

Single-shot thermo, lecture 12

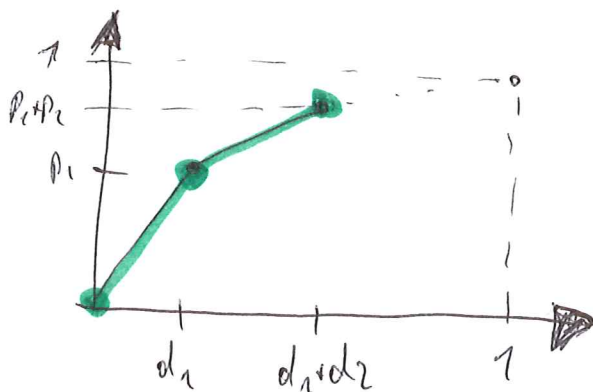
①

- Sorry, exercise sheet 9 went online quite late
→ can give back until Thursday, 17:00 (to me or one of the secretaries in Philosophenweg 19).
- Exercise sheet 10 will be the last one.

Recap last lecture:

- For blockdiagonal states $[H, S] = 0$: can reduce to the classical case.
- Equivalent are, classically:
 - $p \xrightarrow{\text{thermal}} p'$, if initial and final Hamiltonians are such that the Gibbs states are d and d' ,
 - $(p, d) \succ (p', d')$, i.e. \exists stochastic B with $Bp = p'$ and $Bd = d'$.
- Thermal Lorenz curve of p everywhere on or above that of p' .

Sort such that $\frac{p_1}{d_1} \geq \frac{p_2}{d_2} \geq \dots$ & let



• for every convex fn. g ,

$$\sum_i d_i g\left(\frac{p_i}{d_i}\right) \geq$$

$$\sum_i d'_i g\left(\frac{p'_i}{d'_i}\right).$$

→ starting point for Rényi divergences

3.7. The Rényi divergence

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Consider last condition for $g(x) := \pm x^\alpha$

Def (Relative Rényi entropies)

If $p, q \in \mathbb{R}^n$ are probability vectors, and $\alpha \in \mathbb{R} \setminus \{0, 1\}$,

$$D_\alpha(p \parallel q) := \frac{\text{sgn } \alpha}{\alpha - 1} \log \sum_{i=1}^n p_i^\alpha q_i^{1-\alpha}.$$

Furthermore,

$$D_0(p \parallel q) := \lim_{\alpha \rightarrow 0^+} D_\alpha(p \parallel q) = -\log \sum_{\substack{i: \\ p_i \neq 0}} q_i$$

cf. home-work

$$\left\{ \begin{array}{l} D_1(p \parallel q) := \lim_{\alpha \rightarrow 1} D_\alpha(p \parallel q) = \sum_{i=1}^n p_i (\log p_i - \log q_i) \\ \text{standard relative entropy } D(p \parallel q) \end{array} \right.$$

$$D_\infty(p \parallel q) := \lim_{\alpha \rightarrow \infty} D_\alpha(p \parallel q) = \log \max_i \frac{p_i}{q_i}$$

$$D_{-\infty}(p \parallel q) := \lim_{\alpha \rightarrow -\infty} D_\alpha(p \parallel q) = D_\infty(q \parallel p).$$

Thm.: If $(p, d) \succ (p', d')$ then $D_\alpha(p \parallel d) \geq D_\alpha(p' \parallel d')$
 $\forall \alpha$

Notation: $D_{\min} = D_0$, $D = D_1$, $D_{\max} = D_\infty$

$$\Rightarrow D_{\min} \leq D \leq D_{\max}.$$

Caution: $H_{\min} = H_\infty$, $H = H_1$, $H_{\max} = H_0$ (entropies)

$$\Rightarrow H_{\min} \leq H \leq H_{\max}.$$

3.8. Work of formation and the max - relative entropy

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Suppose we want to create a system

How much work does that cost?

→ start with energy eigenstate $|E\rangle$ of another Hamiltonian H .

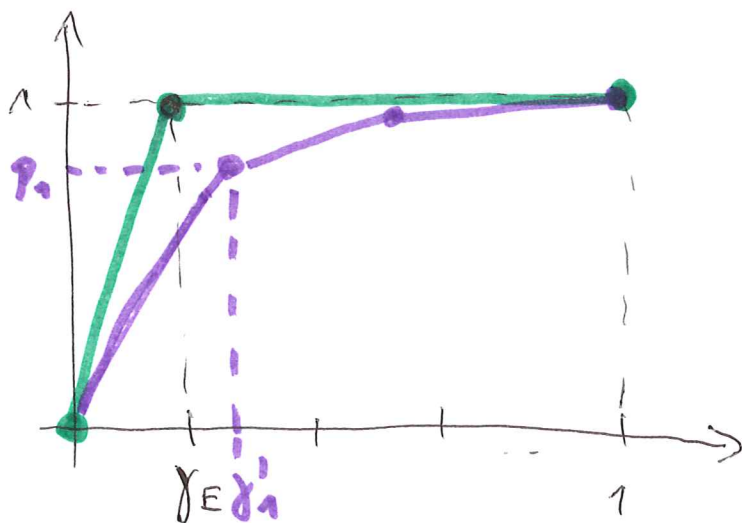
Want $(|E\rangle, H) \xrightarrow{\text{thermal}} (p, H')$.

What's the smallest possible E ?

Thermo-Lowz curve of $|E\rangle$ ($p_E = 1$, all other $p_i = 0$)

$$P = (p, H')$$

prob. vector, resp. eigenvalues of block-diag. density matrix



curve of p

We only have to check the condition for the initial

slope: $\frac{1}{\gamma_E} \geq \frac{p_1}{\gamma'_1} = \max_i \frac{p_i}{\gamma'_i} = e^{D_\infty(p \parallel \gamma')}$

$$\gamma_E = e^{-\beta E} / Z \quad \leftarrow \text{partition f. of Hamiltonian } H$$

$$\Rightarrow E \geq \frac{1}{\beta} D_\infty(p \parallel \gamma') - \frac{1}{\beta} \log Z. \quad (*)$$

(4)

If energy levels of H sufficiently dense, we can get (close to) equality.

Work of formation for thermal state:

$$\log Z = \log \sum_i e^{-\beta E_i} \geq \log e^{-\beta E_{\min}} \geq -\beta E_{\min}$$

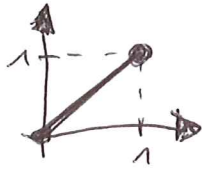
$$\Rightarrow -\frac{1}{\beta} \log Z \leq E_{\min}$$

If $p = \gamma'$ then $D_{\infty}(p \| \gamma') = 0$ (Rényi divergences have this "distance measure" property) \Rightarrow

$$E \geq E_{\min} \geq -\frac{1}{\beta} \log Z,$$

so (*) is automatically satisfied. Thermal states are "free"!

Lorentz curve of γ' :



(*) is equivalent to

$$D_{\infty}(\underbrace{|E \times E|}_{\text{as prob. vector}} \| \gamma) \geq D_{\infty}(p \| \gamma')$$

$$\Leftrightarrow F_{\infty}(|E \times E|) - F_{\infty}(\gamma) \geq F_{\infty}(p) - F_{\infty}(\gamma')$$

$$\text{where } F_{\infty}(p) = \frac{1}{\beta} D_{\infty}(p \| \gamma') - \frac{1}{\beta} \log Z'$$

(if the system is $\mathcal{P} = (p, H')$ with γ' Gibbs state)

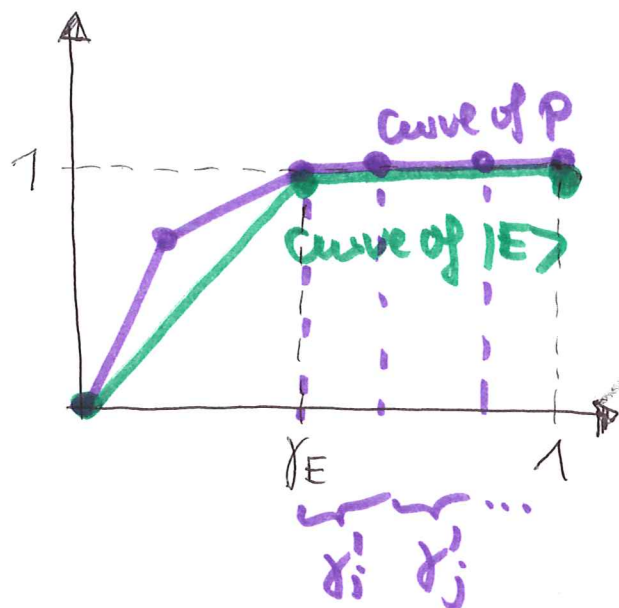
"max-free energy" $F_{\infty} =: F_{\max}$

3.9. Extractable work and the min-relative entropy

(5)

Suppose we are given a system $P = (p, H')$ and want to create a system $(|E\rangle, H)$.

What's the largest possible E we can get?



Need "flat tail" $\hat{=}$ zero entries of P .

green tail length \leq pink tail length

$$\Rightarrow 1 - x_E = 1 - \frac{e^{-\beta E}}{Z} \leq \sum_{\substack{i: \\ p_i = 0}} \frac{e^{-\beta E'_i}}{Z'} = 1 - \sum_{\substack{i: \\ p_i \neq 0}} x'_i$$

$$= 1 - e^{-D_0(p \| x')}$$

$$\Rightarrow E \leq \frac{1}{\beta} D_0(p \| x') - \frac{1}{\beta} \log Z \quad (*)$$

If we choose H suitably, we can achieve equality.

3.10. Allowing finite error probability Smoothing

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Introduce "smoothed Rényi divergences":

$$D_0^\epsilon(p \| q) := \max_{p': D(p, p') \leq \epsilon} D_0(p' \| q),$$

$$D_\infty^\epsilon(p \| q) := \min_{p': D(p, p') \leq \epsilon} D_\infty(p' \| q).$$

Approximate work of formation: Want

$$(|E\rangle, H) \xrightarrow{\text{Thermal}} (p', H') \text{ with } D(p, p') \leq \epsilon.$$

Minimal work needed is therefore

$$W_{\text{form}}^\epsilon = \min_{p': D(p, p') \leq \epsilon} \left(\frac{1}{\beta} D_\infty(p' \| \gamma') - \frac{1}{\beta} \log Z \right)$$

$$\Rightarrow W_{\text{form}}^\epsilon = \frac{1}{\beta} D_\infty^\epsilon(p \| \gamma') - \frac{1}{\beta} \log Z$$

Approximate extractable work:

$$\text{Want } (p, H') \xrightarrow{\text{Thermal}} (q, H) \text{ with } D(q, |E\rangle\langle E|) \leq \epsilon.$$

$$\text{Choose } p' \text{ with } D_0^\epsilon(p \| \gamma') = D_0(p' \| \gamma'), D(p, p') \leq \epsilon.$$

\Rightarrow Can obtain $(|E\rangle\langle E|, H)$ by thermal op. ϕ from p' , where

$$E \leq \frac{1}{\beta} D_0(p' \| \gamma') - \frac{1}{\beta} \log Z = \frac{1}{\beta} D_0^\epsilon(p \| \gamma') - \frac{1}{\beta} \log Z.$$

$$\Rightarrow D(\phi(p), |E\rangle\langle E|) = D(\phi(p), \phi(p')) \leq D(p, p') \leq \epsilon.$$

$$\Rightarrow W_{\text{extr}}^{\varepsilon} \geq \frac{1}{\beta} D_0^{\varepsilon}(p \| \gamma') - \frac{1}{\beta} \log Z.$$

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Exact expression is a bit harder to get;
maybe homework.