

Single-shot thermo, lecture 4

2.4. Majorization and Lorenz curves

Recall from last time: if $p, q \in \mathbb{R}^m$ are probability vectors, (i.e. $p_i, q_i \geq 0$, $\sum_i p_i = \sum_i q_i = 1$), then

$$p \succ q \Leftrightarrow \sum_{i=1}^k p_i^\downarrow \geq \sum_{i=1}^k q_i^\downarrow \text{ for all } k=1, \dots, m,$$

where p_i^\downarrow are the entries of p in non-increasing order.

Furthermore, write $p \xrightarrow{\text{noisy}} q$ if for every $\varepsilon > 0$ there exists q_ε with $\|q - q_\varepsilon\| < \varepsilon$ and a noisy classical operation D_ε with $D_\varepsilon(p) = q_\varepsilon$.

Lemma: Let $p, q \in \mathbb{R}^m$ be probability vectors of the same size. Then $p \xrightarrow{\text{noisy}} q$ if and only if $p \succ q$.

Proof. " \Rightarrow ": Suppose $p \xrightarrow{\text{noisy}} q$. Consider D_ε , then D_ε maps prob. vectors to prob. vectors, and $D_\varepsilon(1_m, \dots, 1_m) = (1_m, \dots, 1_m)$, D_ε linear. Problem 8, Exercise 2 $\Rightarrow D_\varepsilon$ is a bistochastic matrix, and thus $p \succ q_\varepsilon$ for every $q_\varepsilon := D_\varepsilon(p)$. But the set of prob. vectors majorized by p is topologically closed $\Rightarrow p \succ q$.

← Suppose $p > q$. Problem 6, Exercises 2 \Rightarrow
 q can be obtained from p as a mixture of permutations.

$$q = \sum_{j=1}^l \alpha_j \pi_j(p), \quad \alpha_j \geq 0, \quad \sum_j \alpha_j = 1, \quad \text{all } \pi_j \text{ permutations of entries of entries}$$

($\pi_j(p)_k = p_{\pi_j(k)}$)

Consider only special case $l=2$ ($l \geq 3$ analogous).

$$\Rightarrow q = \alpha \pi_1(p) + (1-\alpha) \pi_2(p), \quad 0 \leq \alpha \leq 1.$$

Construct a noisy operation with $p \rightarrow q_\epsilon$:

$$p = (p_1, \dots, p_m)$$

bring on max. mixed state \downarrow

$$p \otimes \gamma_N = \frac{1}{N} (\overbrace{p_1, \dots, p_1}^{N_1}, \overbrace{p_2, \dots, p_2}^{N_2}, \dots, \overbrace{p_m, \dots, p_m}^{N_m})$$

($\frac{1}{N}, \dots, \frac{1}{N}$)

global permutation \downarrow

$$\pi[p \otimes \gamma_N] := \frac{1}{N} (\overbrace{P_{\pi_1(1)}, \dots, P_{\pi_1(N_1)}}^{N_1}, \overbrace{P_{\pi_2(1)}, \dots, P_{\pi_2(N_2)}}^{N_2}, \dots, \overbrace{P_{\pi_1(N_1)}, \dots, P_{\pi_1(N_1)}}^{N_1}, \overbrace{P_{\pi_2(N_2)}, \dots, P_{\pi_2(N_2)}}^{N_2})$$

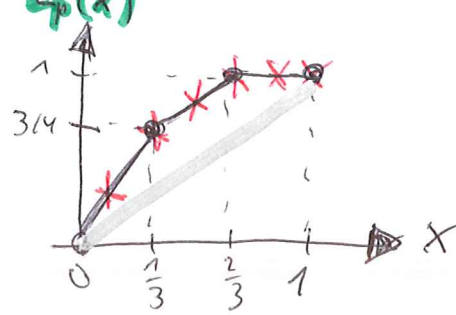
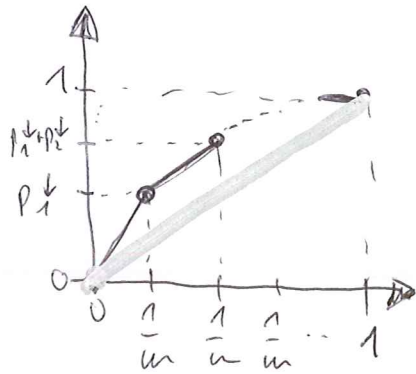
marginal (sum over ancilla) \downarrow
 \rightarrow sum over N -blocks

$$(\pi[p \otimes \gamma_N])_{\text{out}} = \frac{1}{N} (N_1 P_{\pi_1(1)} + N_2 P_{\pi_2(1)}, \dots, N_1 P_{\pi_1(N_1)} + N_2 P_{\pi_2(N_1)})$$

$$= \frac{N_1}{N} \pi_1(p) + \frac{N_2}{N} \pi_2(p).$$

By making $N \rightarrow \infty$ and choosing $N_i = N_1(N)$ appropriately, we can approximate α to arbitrary accuracy. \square

Graphical tool: Lorentz curve: linear interpolation of Prob. vector $p \in \mathbb{R}^m$. Ex: $p = (0, \frac{3}{4}, \frac{1}{4})$ $(\sum_{i=1}^k p_i^\downarrow)_{k=0, \dots, m}$



Observation: p and $p \otimes \gamma_N = p \otimes (\frac{1}{N}, \dots, \frac{1}{N})$ have identical Lorentz curves

e.g. $N=2$: $p \otimes (\frac{1}{2}, \frac{1}{2}) = (0, 0, \frac{3}{8}, \frac{3}{8}, \frac{1}{8}, \frac{1}{8})$

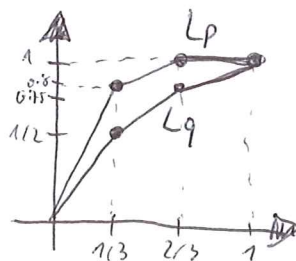
Lemma: $p \xrightarrow{\text{noisy}} q$ if and only if $L_p(x) \geq L_q(x)$ for all x , even if p and q do not have the same dimensionalities (i.e. $p \in \mathbb{R}^m, q \in \mathbb{R}^n$).

Proof: If $m=n$ then $L_p \geq L_q \Leftrightarrow p \geq q$ by definition.

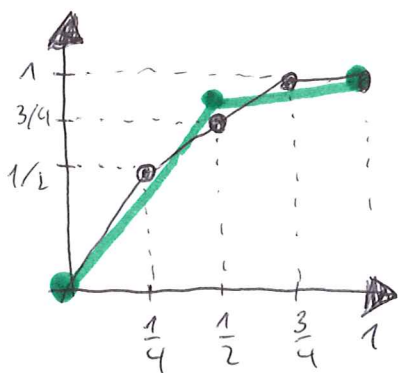
Otherwise: $p \xrightarrow{\text{noisy}} q \Leftrightarrow p \otimes \gamma_n \xrightarrow{\text{noisy}} q \otimes \gamma_m$
 $\Leftrightarrow L_{p \otimes \gamma_n} \geq L_{q \otimes \gamma_m} \Leftrightarrow L_p \geq L_q$

All Lorentz curves are concave, and lie above the Lorentz curve of $\gamma_n = (\frac{1}{n}, \dots, \frac{1}{n})$.

$\begin{pmatrix} 1/5 \\ 0 \\ 4/5 \end{pmatrix} \succ \begin{pmatrix} 1/2 \\ 1/4 \\ 1/4 \end{pmatrix}$
 $\underbrace{\quad}_p \quad \underbrace{\quad}_q$



p is a more valuable resource than q .



$$p = \left(\frac{1}{2}, \frac{1}{4}, \frac{1}{4}, 0\right),$$

$$q = \left(\frac{4}{5}, \frac{1}{5}\right)$$

$$\Rightarrow p \not\rightarrow q \text{ and } q \not\rightarrow p$$

Neither resource is "more valuable" than the other.

2.5. Nonuniformity monotones

Let \mathcal{P}_n be the set of probability vectors in \mathbb{R}^n .

Def. A function $F: \mathcal{P}_n \rightarrow \mathbb{R}$ is Schur-convex (on \mathcal{P}_n) if $p \succ q \Rightarrow F(p) \geq F(q)$

(or Schur-concave if $p \succ q \Rightarrow F(p) \leq F(q)$).

Example: $H_\infty(p) := -\log \max_i p_i$

$$p \succ q \Rightarrow p_1^\downarrow \geq q_1^\downarrow \Rightarrow \max_i p_i \geq \max_i q_i$$

$$\Rightarrow H_\infty(p) \leq H_\infty(q)$$

H_∞ is Schur-concave.

We will now prove the intuition that "noisy operations increase the entropy".

Lemma: For every convex (concave) function $f: [0,1] \rightarrow \mathbb{R}$, the function $F: \mathcal{P}_n \rightarrow \mathbb{R}$ defined by

$$F(p) := \sum_{i=1}^n f(p_i)$$

is Schur-convex (Schur-concave).

(f convex $\Leftrightarrow f(\lambda x + (1-\lambda)y) \leq \lambda f(x) + (1-\lambda)f(y)$
 $\forall \lambda \in [0,1]$.)
 If f is twice differentiable, then $f'' \geq 0 \Rightarrow$ ~~convex~~ f convex.

Proof: exercise.

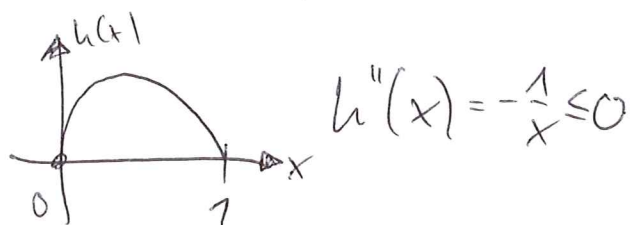
Now consider Shannon entropy

$$H(p) := - \sum_{i=1}^n p_i \log p_i \quad (0 \log 0 := 0).$$

$$= \sum_{i=1}^n h(p_i), \quad h(p_i) = -p_i \log p_i.$$

h concave $\Rightarrow H$ Schur-concave.

$$p \succ q \Rightarrow H(p) \leq H(q)$$



\Rightarrow Noisy operators $D: \mathbb{R}^n \rightarrow \mathbb{R}^n$ ~~never~~ never decrease the entropy.

Reminiscent of "Second Law".

For different dimensionalities: $p \in \mathbb{R}^m, q \in \mathbb{R}^n$

$$\boxed{p \xrightarrow{\text{noisy}} q} \Leftrightarrow p \otimes \delta_n \xrightarrow{\text{noisy}} q \otimes \delta_n \Leftrightarrow p \otimes \delta_n \succ q \otimes \delta_n$$

$$\Rightarrow H(p \otimes \delta_n) \leq H(q \otimes \delta_n)$$

$$\Rightarrow H(p) + \log n \leq H(q) + \log n$$

$$\Rightarrow I(p) \geq I(q)$$

for $I(p) := \log d_p - H(p)$ "negentropy"
 $d_p = \dim p$

This is necessary, but is it sufficient? No.

Recall $p = (\frac{1}{2}, \frac{1}{4}, \frac{1}{4}, 0)$ $q = (\frac{4}{5}, \frac{1}{5})$ (for now, $\log = \log_2$)

$$\Rightarrow I(p) = \frac{1}{2}, \quad I(q) \approx 0.28 \quad (\text{and } q \xrightarrow{\text{noise}} p).$$

But [redacted] $p \not\xrightarrow{\text{noise}} q$. More generally

$\xrightarrow{\text{noise}}$ is a preorder, but [redacted] I (like every other monotone) introduces an order:

$$p \succeq q \Leftrightarrow I(p) \geq I(q), \text{ all states comparable.}$$

Def: A nonuniformity monotone is a function

$$F: \bigcup_{n \in \mathbb{N}} \mathcal{P}_n \rightarrow \mathbb{R} \text{ such that}$$

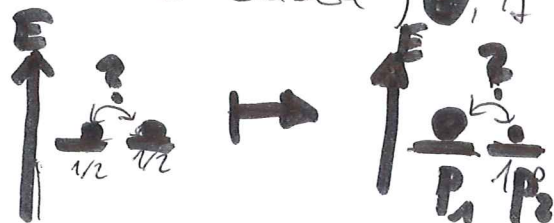
$$p \xrightarrow{\text{noise}} q \Rightarrow F(p) \geq F(q).$$

(Example: negentropy I .)

Because of $(*)$, no single monotone can fully characterize the possible state transitions.

2.6. Extracting work from states: distillable nonuniformity and uniformity of formation

How much work does it cost to create a state (as a Landauer eraser) [redacted] , if we want to succeed with probability one?



We will show:

$$(\log d_p - H_\infty(p)) / k_B T \ln 2$$

How much work can we extract from a state (in a Szilard engine)?

We show: $(\log d_p - H_{\bullet}(p)) k_B T \ln 2$

where

$$\underbrace{H_{\infty}}_{\text{"min entropy"}} \leq \underbrace{H}_{\text{Shannon entropy}} \leq \underbrace{H_0}_{\text{"max entropy"}}$$

- Fundamental irreversibility: have to invest more work than we can get out of ~~the state~~ the state later on.
- But we will show: reversibility is recovered in the thermodynamic limit (work extraction "on average" for many copies).

Def. (Rényi entropies): Let $p \in \mathbb{R}^n$ be any probability vector.

For $\alpha > 0, \alpha \neq 1$, define $H_{\alpha}(p) := \frac{1}{1-\alpha} \log \sum_{i=1}^n p_i^{\alpha}$

By taking the limits, we also get

$$H_0(p) := \log \text{rank}(p), \text{ where } \text{rank}(p) := \# \text{ of nonzero entries of } p,$$

$$H_1(p) = H(p) = - \sum_{i=1}^n p_i \log p_i \quad \text{Shannon entropy}$$

$$H_{\infty}(p) = -\log \max_i p_i.$$

Analogously for quantum states S : $H_{\alpha}(S) = \frac{1}{1-\alpha} \log \text{tr}(S^{\alpha})$

$$H_0(S) = \log \text{rank}(S), \quad H_1(S) = S(S) \quad \text{von Neumann entropy}$$

$$H_{\infty}(S) = -\log \underbrace{\lambda_{\max}(S)}_{\text{largest eigenvalue}}$$