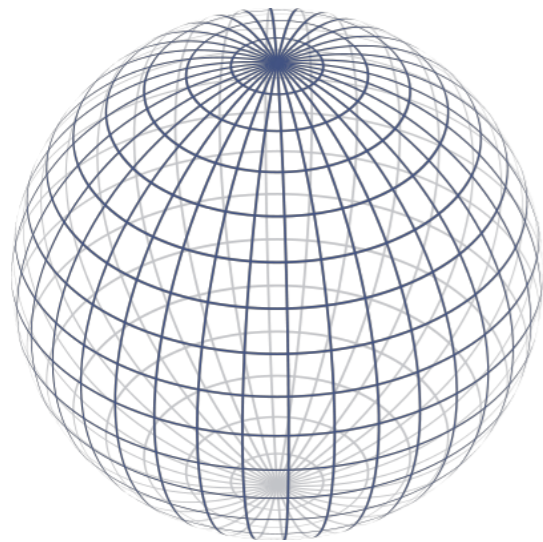
$\Omega_2 =$ 

If $\dim(\Omega_2) = 1$ then the theory is **CPT** (easy):

$\Omega_N =$  $\mathcal{G}_N =$ permutation group.

N-simplex

4. Derivation of the Hilbert space formalism



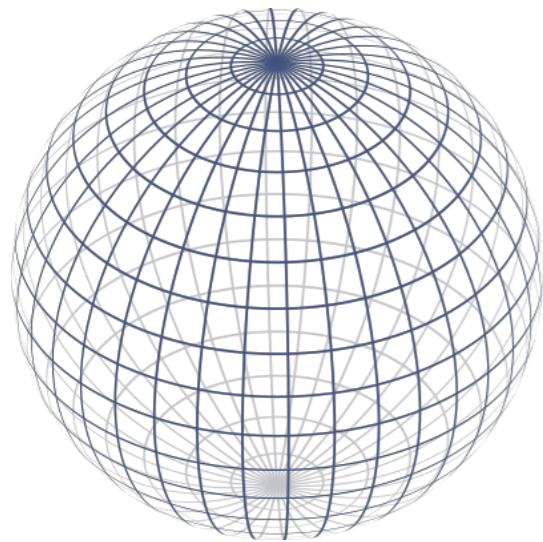
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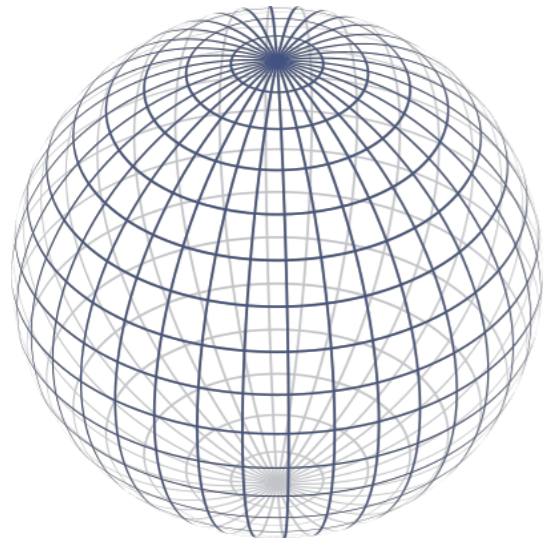
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- if **d =even**, then many possibilities (like $SU(d/2)$),
- if **d =odd and $d \neq 7$** : only $SO(d)$,
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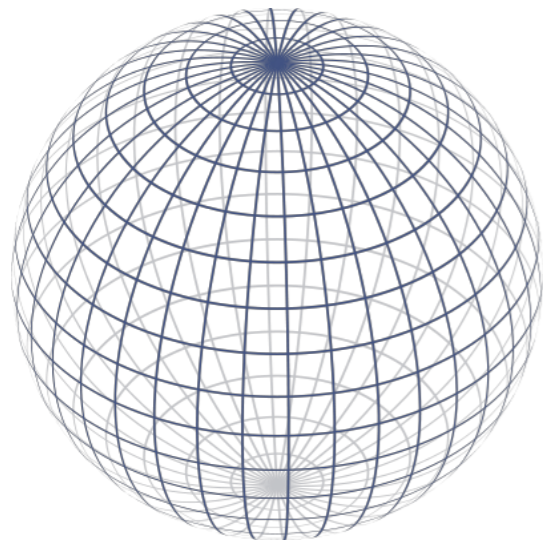
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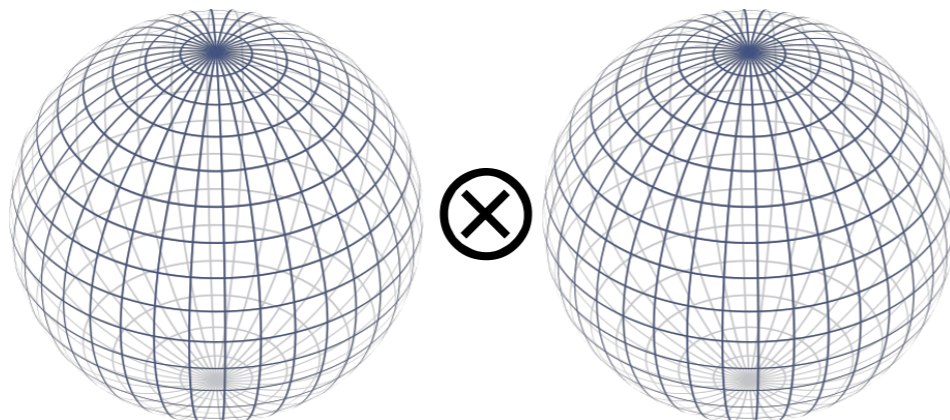
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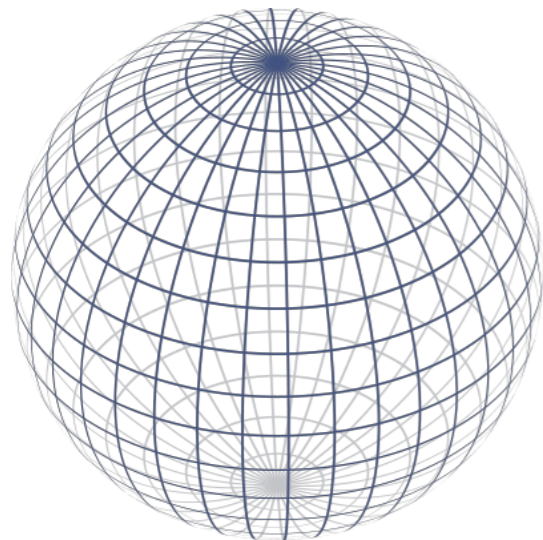
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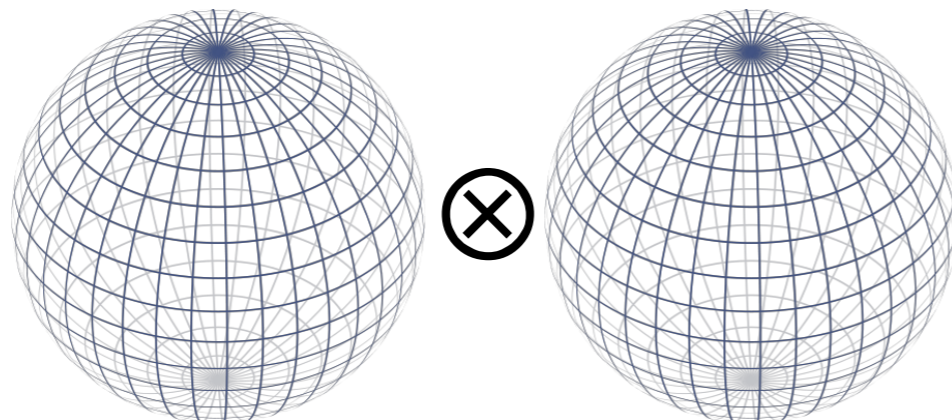
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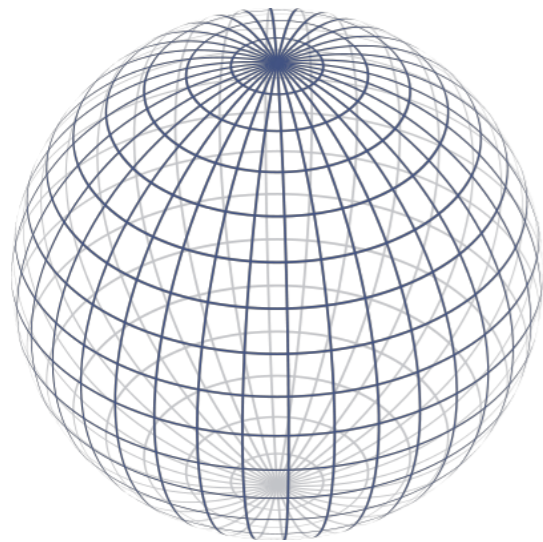
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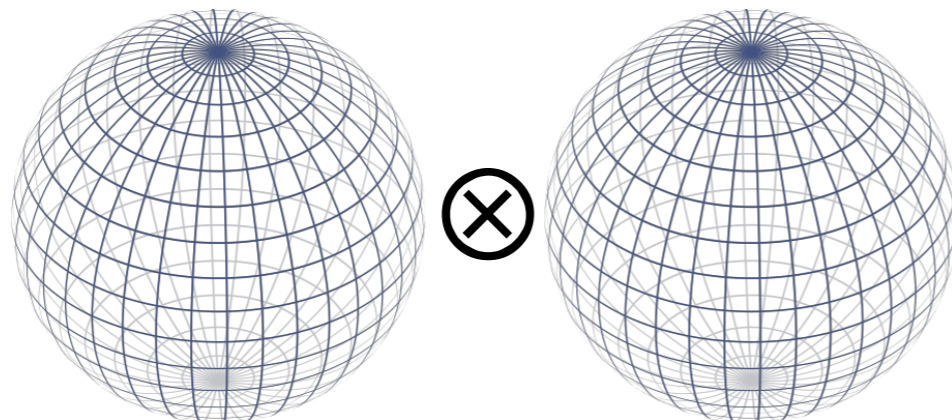
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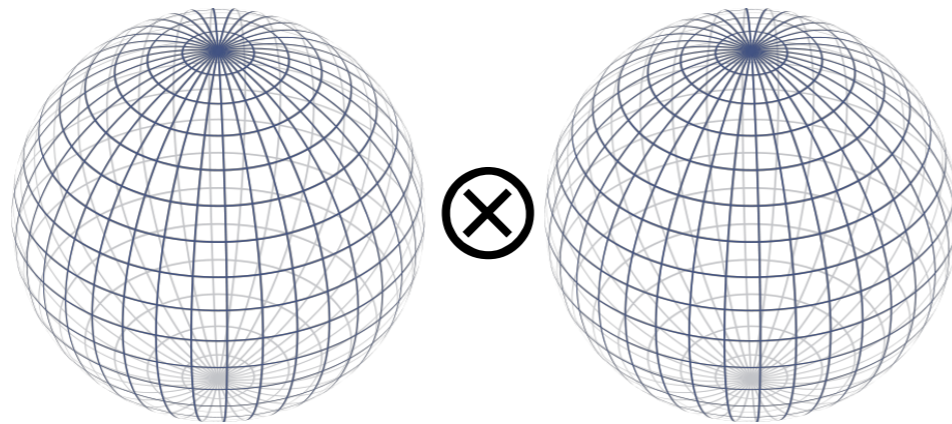
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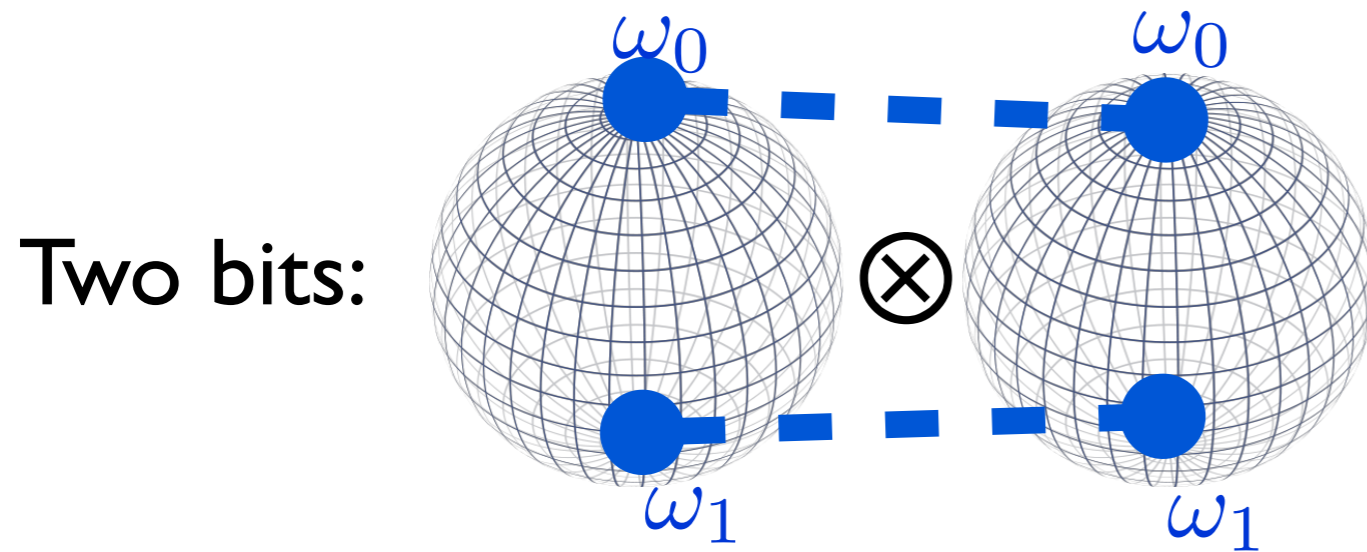
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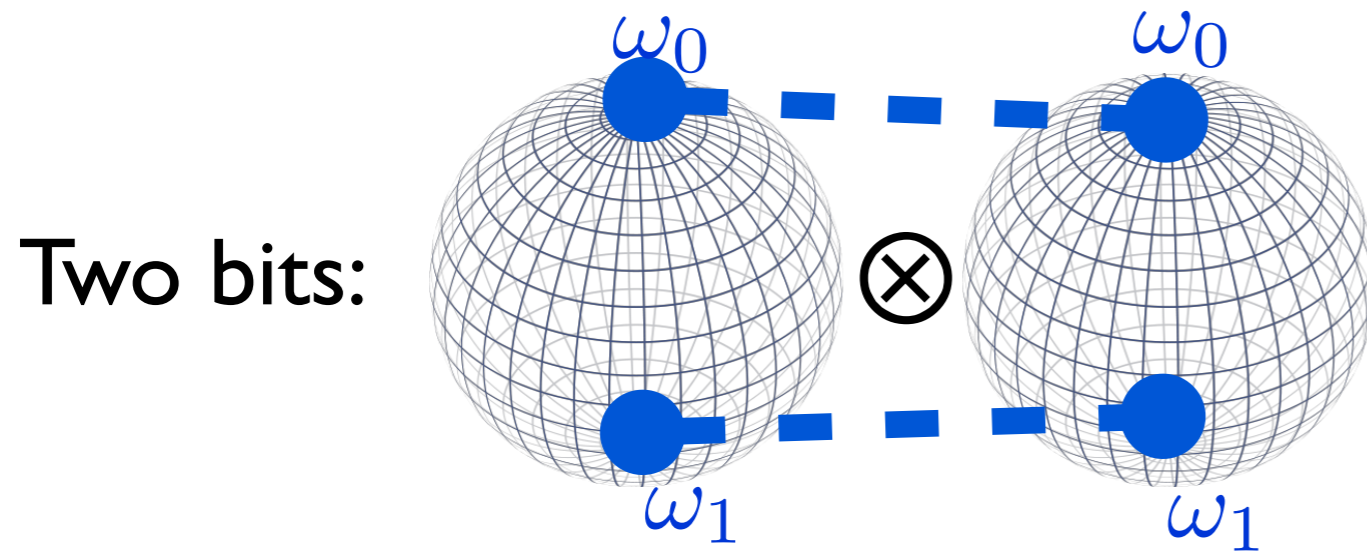
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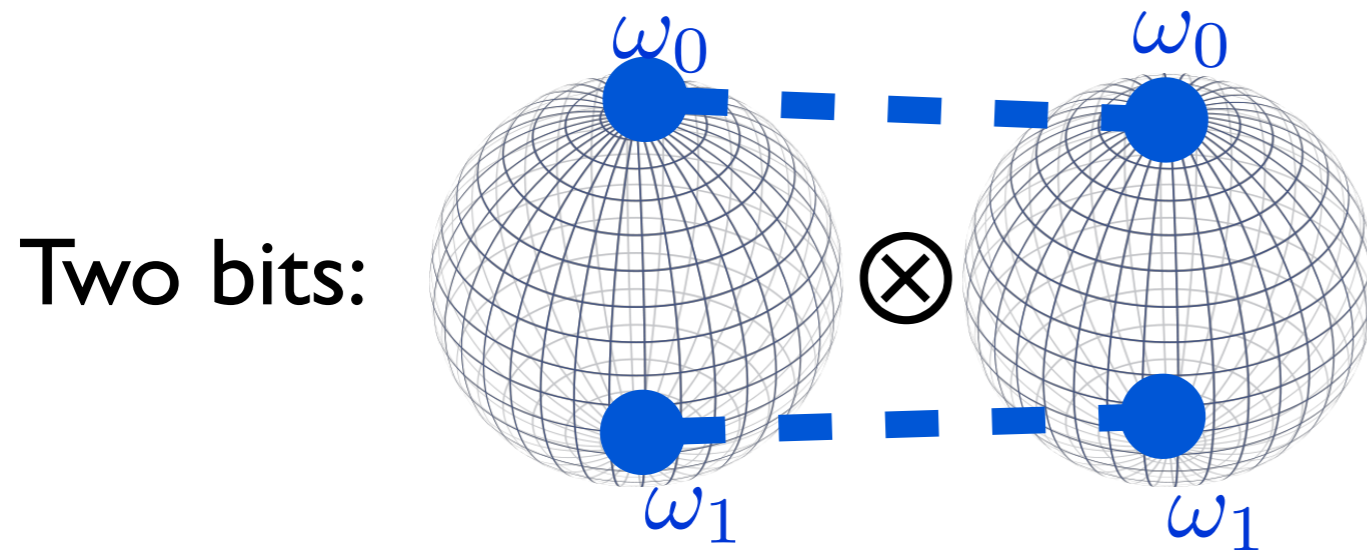


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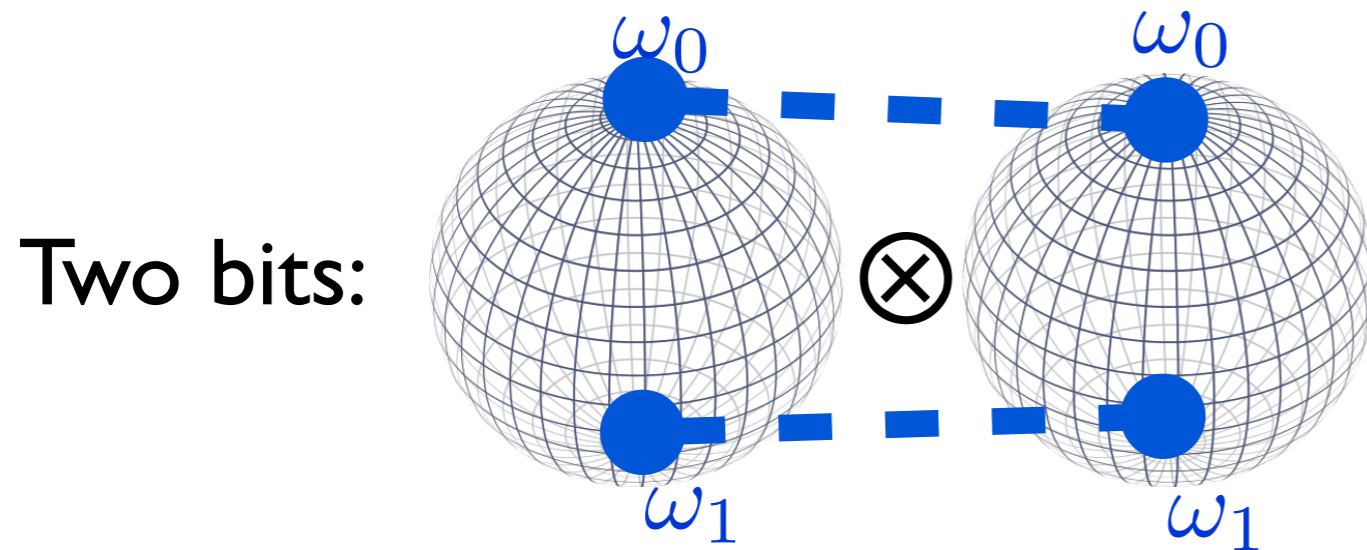
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4. Derivation of the Hilbert space formalism

Map 3-vectors to Hermitian matrices: $L(\omega) := \frac{1}{2} \left(\mathbf{1} + \sum_{i=1}^3 \omega_i \sigma_i \right)$

- Facts on **universal quantum computation**,
- **Wigner's theorem**
- some other tricks

prove:

4. Derivation of the Hilbert space formalism

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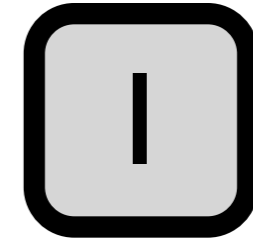
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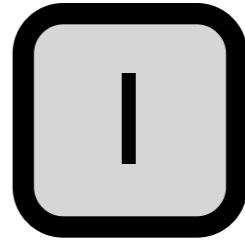
Theorem: Every theory satisfying Axioms I-V (rather than CPT) is equivalent to $(\Omega_N, \mathcal{G}_N)$, where

- Ω_N are the density matrices on \mathbb{C}^N ,
- \mathcal{G}_N is the group of unitaries, acting by conjugation,
- the measurements are exactly the POVMs.

5. Some new developments



5. Some new developments



The Axioms:

I. Local tomography

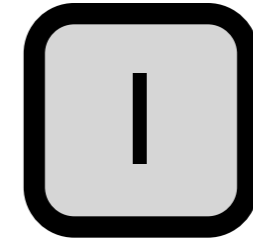
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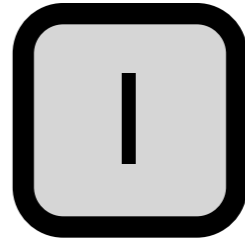
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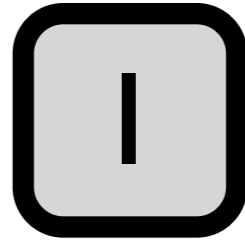


Conjecture: All state spaces satisfying I, II, IV are quantum systems.

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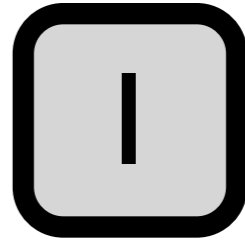
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LI. Masanes, MM, D. Pérez-García,
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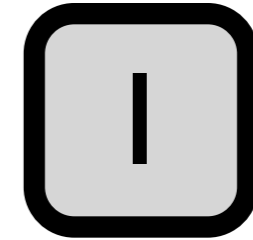
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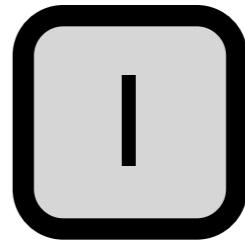


5. Some new developments



LI. Masanes, MM, R. Augusiak, and D. Pérez-García, *A digital approach to quantum theory*, arXiv:1208.0493

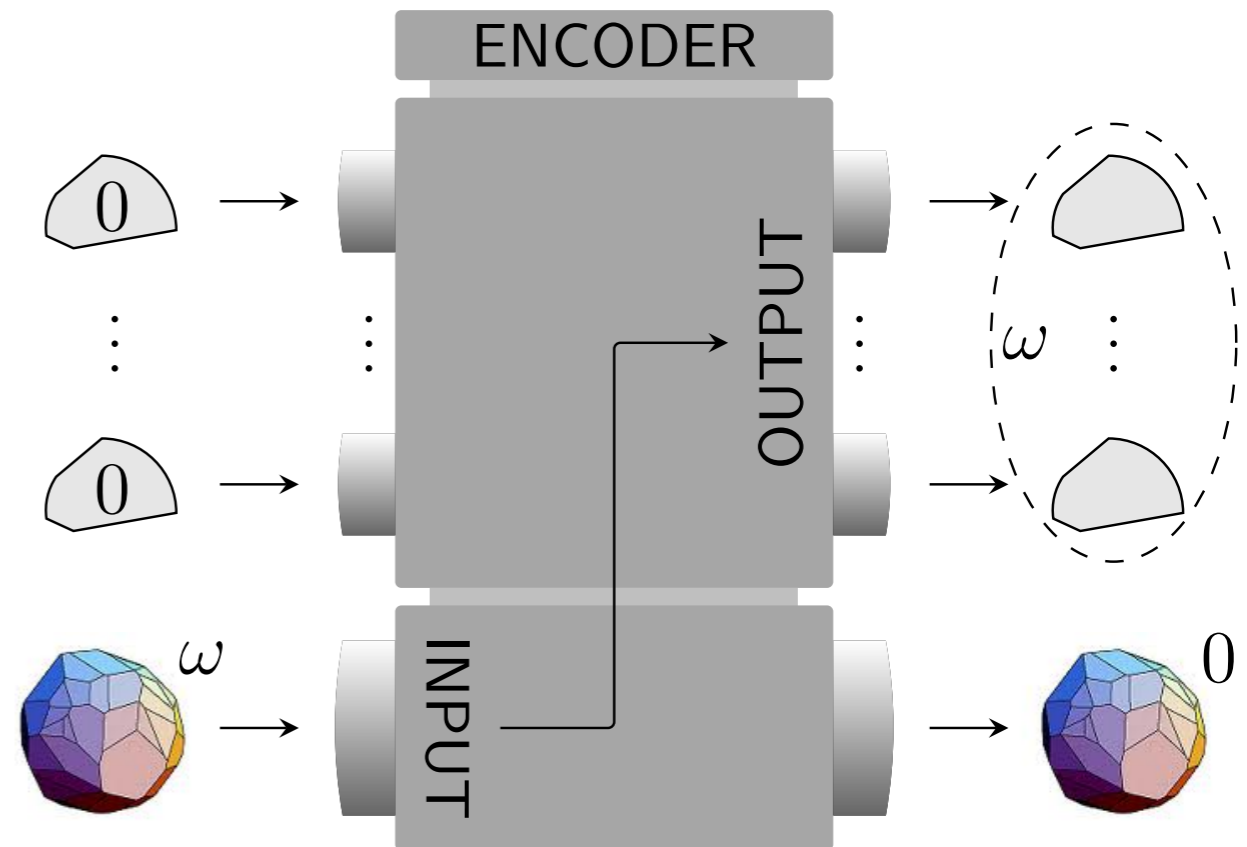
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Quantum theory follows from

- Local tomography,
- Continuous reversibility,
- Existence of an information unit:
there is “nice” binary system (“gbit”) such that the state of any system can be reversibly encoded in a sufficiently large number of gbits.



5. Some new developments



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Science 23 July 2010:

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DOI: 10.1126/science.1190545

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REPORT

Ruling Out Multi-Order Interference in Quantum Mechanics

Urbasi Sinha^{1,*}, Christophe Couteau^{1,2}, Thomas Jennewein¹, Raymond Laflamme^{1,3}, Gregor Weihs^{1,4,*}

[±](#) Author Affiliations

[↵](#)*To whom correspondence should be addressed. E-mail: usinha@iqc.ca, gregor.weihs@uibk.ac.at

ABSTRACT

Quantum mechanics and gravitation are two pillars of modern physics. Despite their success in describing the physical world around us, they seem to be incompatible theories. There are suggestions that one of these theories must be generalized to achieve unification. For example, Born's rule—one of the axioms of quantum mechanics—could be violated. Born's rule predicts that quantum interference, as shown by a double-slit diffraction experiment, occurs from pairs of paths. A generalized version of quantum mechanics might allow multipath (i.e., higher-order) interference, thus leading to a deviation from the theory. We performed a three-slit experiment with photons and bounded the magnitude of three-path interference to less than 10^{-2} of the expected two-path interference, thus ruling out third- and higher-order interference and providing a bound on the

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1



2



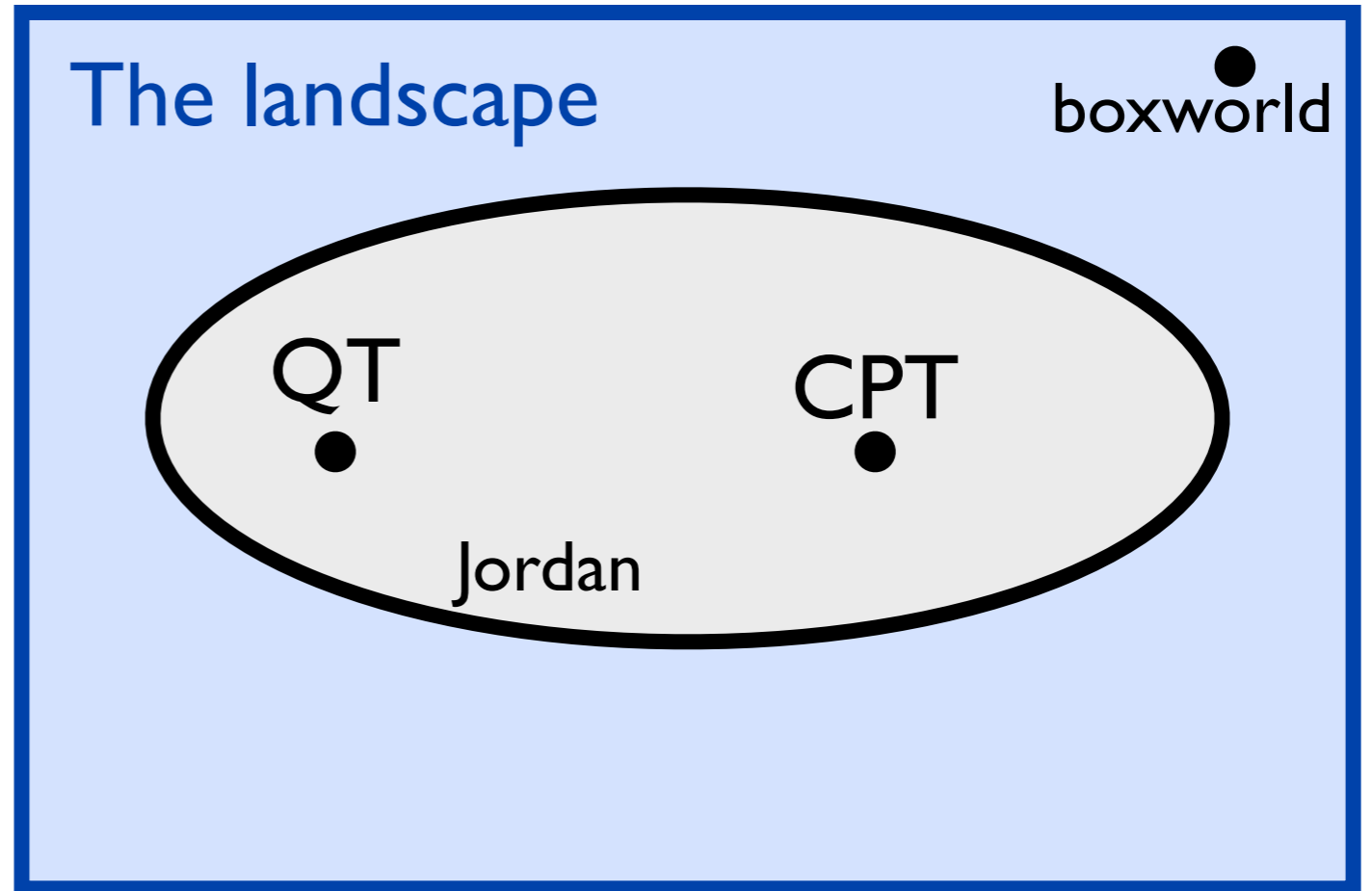
3



Quantum theory: $P_{123} - P_{12} - P_{23} - P_{13} + P_1 + P_2 + P_3 = 0$

⇒ no 3rd-order interference (R. Sorkin, Mod. Phys. Lett. **A9**, 3119 (1994))

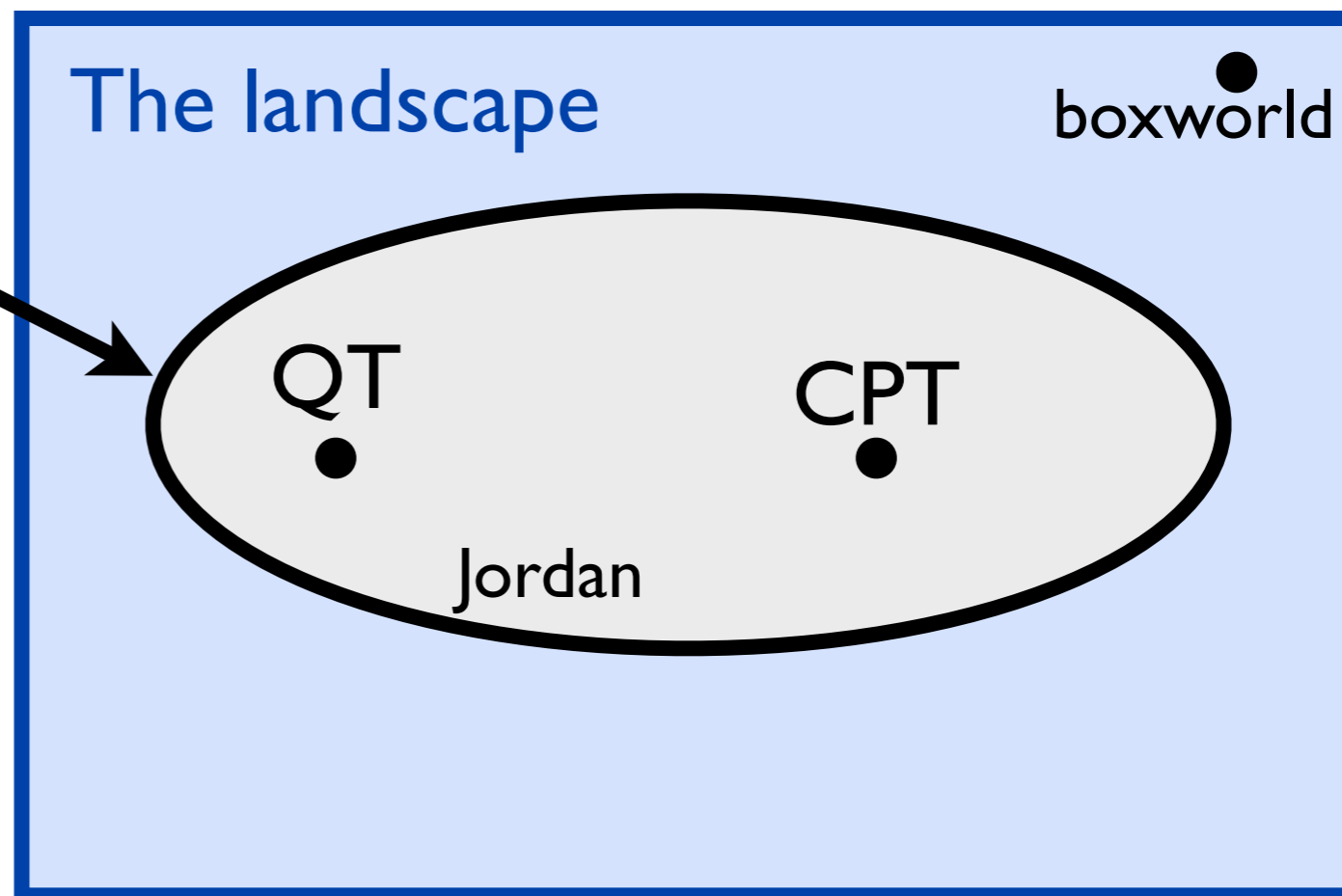
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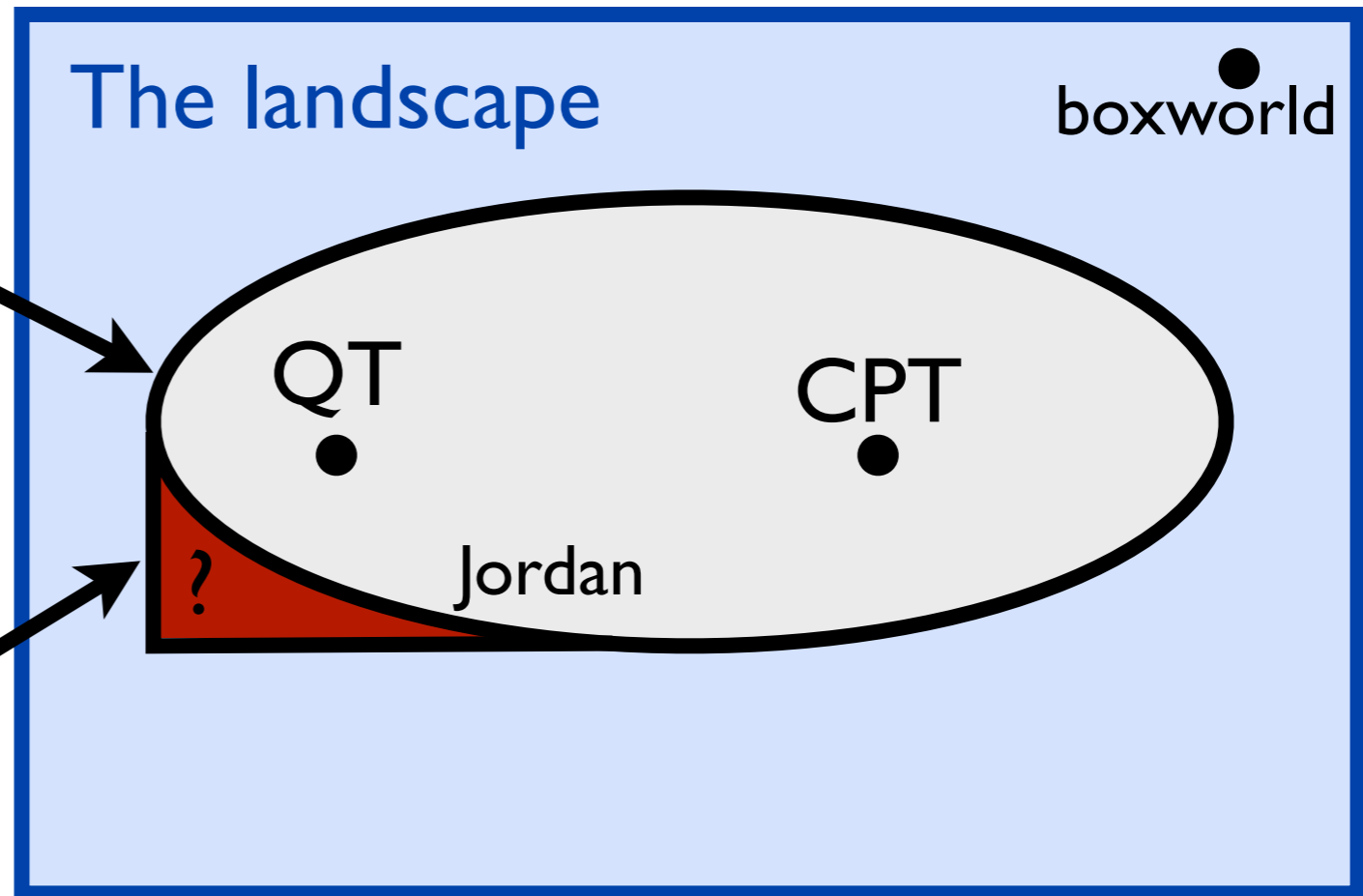
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(real, complex, quaternionic
QM, octonionic 3-level QM,
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2nd order, but no 3rd-order
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5. Some new developments



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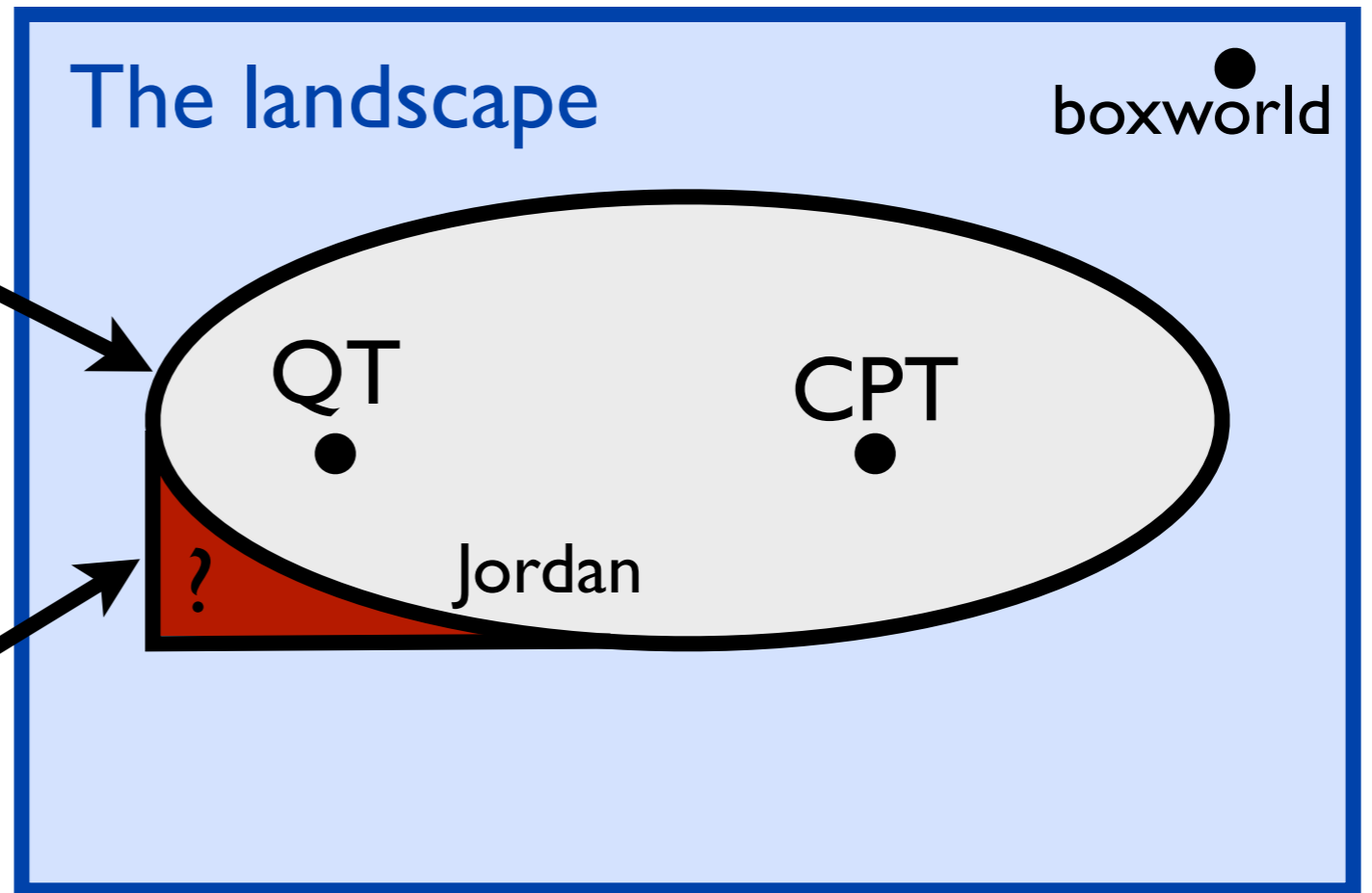


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5. Some new developments



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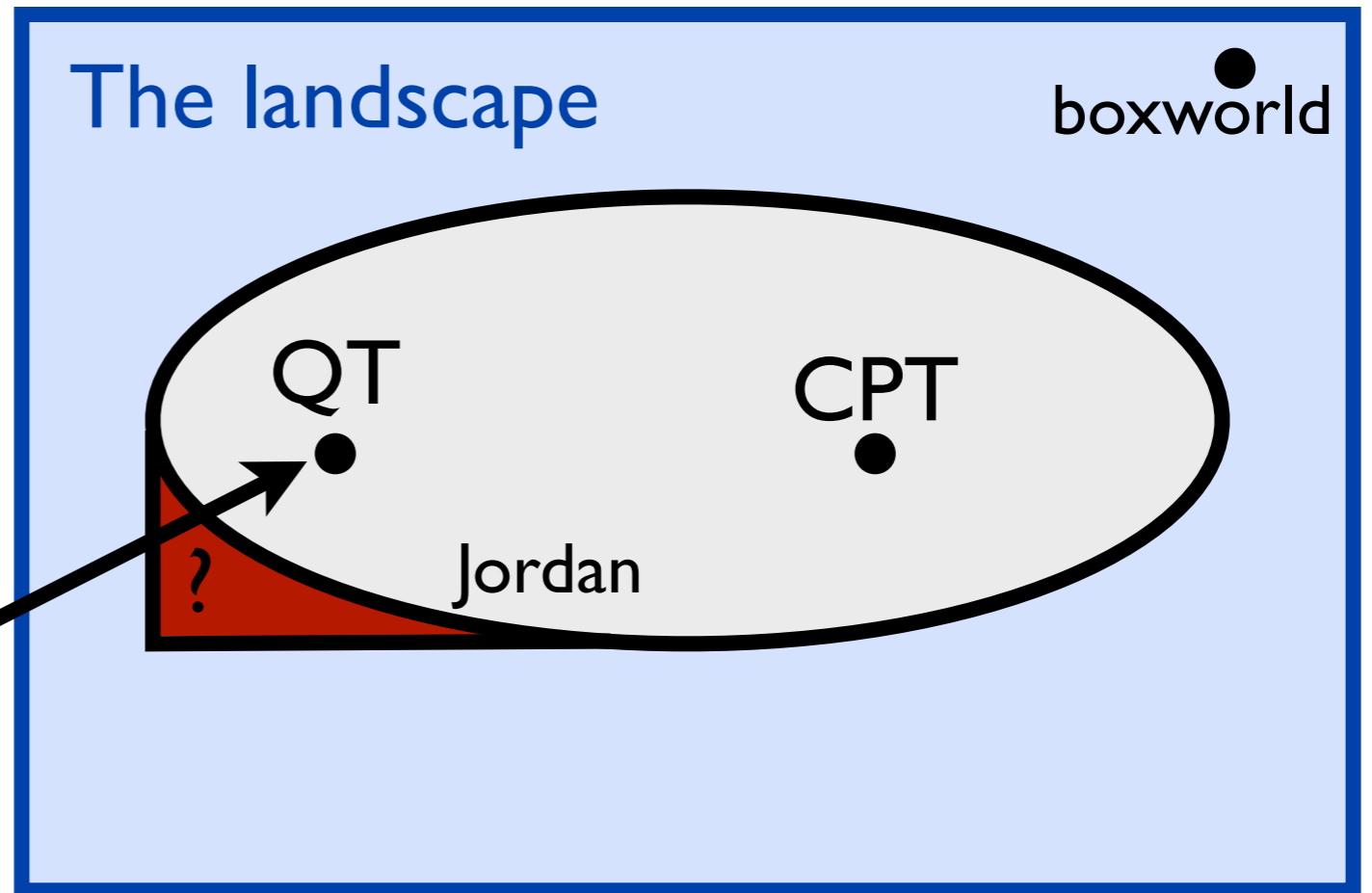
Def.: $\omega_1, \dots, \omega_n$ pure & perfectly distinguishable states are called a **frame**.

5. Some new developments



Joint work with Howard Barnum & Cozmin Ududec:

1. Every state is in the convex hull of some frame.
2. All frames of the same size are related by reversible transformations.
3. Local tomography
 \Rightarrow QT + CPT

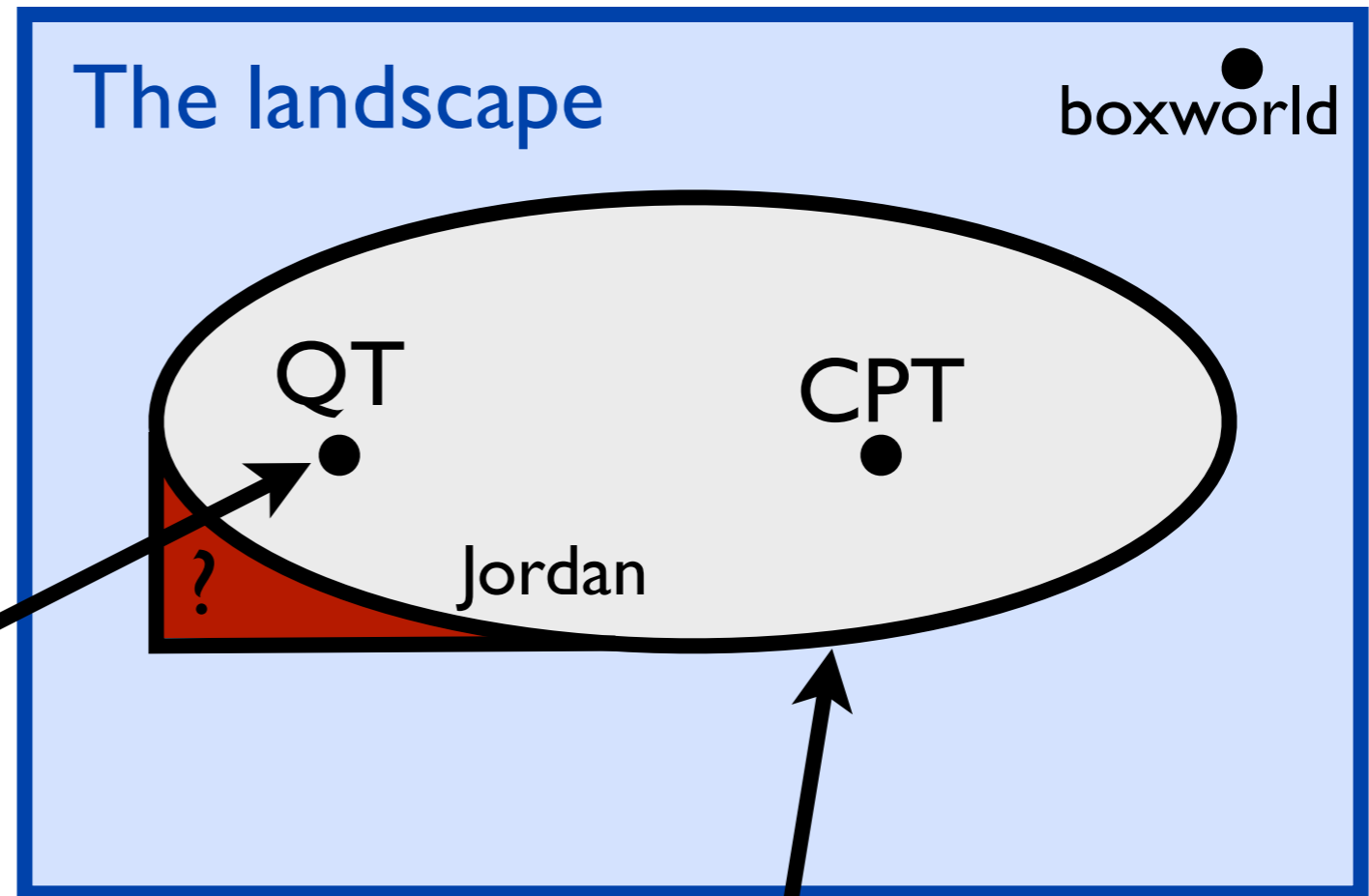


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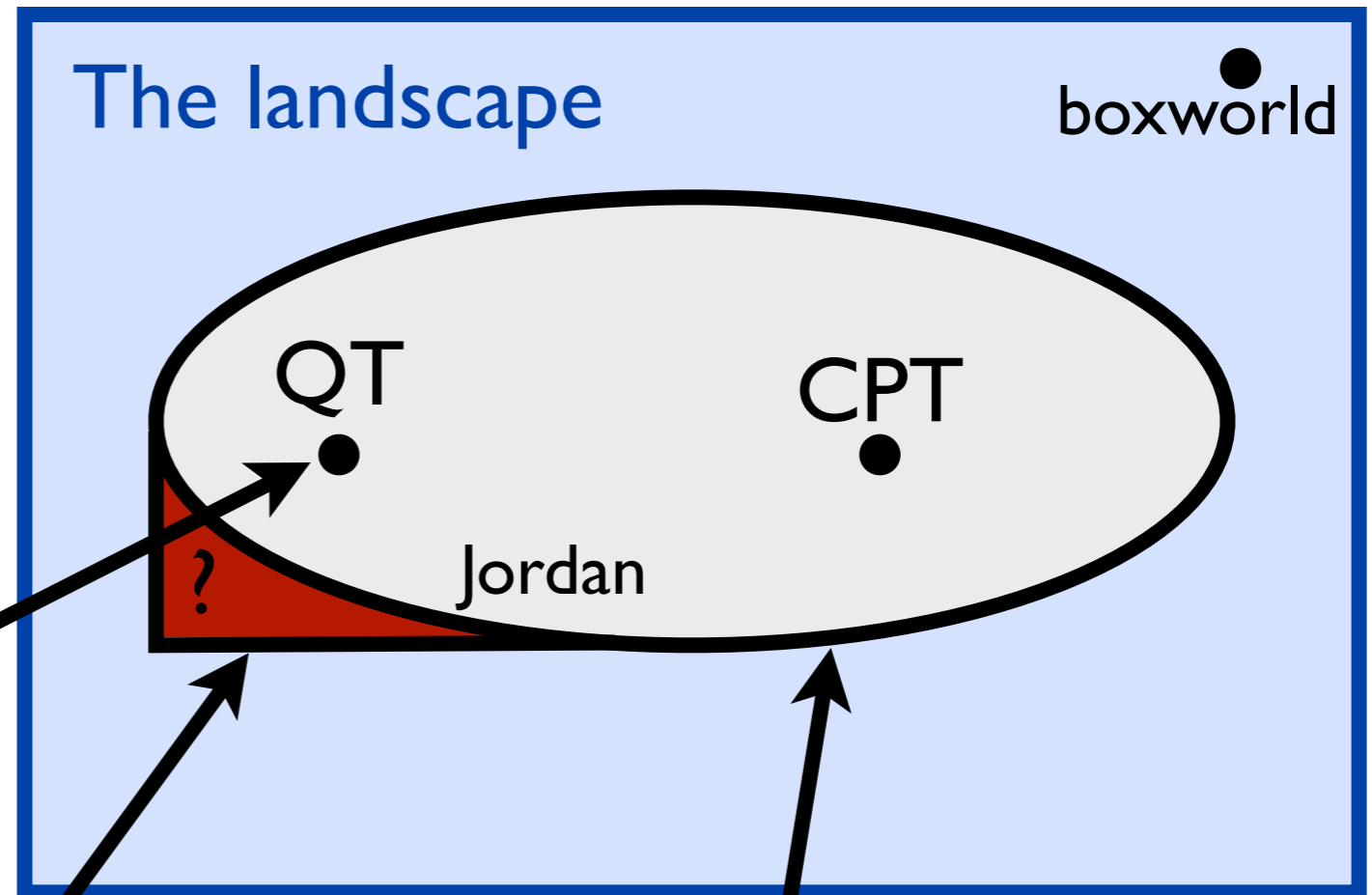
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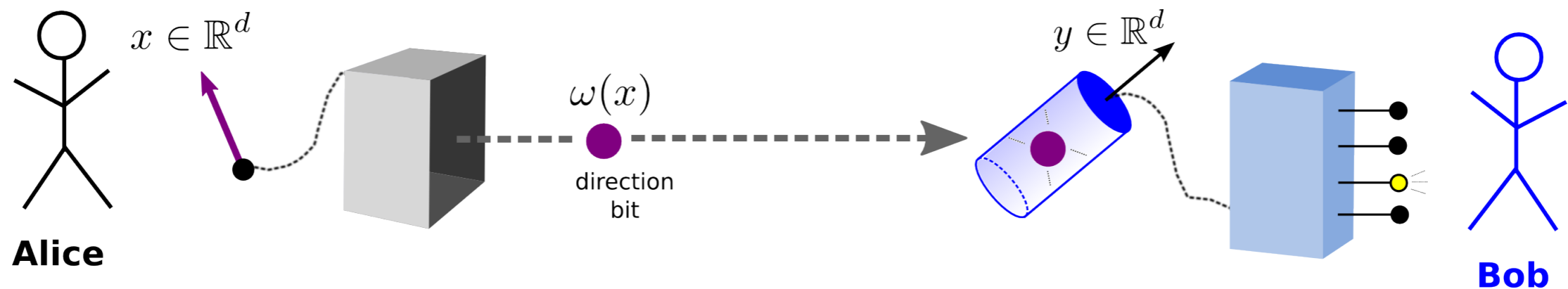
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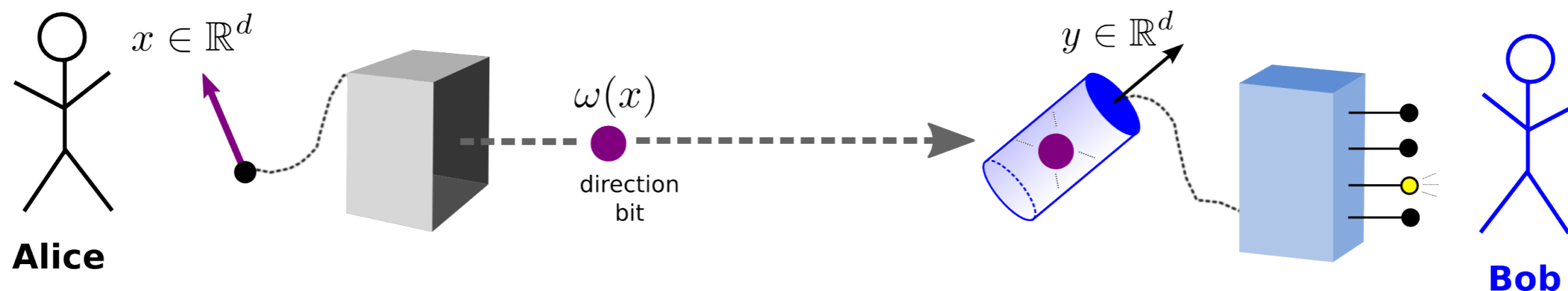
MM and LI. Masanes, *Three-dimensionality of space and the quantum bit: how to derive both from information-theoretic postulates*, arXiv:1206.0630



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Invitation: [Q+ Hangout Talk](#) (online), Tuesday, Oct 23.
More info later at: mattleifer.info

Reference:

[arXiv:1004.1483](#)

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More references:

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Thank you!